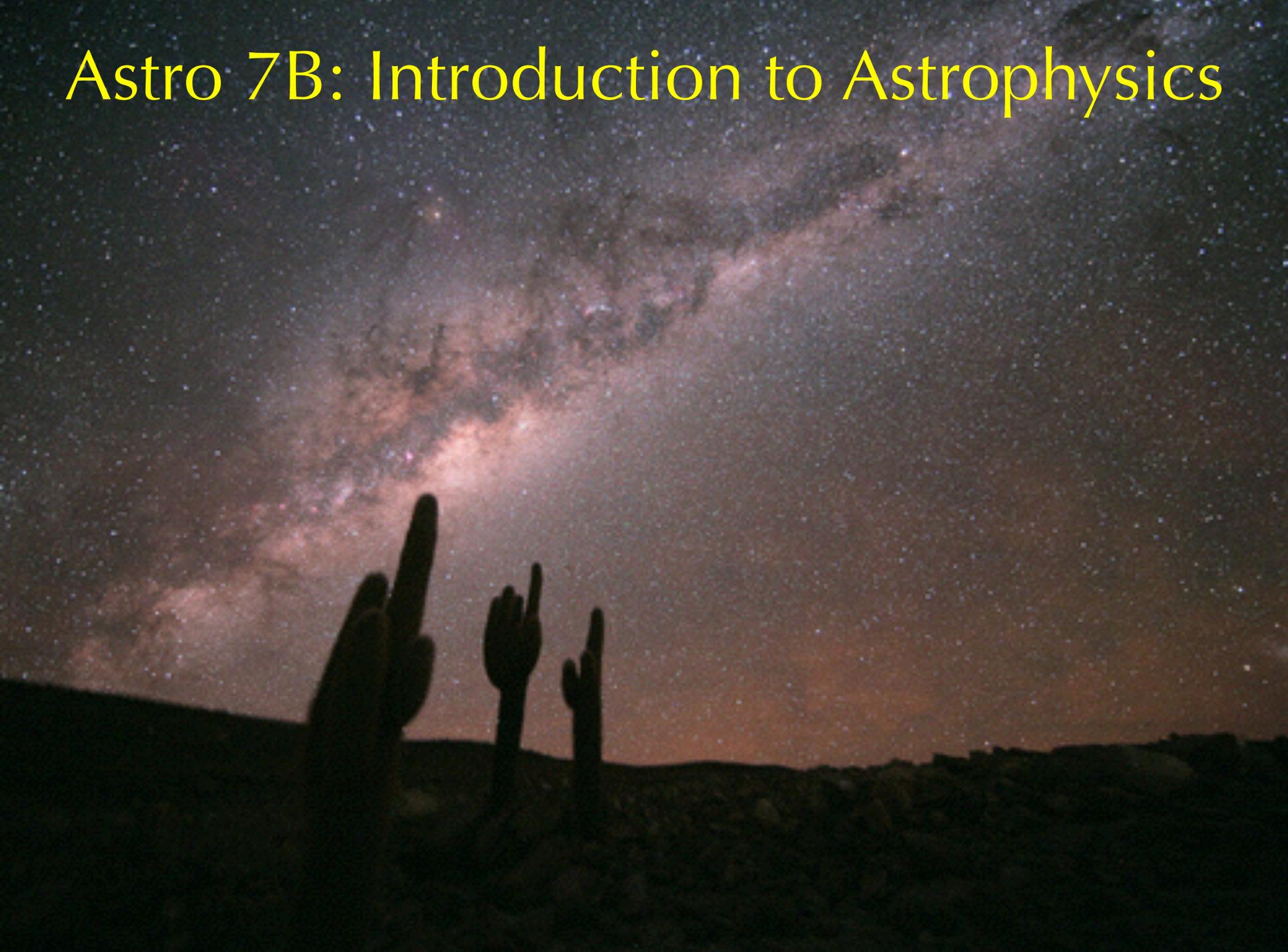


# Astro 7B: Introduction to Astrophysics



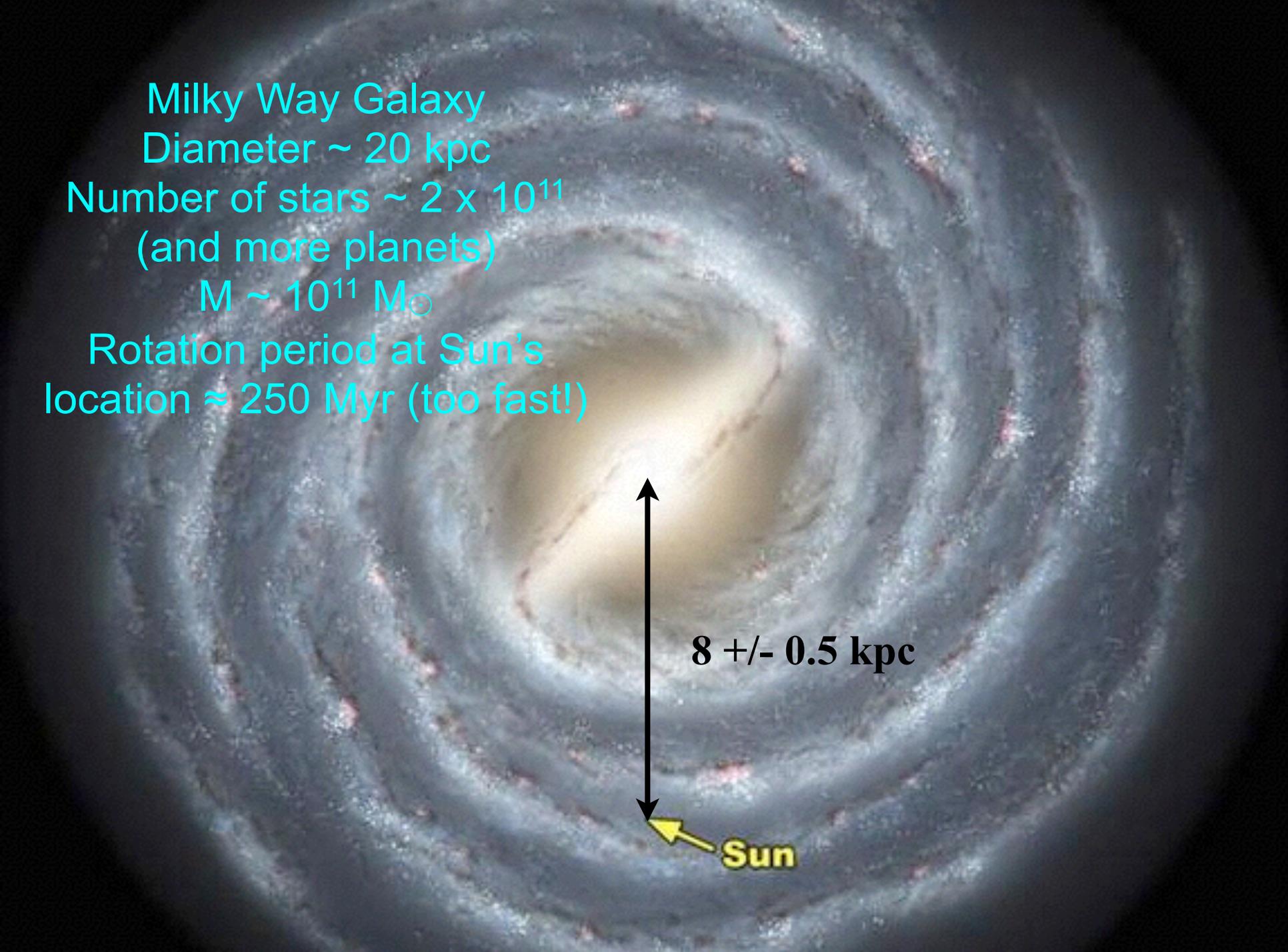
Milky Way Galaxy

Diameter  $\sim 20$  kpc

Number of stars  $\sim 2 \times 10^{11}$   
(and more planets)

$M \sim 10^{11} M_{\odot}$

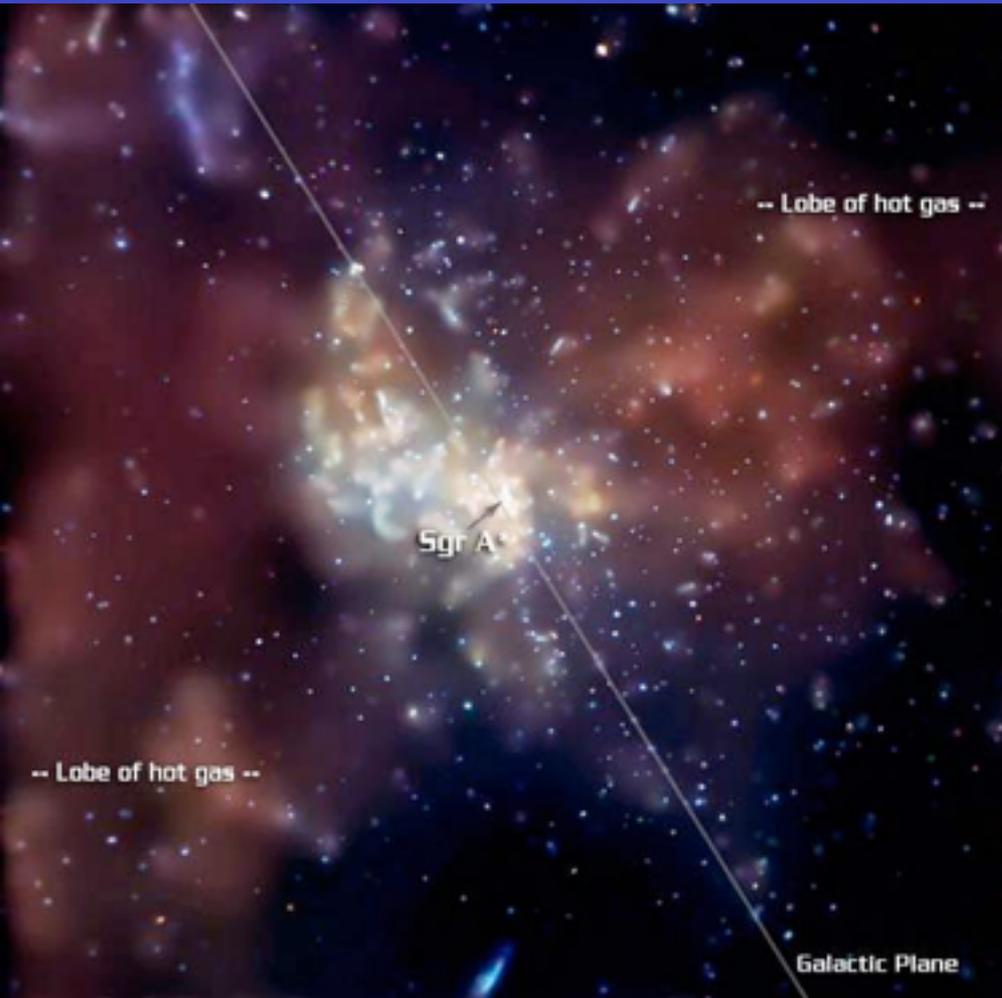
Rotation period at Sun's  
location  $\approx 250$  Myr (too fast!)



**8 +/- 0.5 kpc**

**Sun**

# Supermassive Black Hole and Young Stars at Galactic Center



Central  $\sim 10$  pc

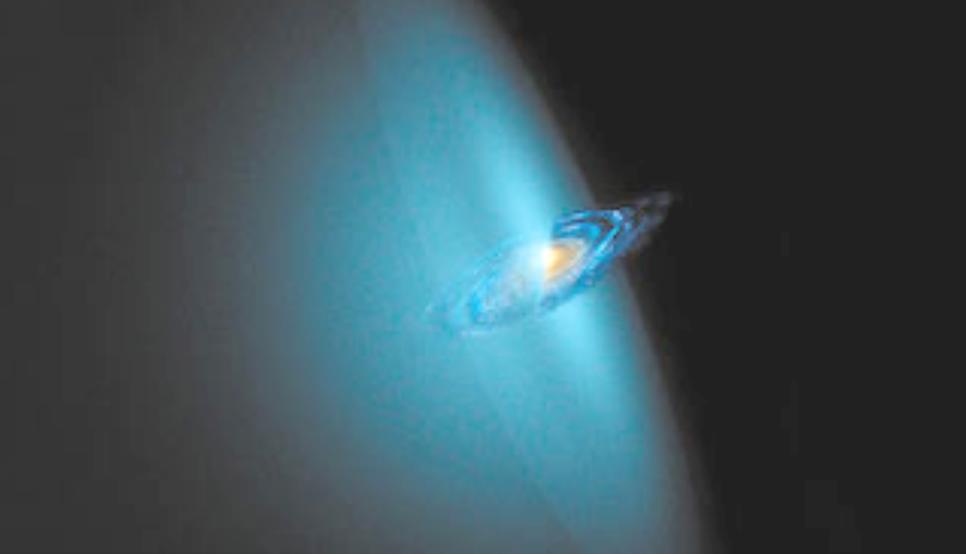


Central  $\sim 1$  pc

# M87: Elliptical galaxy with Active Galactic Nucleus



All galaxies have black hole nuclei

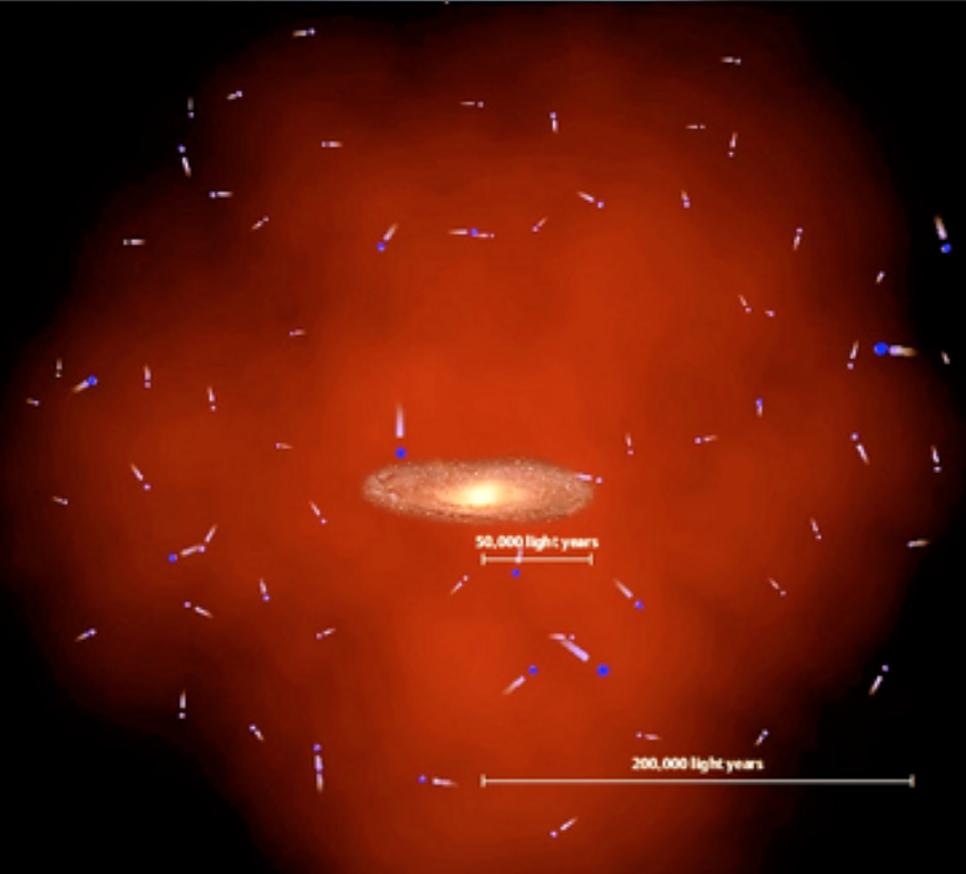


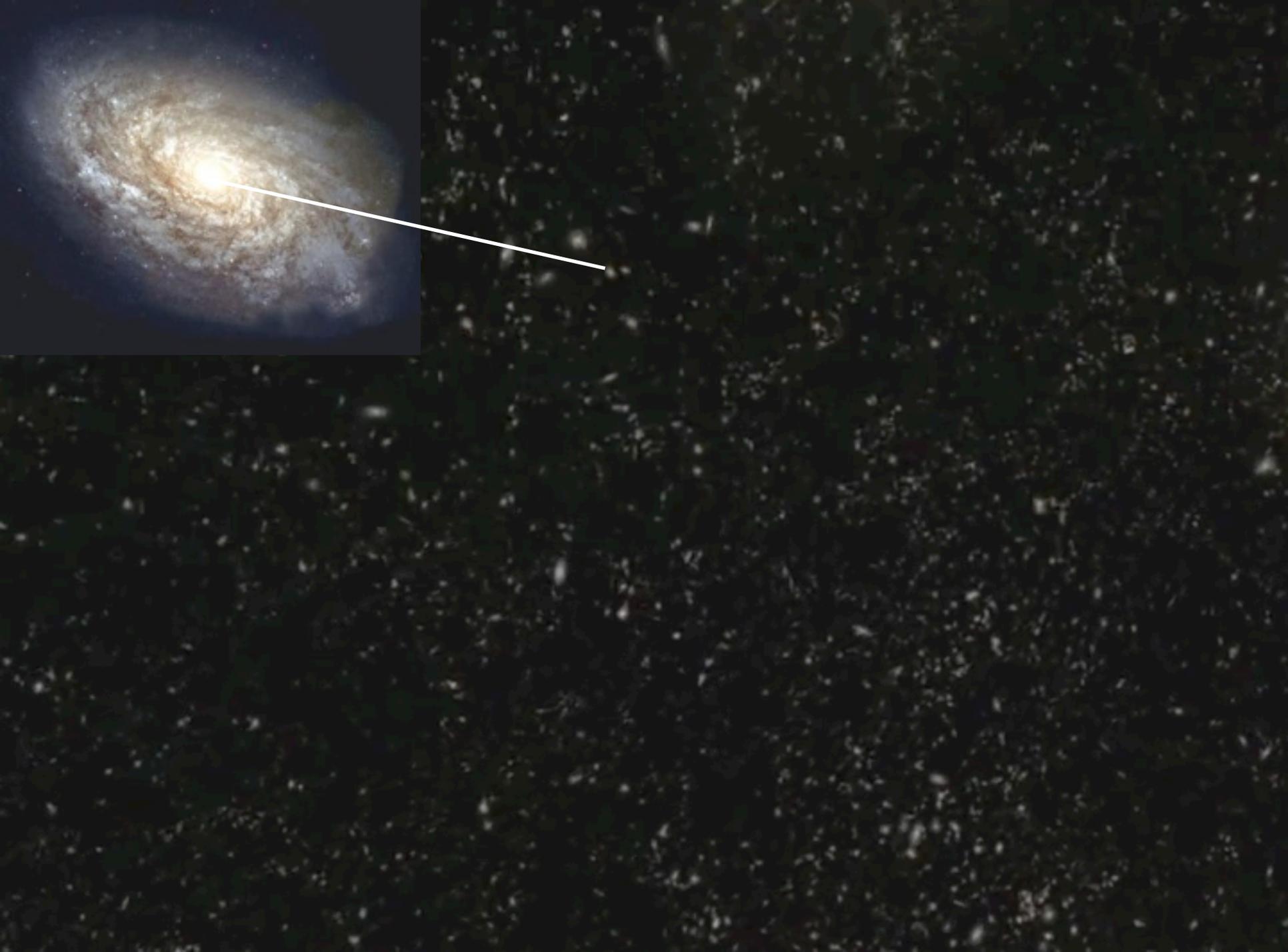
# Dark Matter Halos

$$M \sim 20 \times M_{\text{stars}}$$

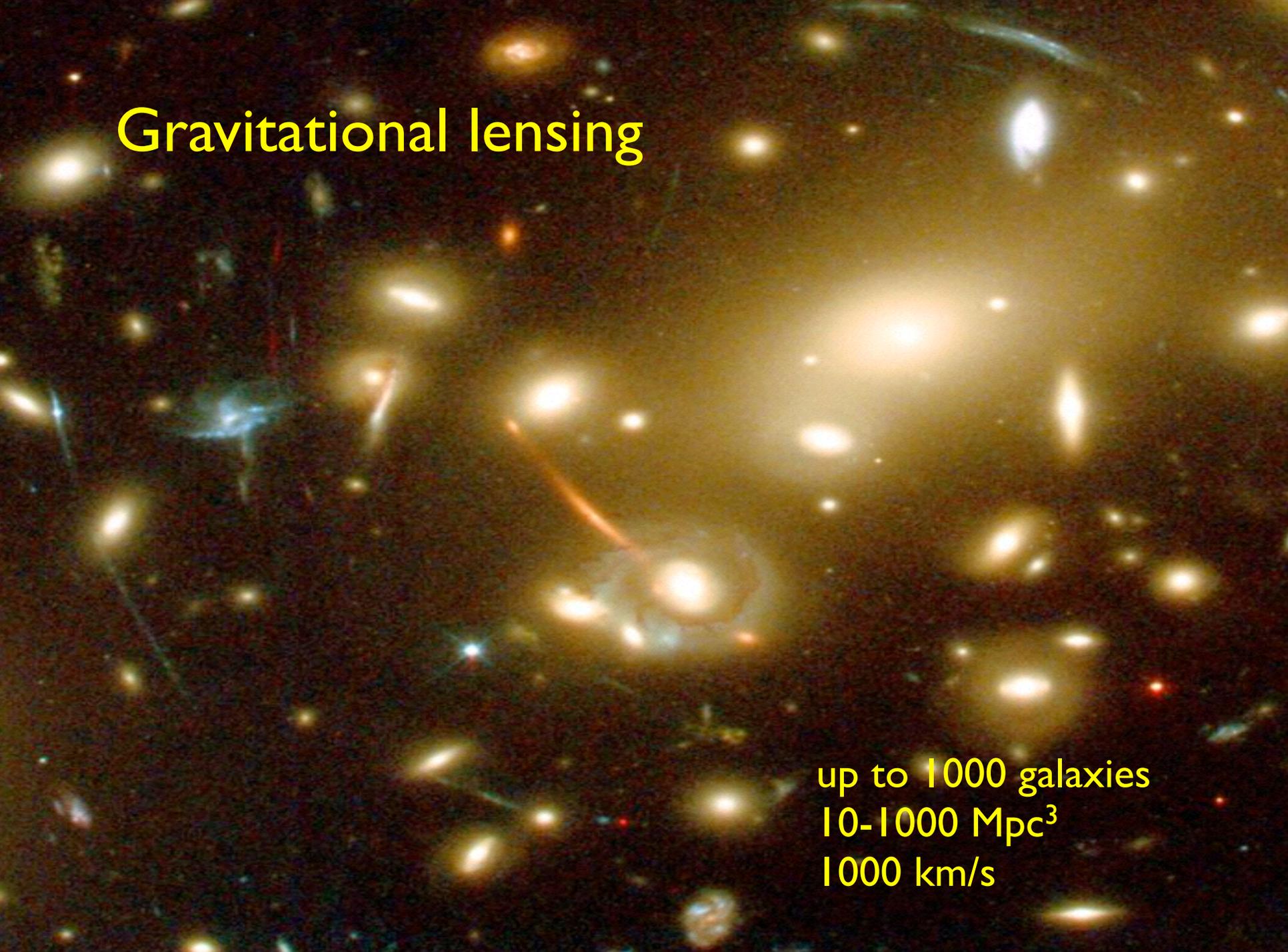
$$R \sim \text{Mpc}$$

~80% of matter consists neither of protons nor neutrons, but some “cold” (non-relativistic) particle, to be identified



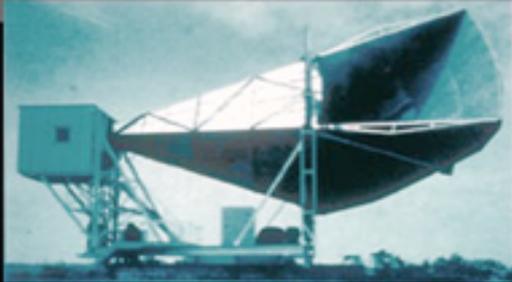


# Gravitational lensing

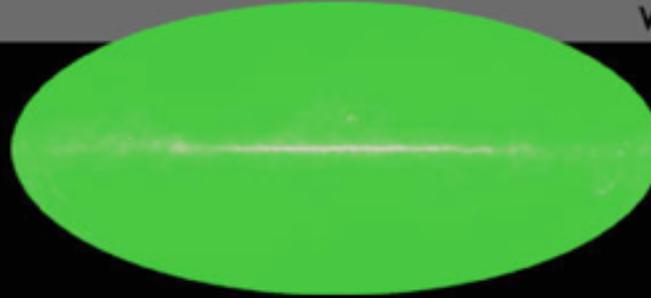


up to 1000 galaxies  
10-1000 Mpc<sup>3</sup>  
1000 km/s

1965



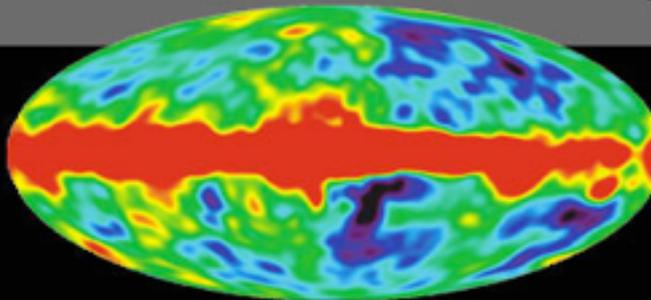
Penzias and  
Wilson



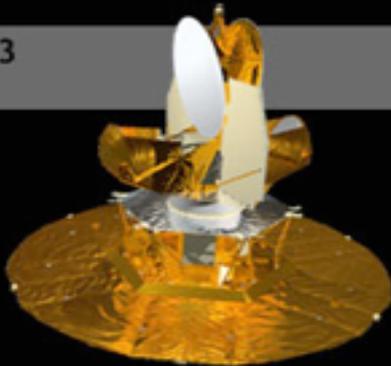
1992



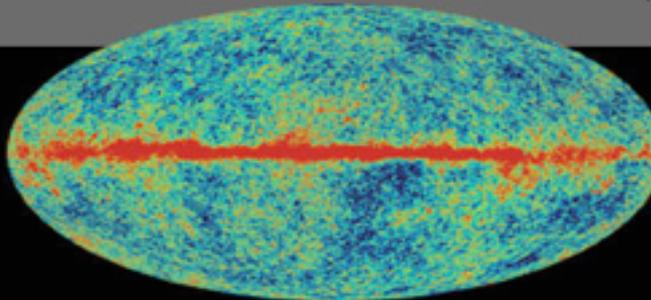
COBE



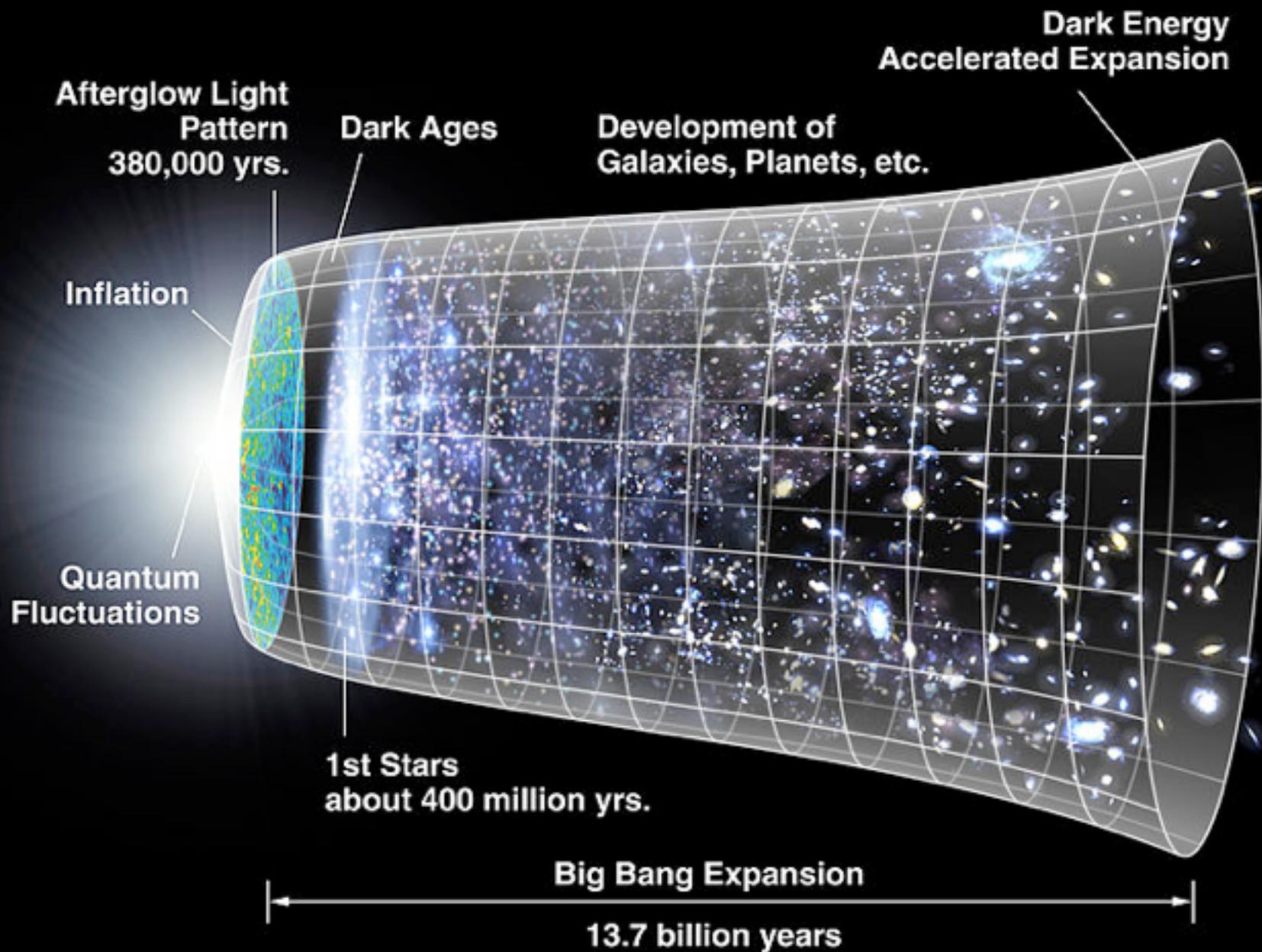
2003



WMAP



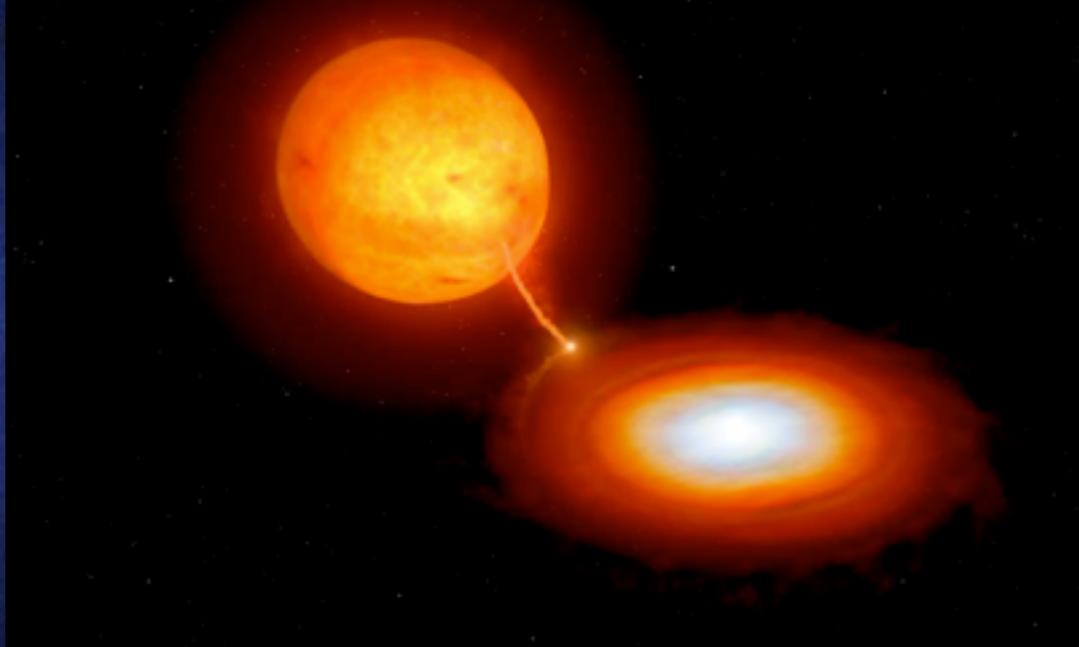
Universe is flat in space and has a finite age (14 Gyr)



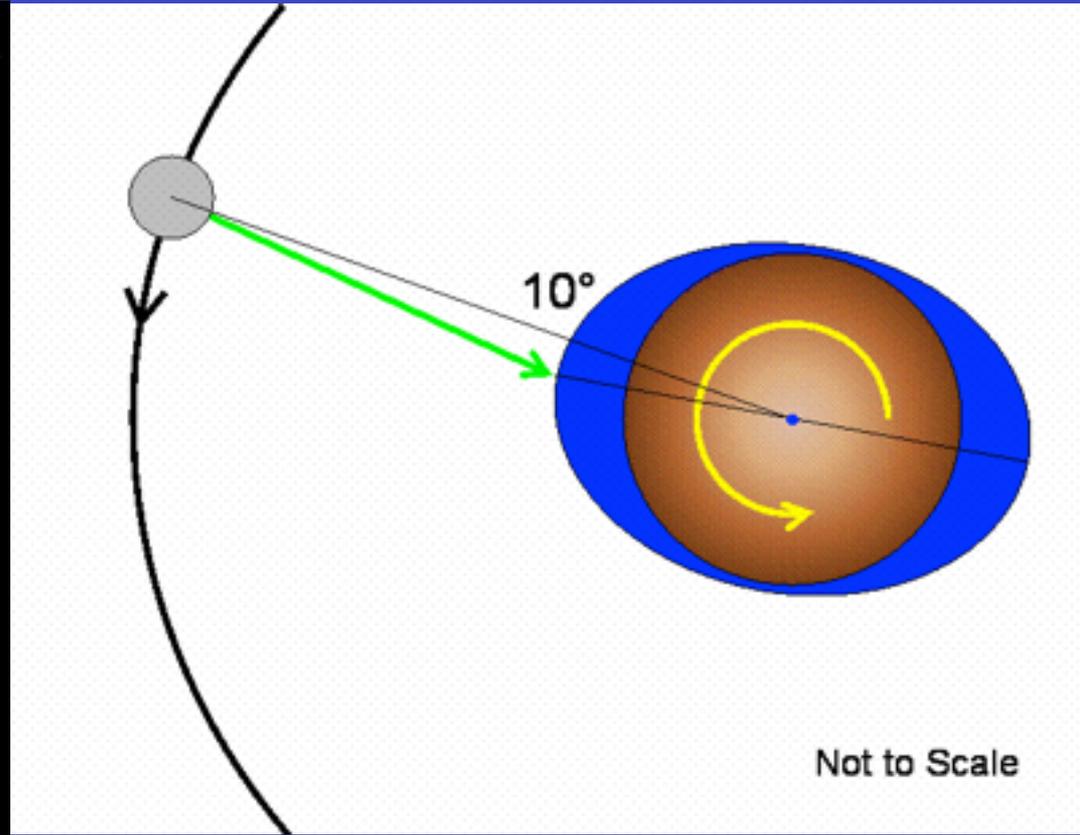
# Kepler 2-body problem



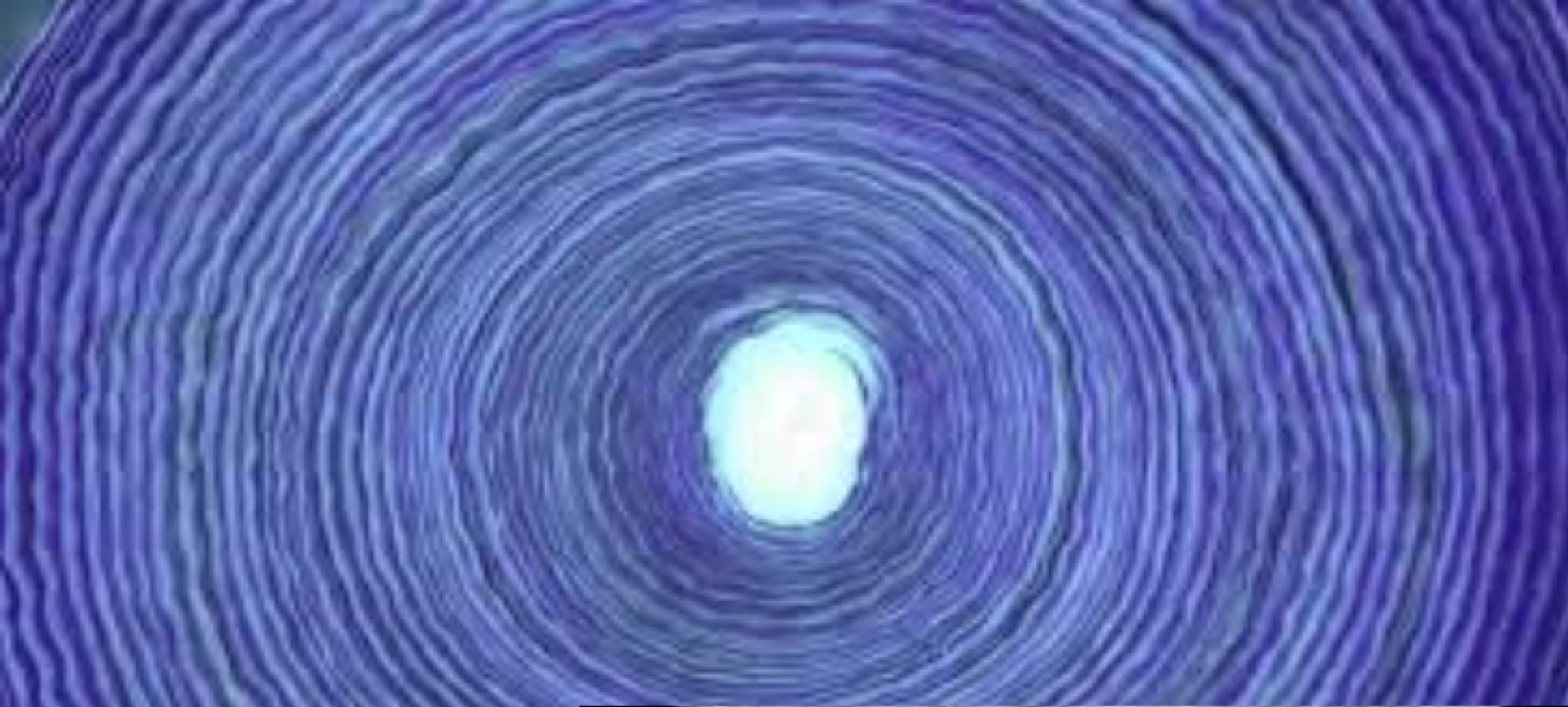
# Binary Stars: Some Interacting



# Shifted tidal bulge → Tidal torque



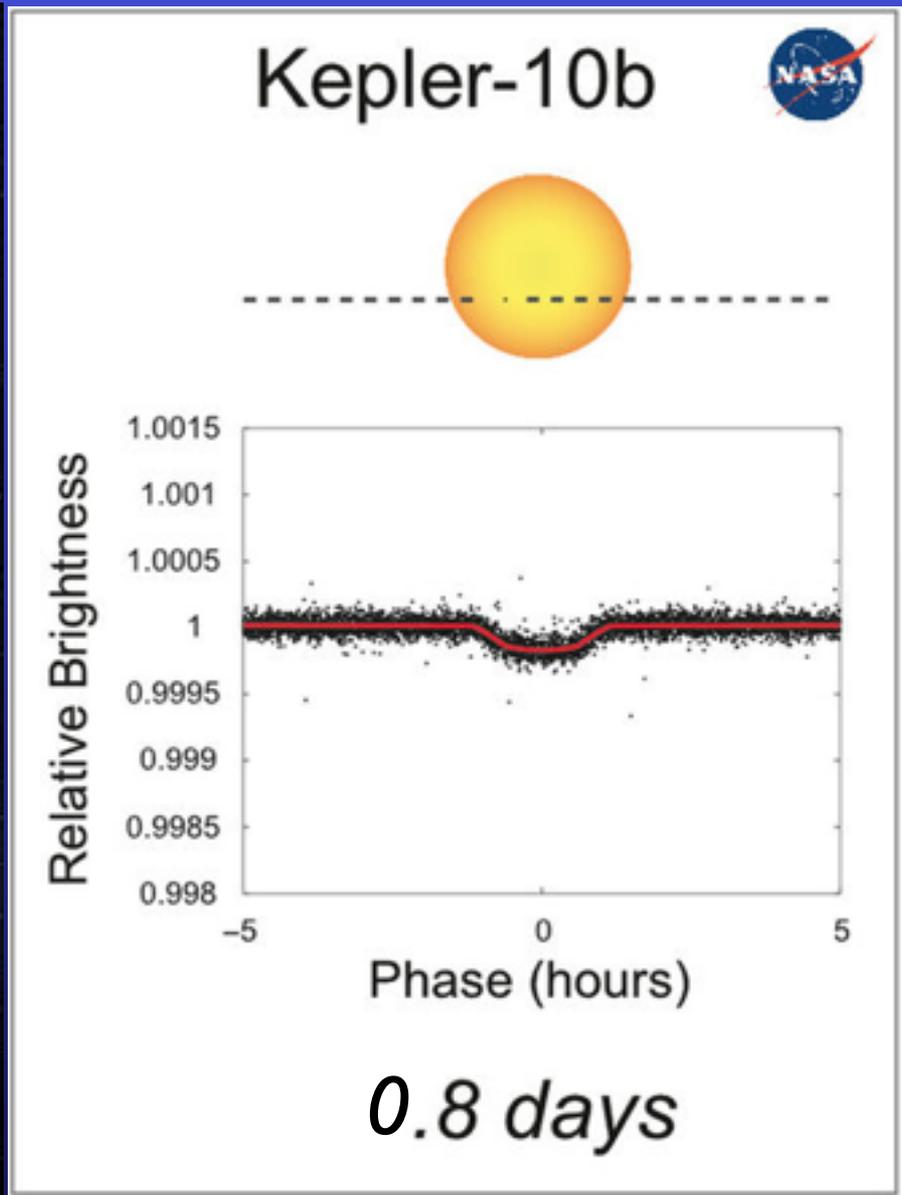
Earth is not yet synchronous with orbit  
→ tidal torque and continued tidal spin-orbit evolution  
Moon has already synchronized



# Galilean moons of Jupiter



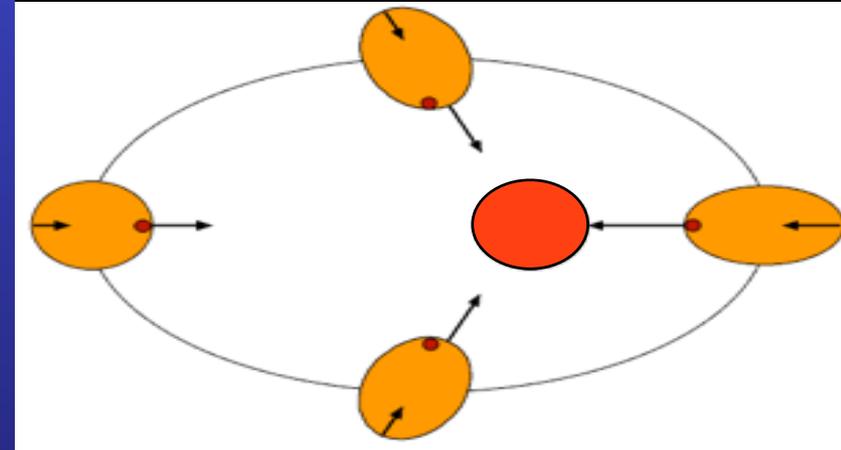
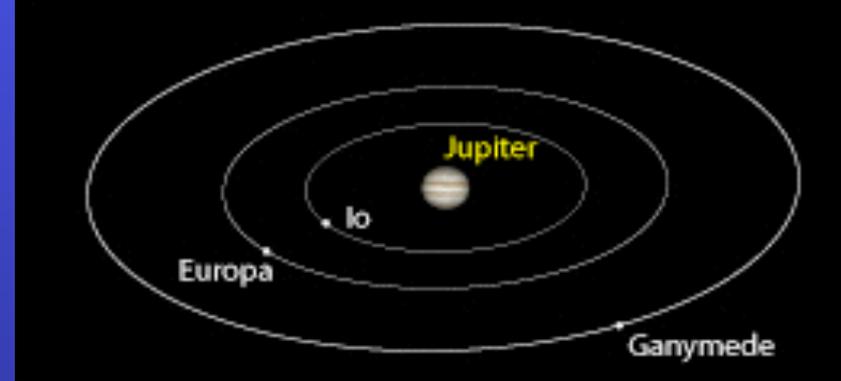
Moons have synchronized.  
Jupiter has not.



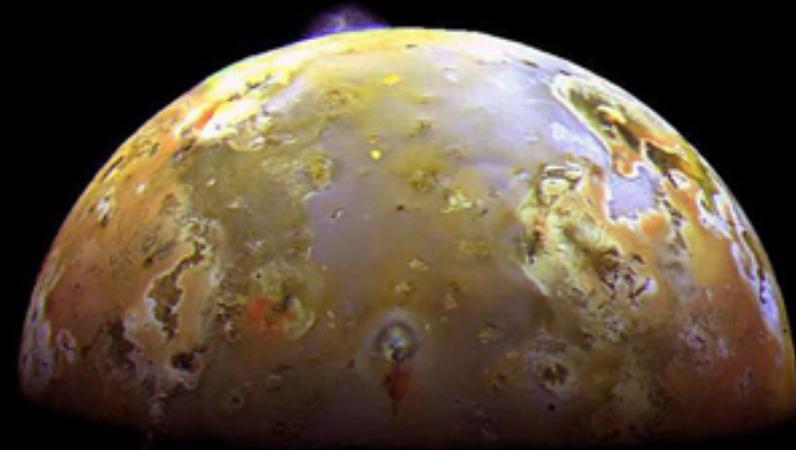
A (Probably) Tidally Locked Extrasolar Planet

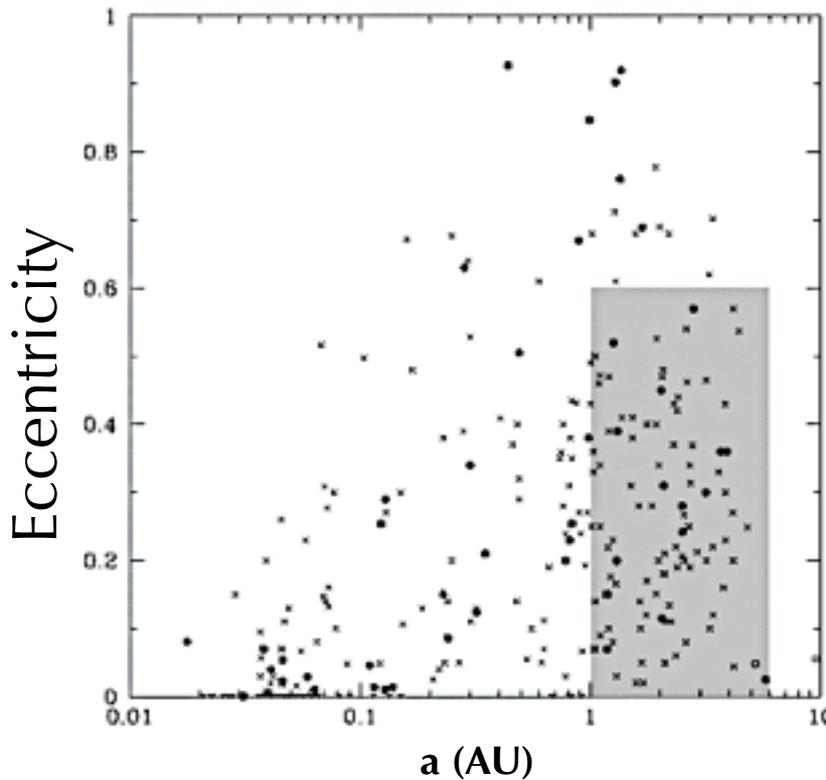
# Tidal melting of Io

1. Jupiter's asynchronous tide repels Io
2. Io approaches Europa
3. Europa forces Io onto eccentric orbit
4. Io's eccentricity tide stresses Io

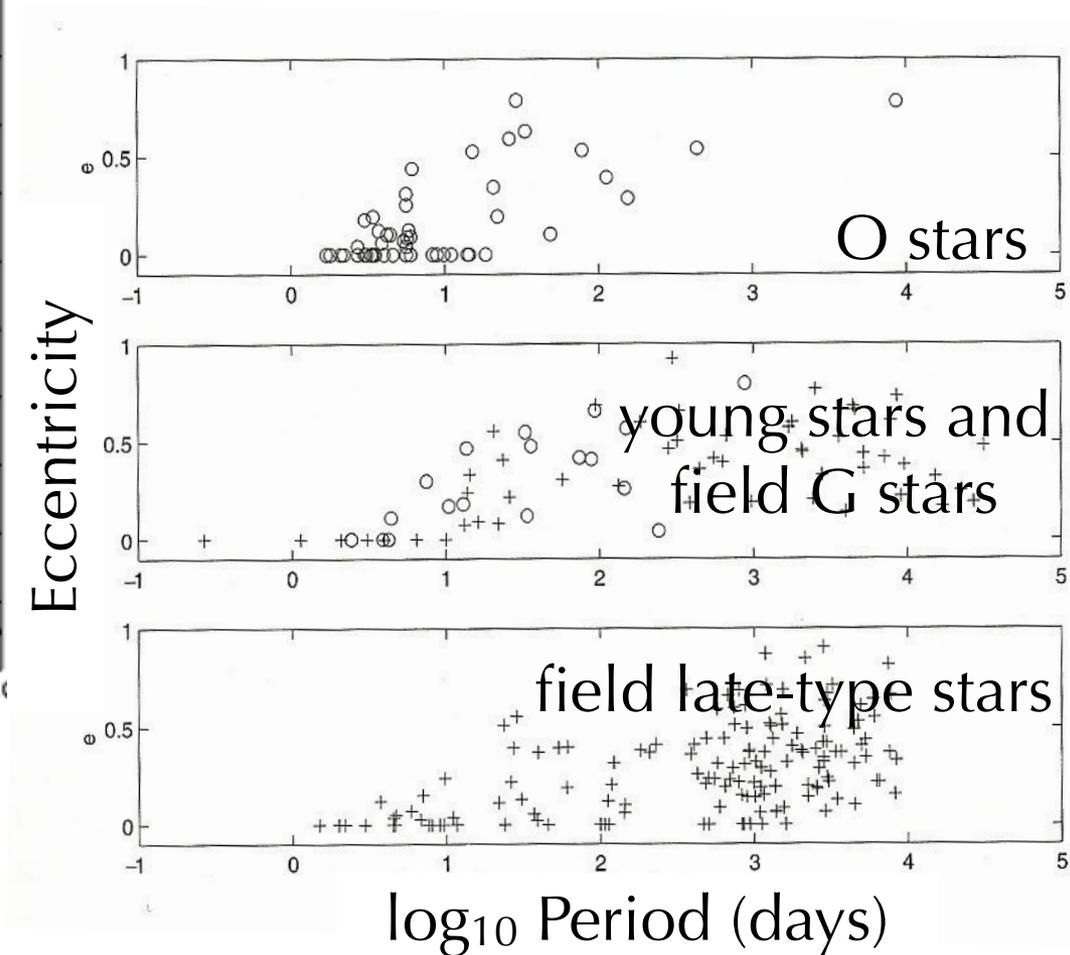


BBC TWO



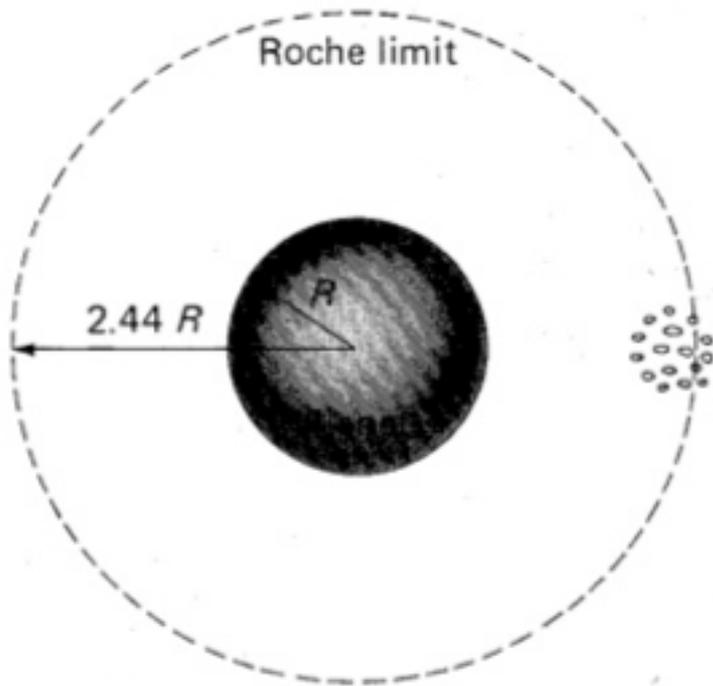


Extrasolar planets detected  
 by radial velocity  
 (> Neptune mass)



Tidal interactions (circularization +  
 synchronization) most important close-in

■ **Figure 18.5** The Roche limit. A satellite (far right) in orbit around a planet is slightly deformed by tides. If it moves closer to the planet, it is deformed even more and finally disintegrates when it passes within the Roche limit (*dotted line*), located 2.44 times the planet's radius from its center.



tidal disruption



Satellite



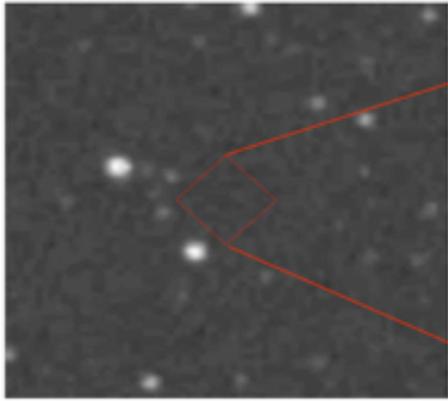
planet and star formation

# Comet Shoemaker-Levy 9: Tidally disrupted by Jupiter

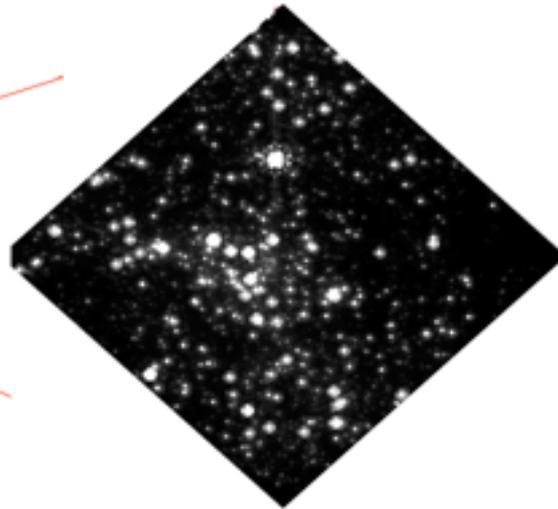


Planetary rings are *inside* the Roche disruption zone.

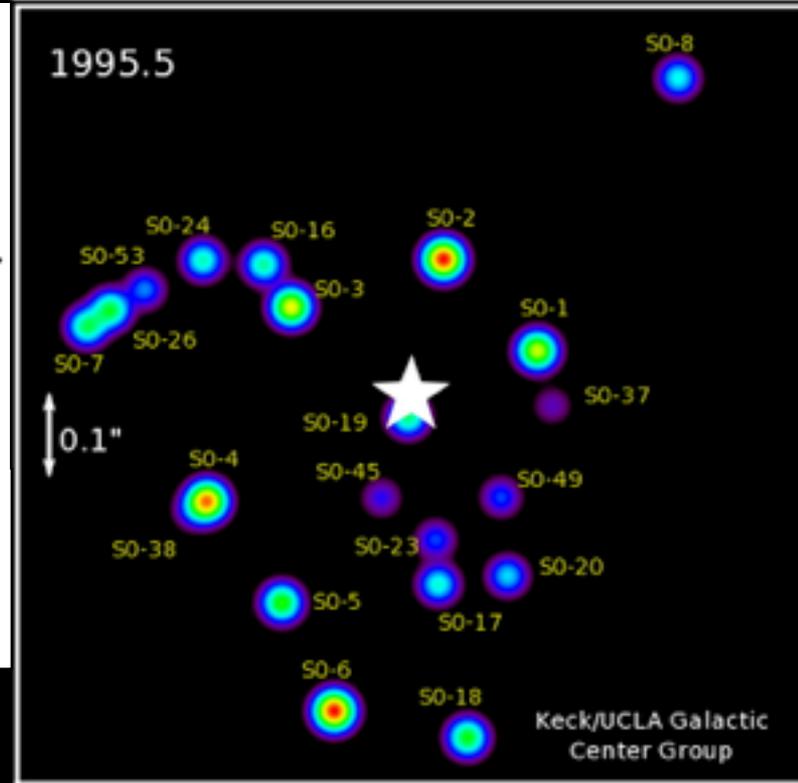
Satellites are always *outside*.



Visible Light



Infrared Light



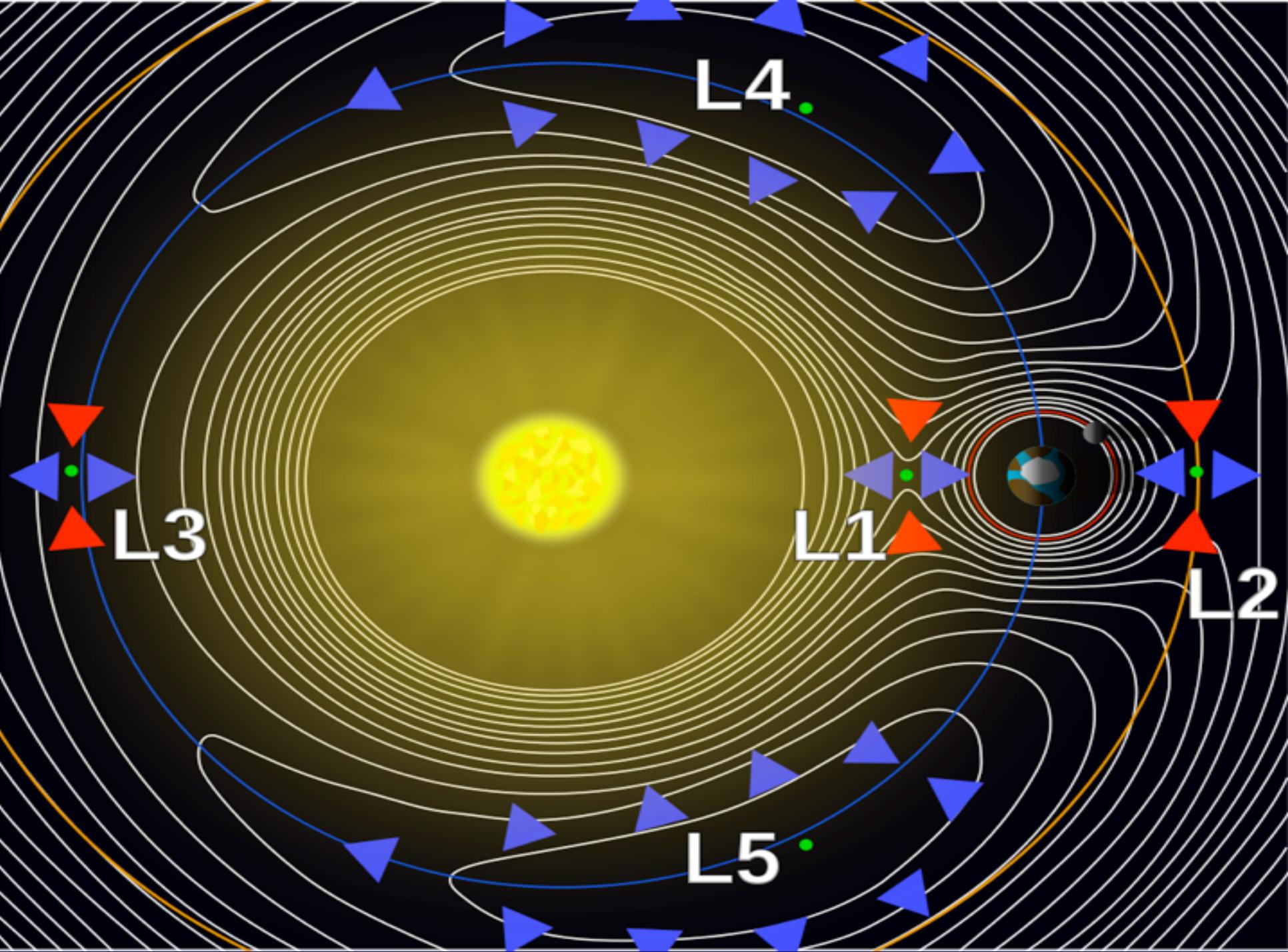
## The Supermassive Black Hole at the Galactic Center

$$M \approx 3.5 \times 10^6 M_{\odot}$$

$$R \approx 0.07 \text{ AU}$$

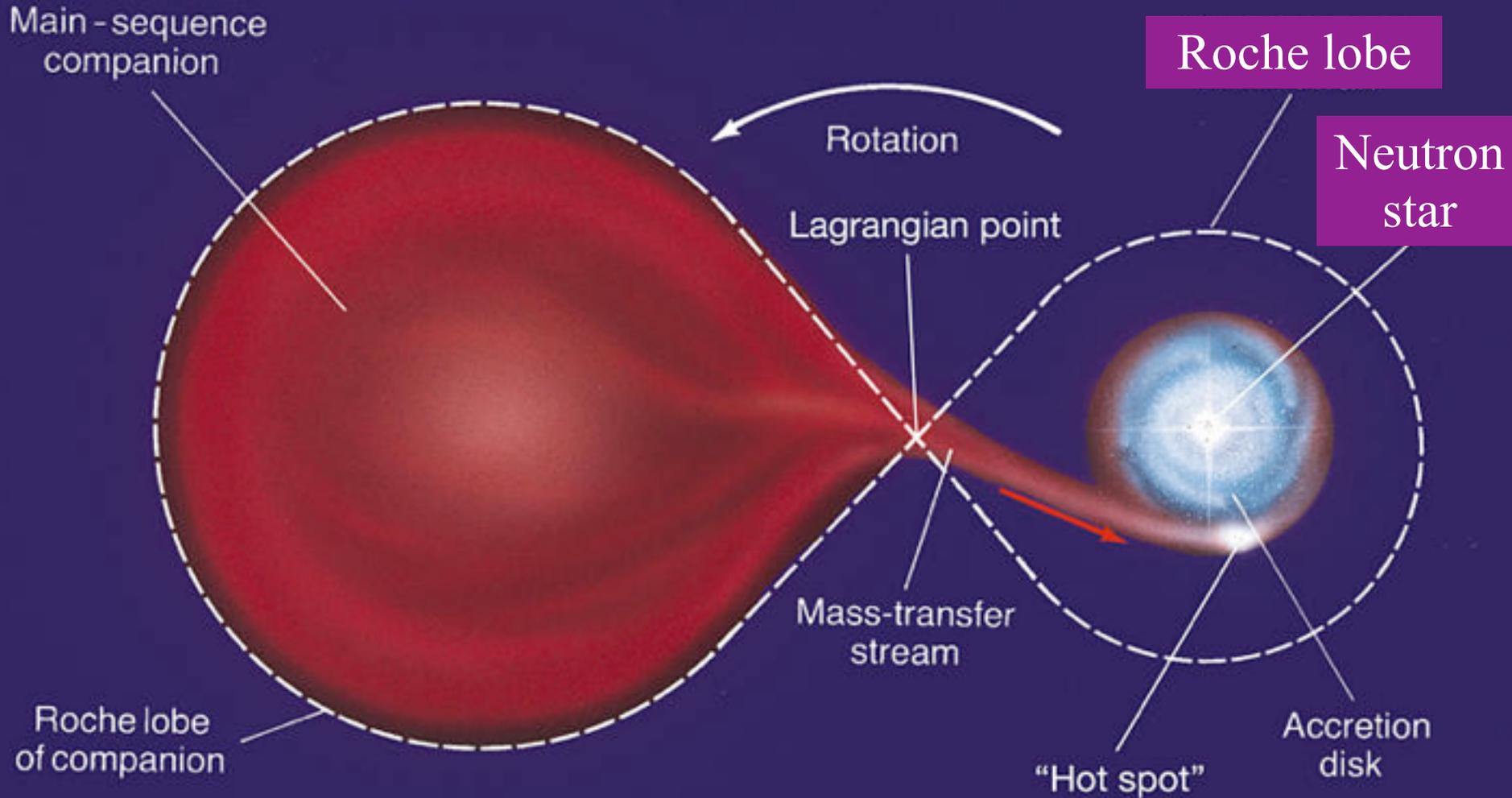


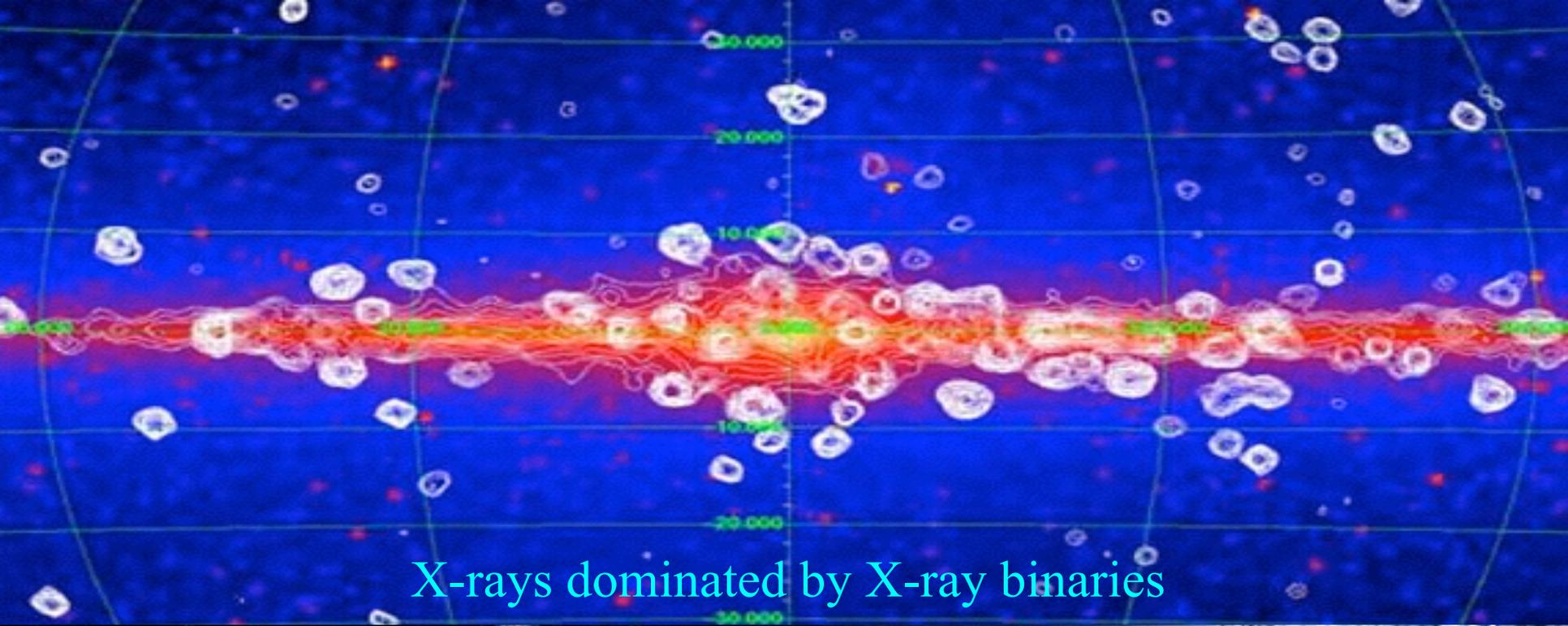






# LMXB = "Low-Mass X-ray Binary"



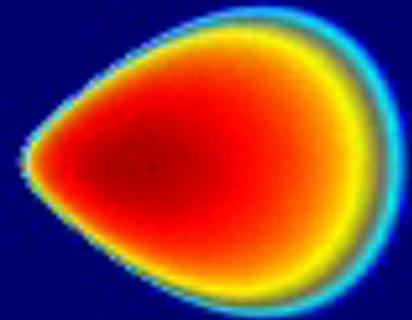
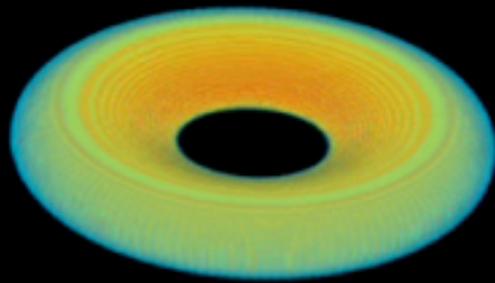
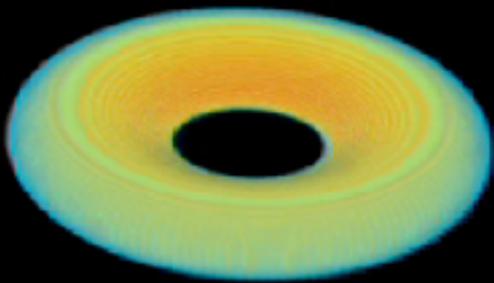


X-rays dominated by X-ray binaries



# Accretion Disk

Turbulence driven by Magneto-Rotational Instability



Spin

Parameter

$$a^* = cJ / GM^2$$

( $0 < a^* < 1$ )

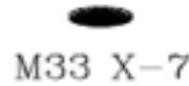
$a^* \sim 0.25$



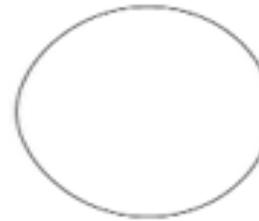
LMC X-3



LMC X-1



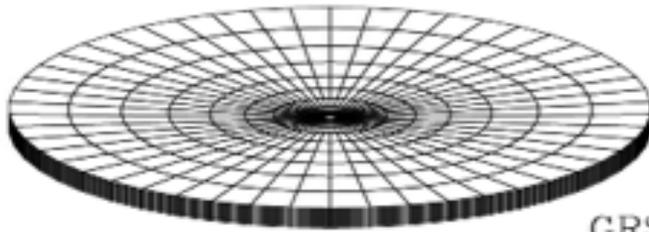
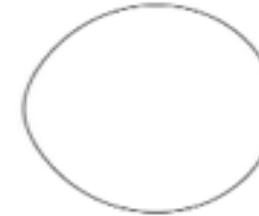
M33 X-7



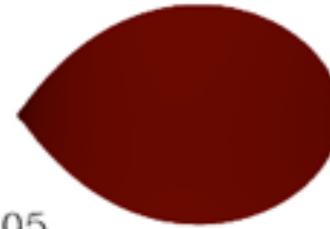
$a^* = 0.77 \pm 0.05$



Cyg X-1



GRS 1915+105



$a^* = 0.98 - 1.0$

XTE J1650-500

XTE J1118+480

XTE J1859+226

GRS 1009-45

GRS 1124-683

SAX J1819.3-2525

GS 2000+25

H1705-250

GRO J1655-40

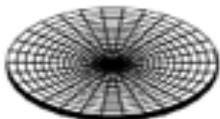
A0620-00

GRO J0422+32

XTE J1550-564

$a^* = 0.65 - 0.75$

$(a^* \sim 0.5)$



GS 2023+338



GS 1354-64



GX 339-4



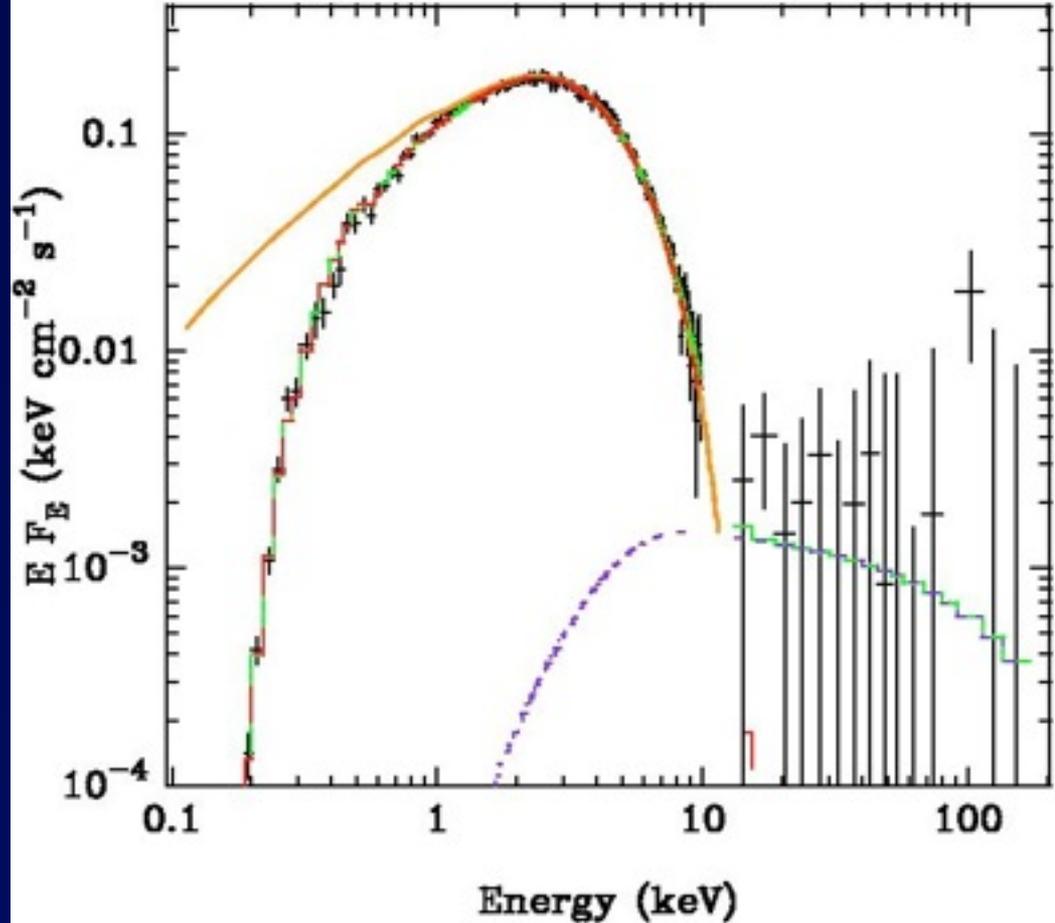
4U 1543-47

$a^* = 0.75 - 0.85$

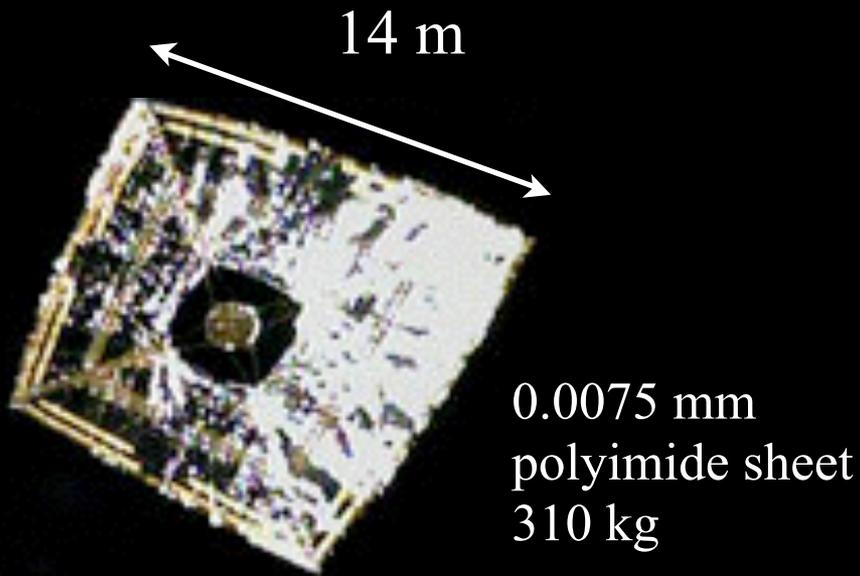
# Blackbody-Like Spectral State in BH Accretion Disk

LMC X-3: Beppo-SAX  
(Davis, Done & Blaes 2006)

Up to 10 keV, the only component seen is the disk.  
Beyond that, a weak PL tail



- Perfect for estimating inner radius of accretion disk → BH spin
- Just need to estimate  $L_x$ ,  $T_x$  (and  $N_H$ ) from X-ray continuum
- Use full relativistic model (Novikov-Thorne 1973; KerrBB, Li et al. 2005)

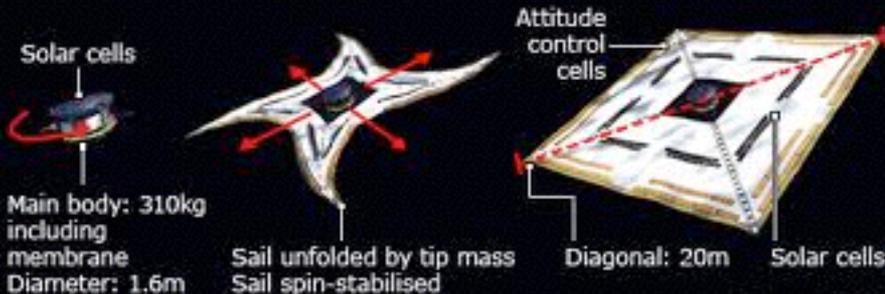


© JAXA

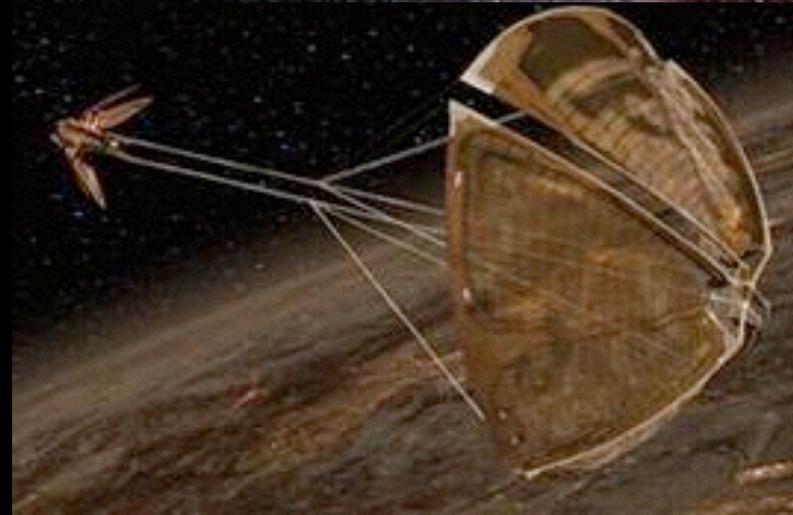
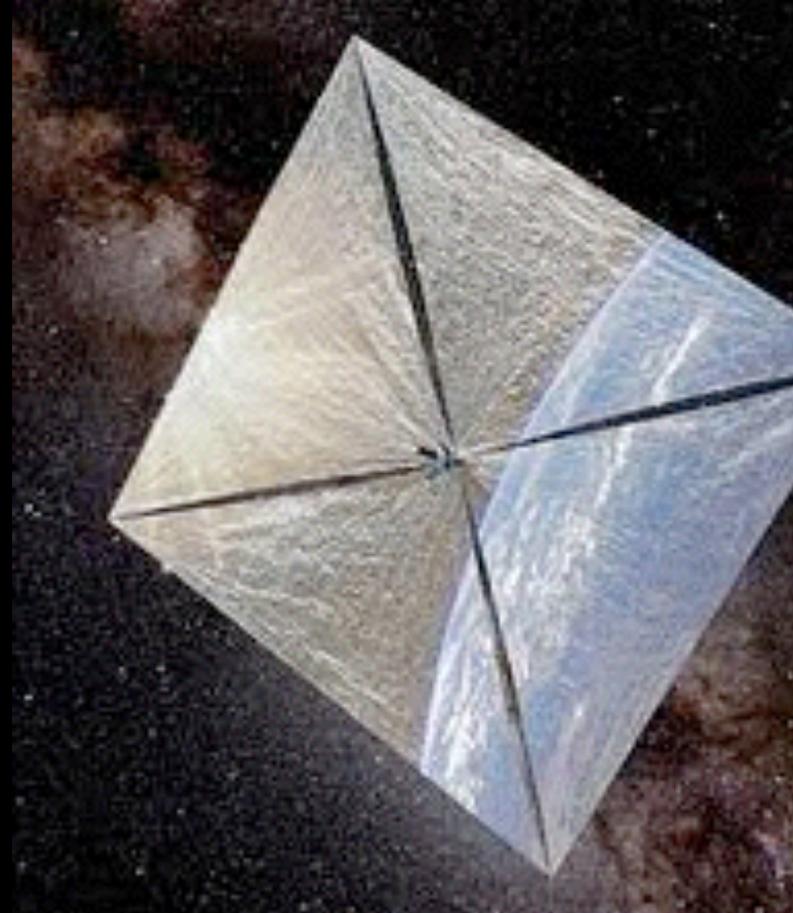
## IKAROS

“Interplanetary Kite-craft  
Accelerated by Radiation of the Sun”

- 1 Launch vehicle    2 Solar sail deployment    3 Fully deployed



Source: JAXA





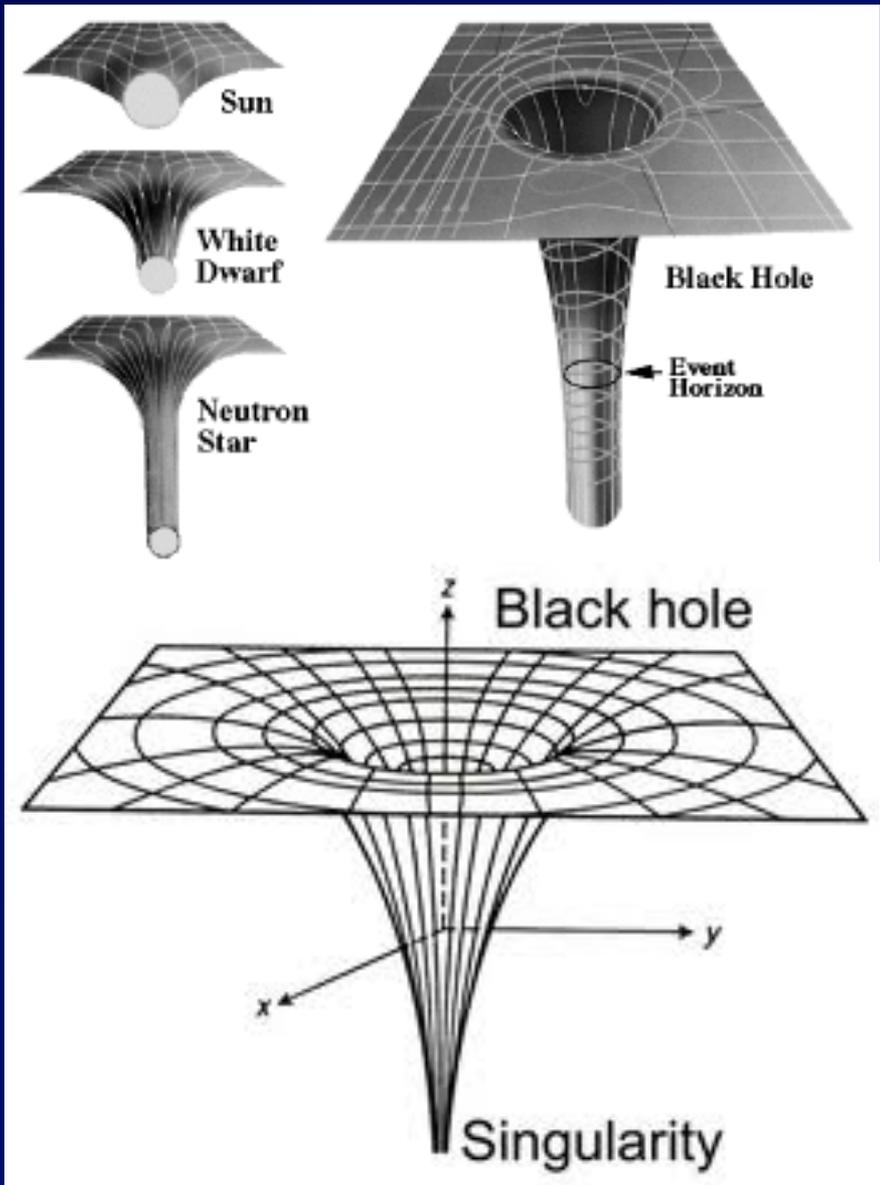
“Superstar” Eta Carina:  $100\text{-}150 M_{\odot}$   
lost  $\sim 30 M_{\odot}$  in previous eruptions  
Eruptions driven by radiation pressure

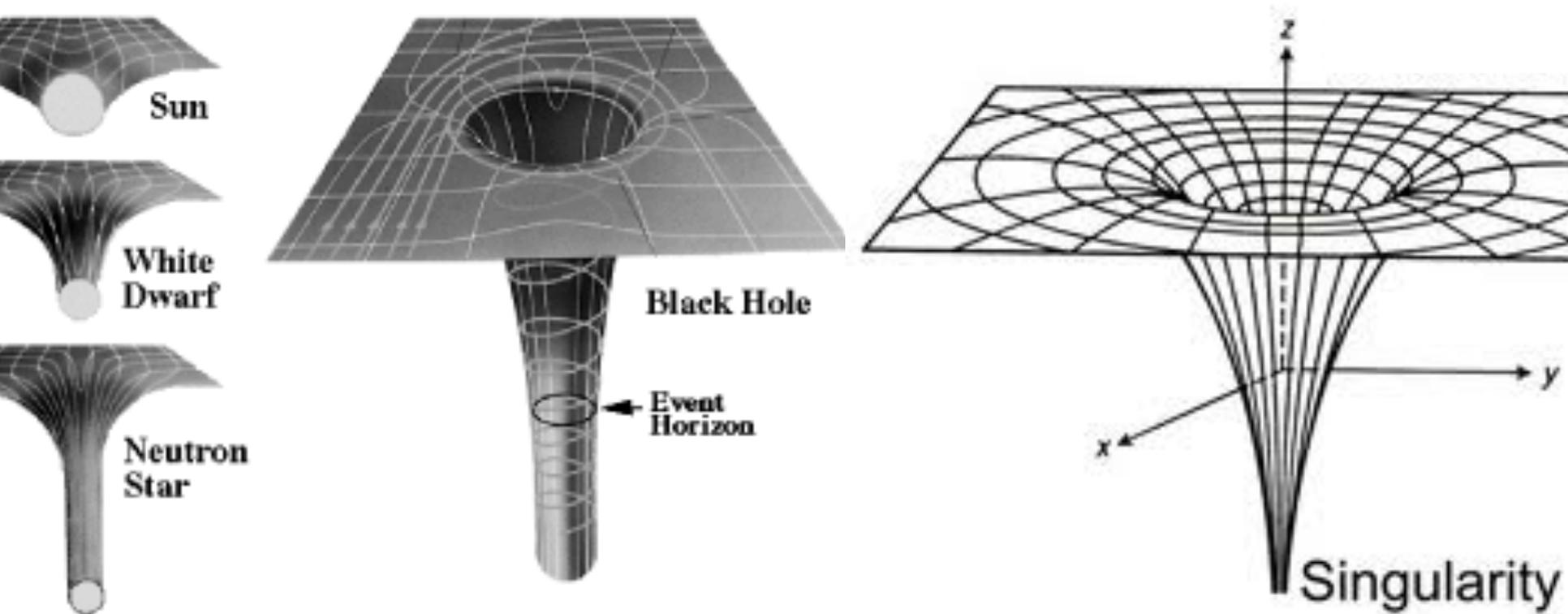
29. If there should really exist in nature any bodies, whose density is not less than that of the sun, and whose diameters are more than 500 times the diameter of the sun, since their light could not arrive at us; or if there should exist any other bodies of a somewhat smaller size, which are not naturally luminous; of the existence of bodies under either of these circumstances, we could have no information from sight; yet, if any other luminous bodies should happen to revolve about them we might still perhaps from the motions of these revolving bodies infer the existence of the central ones with some degree of probability, as this might afford a clue to some of the apparent irregularities of the revolving bodies, which would not be easily explicable on any other hypothesis; but as the consequences of such a supposition are very obvious, and the consideration of them somewhat beside my present purpose, I shall not prosecute them any farther.

John Michell

$$G_{\mu\nu} + g_{\mu\nu}\Lambda = \frac{8\pi G}{c^4}T_{\mu\nu}.$$

Albert Einstein

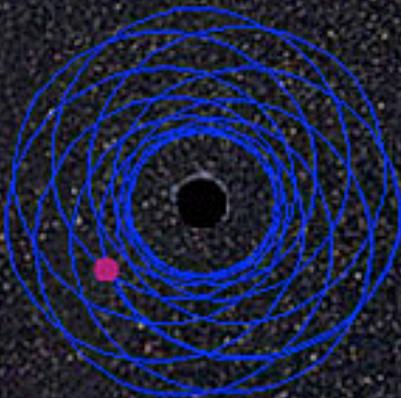




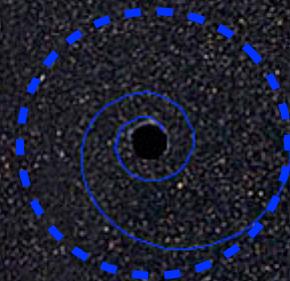
stable and closed



stable and not closed

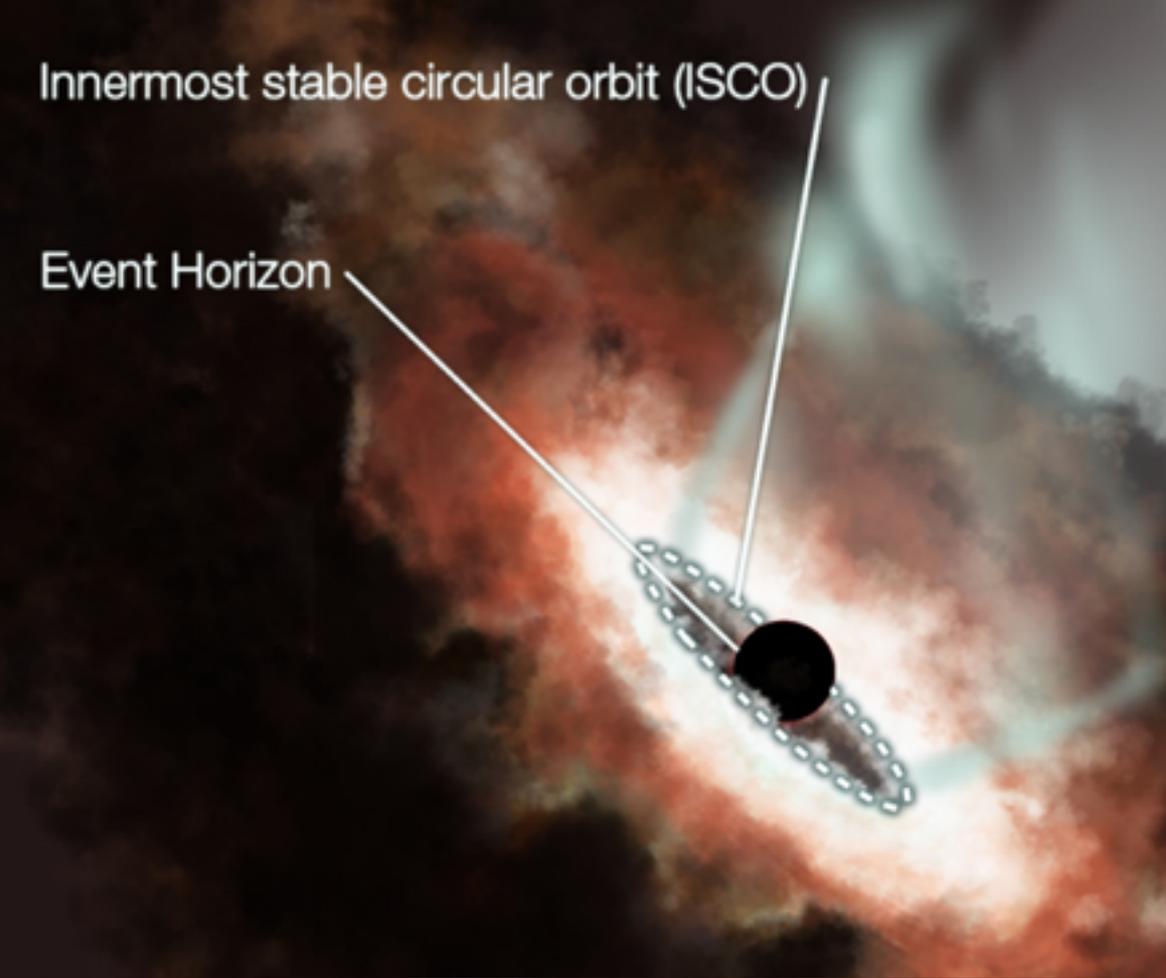


unstable

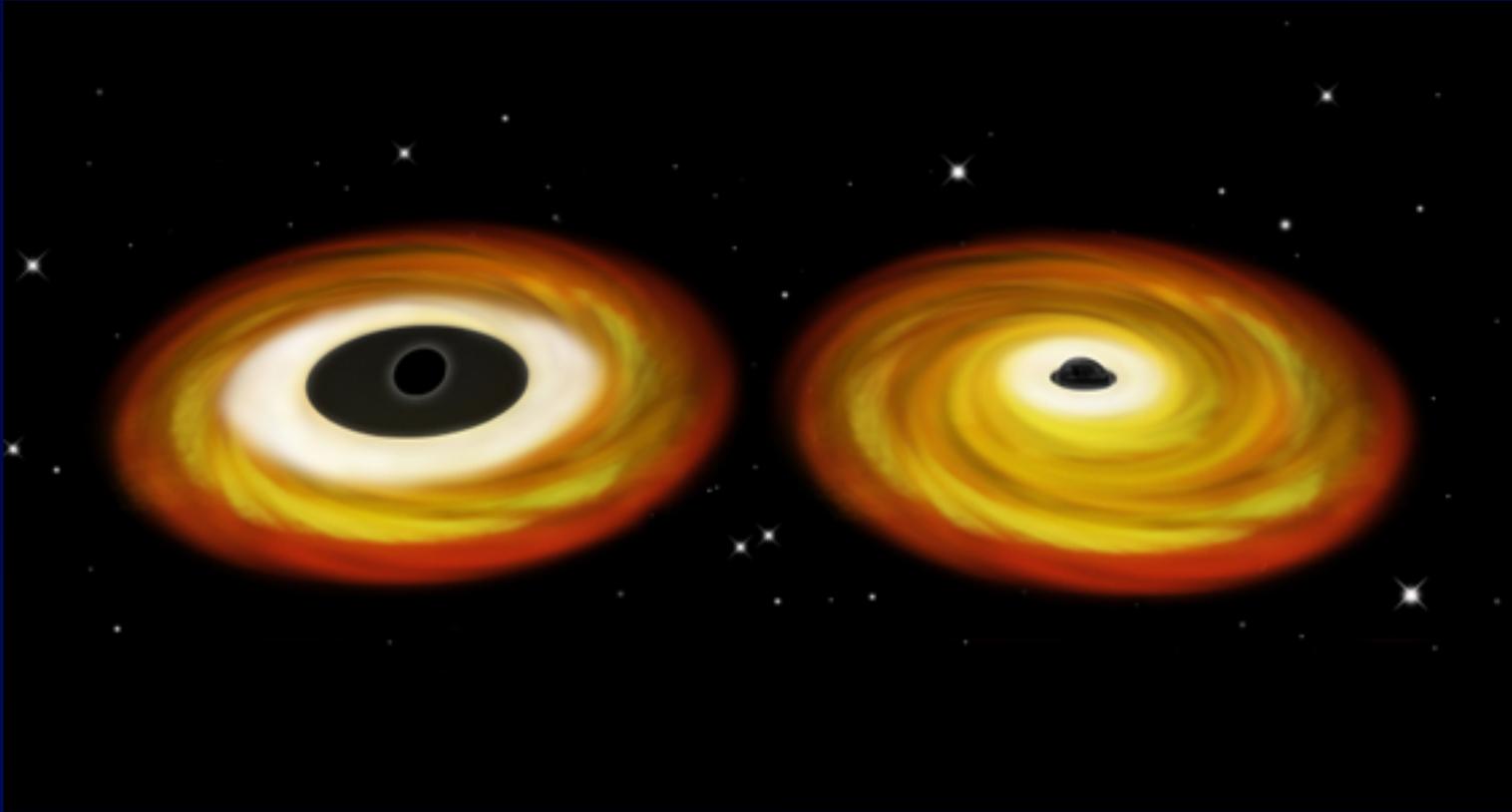


Innermost stable circular orbit (ISCO)

Event Horizon



# Measuring spin for stellar-mass black holes



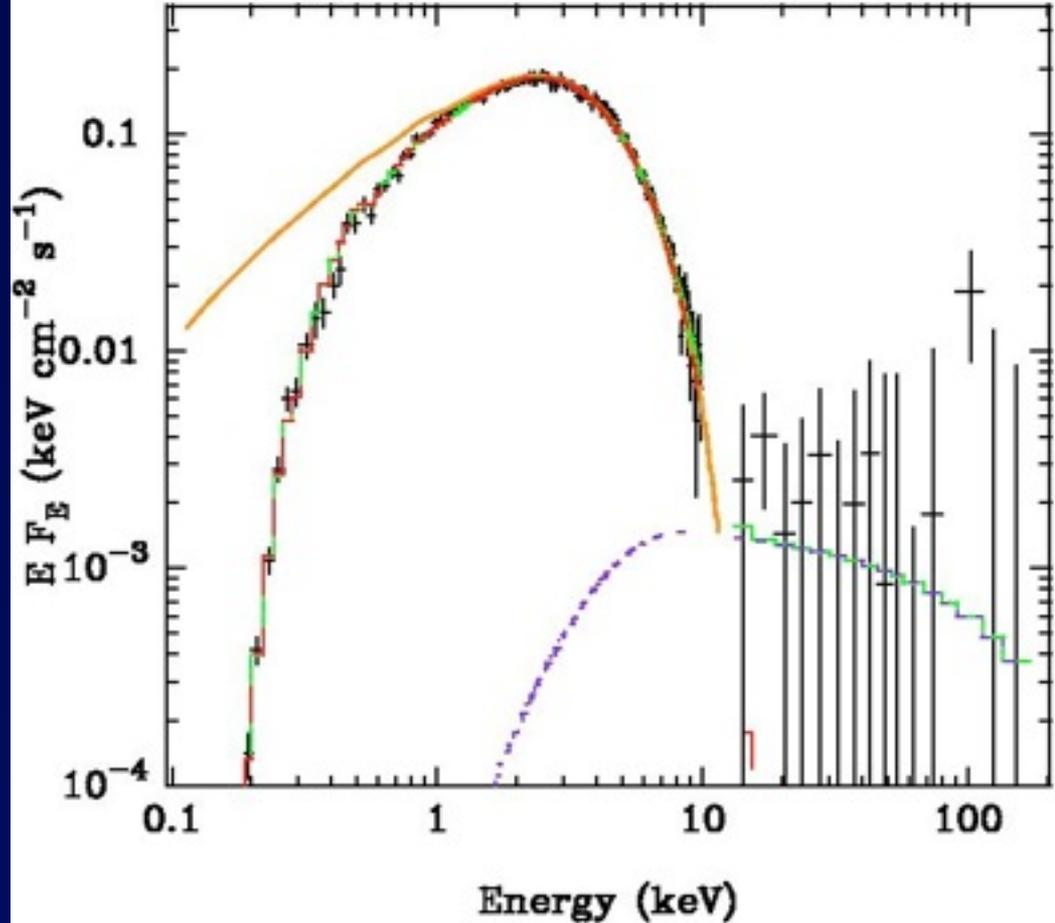
Measure **Radius of the Hole** in the disk by estimating the area of the bright inner disk using **X-ray Data** in the **Thermal State**:  $L_x$  and  $T_x$

Zhang et al. (1997); Shafee et al. (2006); Davis et al. (2006);  
McClintock et al. (2006); Middleton et al. (2006); Liu et al. (2008)

# Blackbody-Like Spectral State in BH Accretion Disk

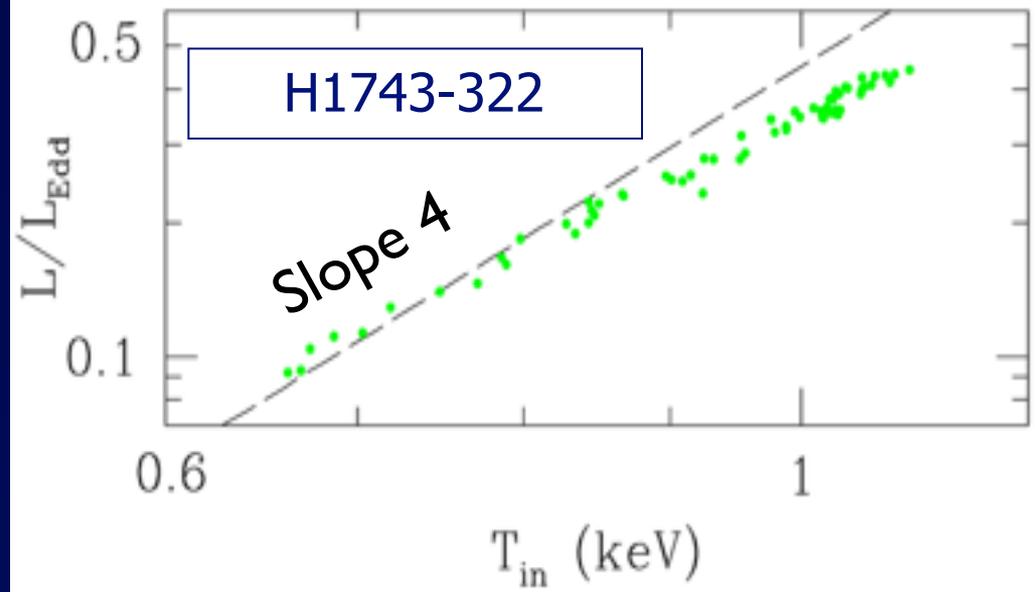
LMC X-3: Beppo-SAX  
(Davis, Done & Blaes 2006)

Up to 10 keV, the only component seen is the disk.  
Beyond that, a weak PL tail



- Perfect for estimating inner radius of accretion disk → BH spin
- Just need to estimate  $L_x$ ,  $T_x$  (and  $N_H$ ) from X-ray continuum
- Use full relativistic model (Novikov-Thorne 1973; KERRBB, Li et al. 2005)

# A Test of the Blackbody Assumption



- For a blackbody,  $L$  scales as  $T^4$
- BH accretion disks vary a lot in their luminosity
- If a disk is a good blackbody,  $L$  should vary as  $T^4$
- **Looks reasonable**

Kubota et al. (2002)

McClintock et al. (2008)

$$L = A\sigma T^4$$

Spin

Parameter

$$a^* = cJ / GM^2$$

( $0 < a^* < 1$ )

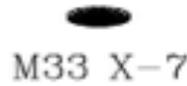
$a^* \sim 0.25$



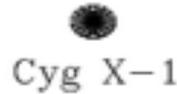
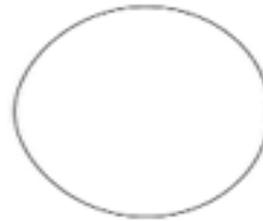
LMC X-3



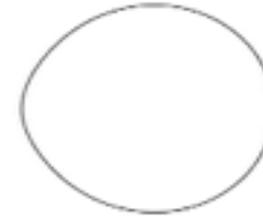
LMC X-1



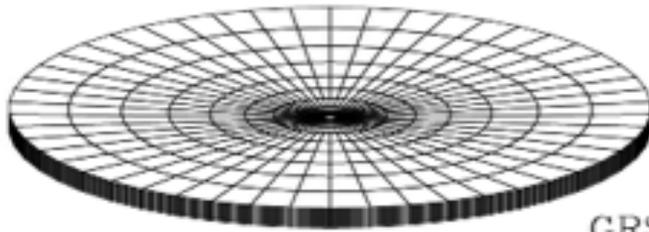
M33 X-7



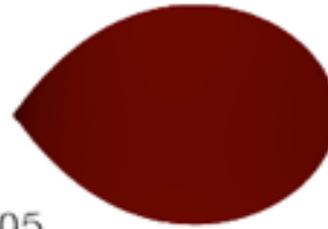
Cyg X-1



$a^* = 0.77 \pm 0.05$



GRS 1915+105



$a^* = 0.98 - 1.0$

XTE J1650-500

XTE J1118+480

XTE J1859+226

GRS 1009-45

GRS 1124-683

SAX J1819.3-2525

GS 2000+25

H1705-250

GRO J1655-40

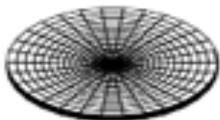
A0620-00

GRO J0422+32

XTE J1550-564

$a^* = 0.65 - 0.75$

$(a^* \sim 0.5)$



GS 2023+338



GS 1354-64

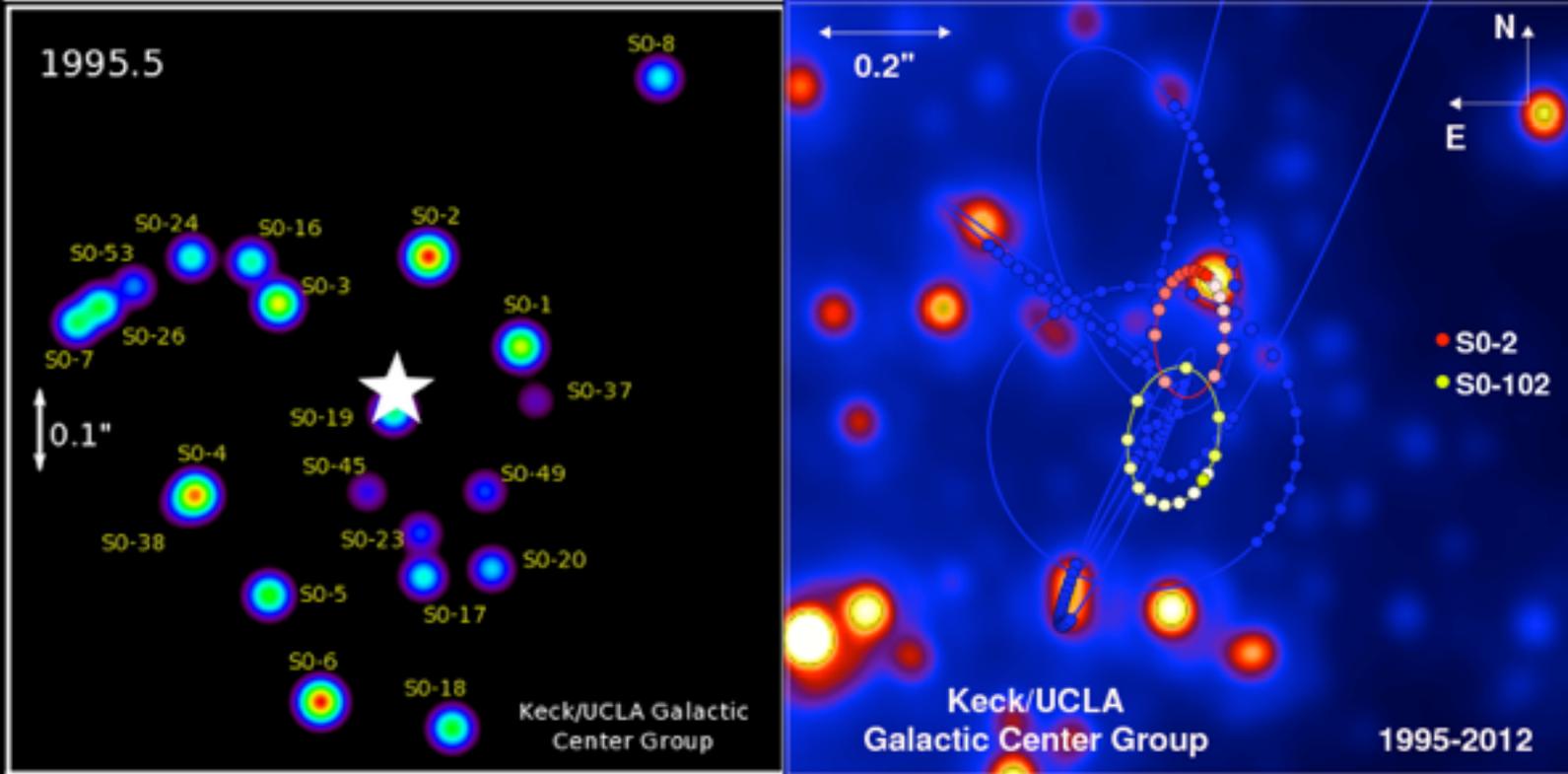


GX 339-4



4U 1543-47

$a^* = 0.75 - 0.85$



Sgr A\*

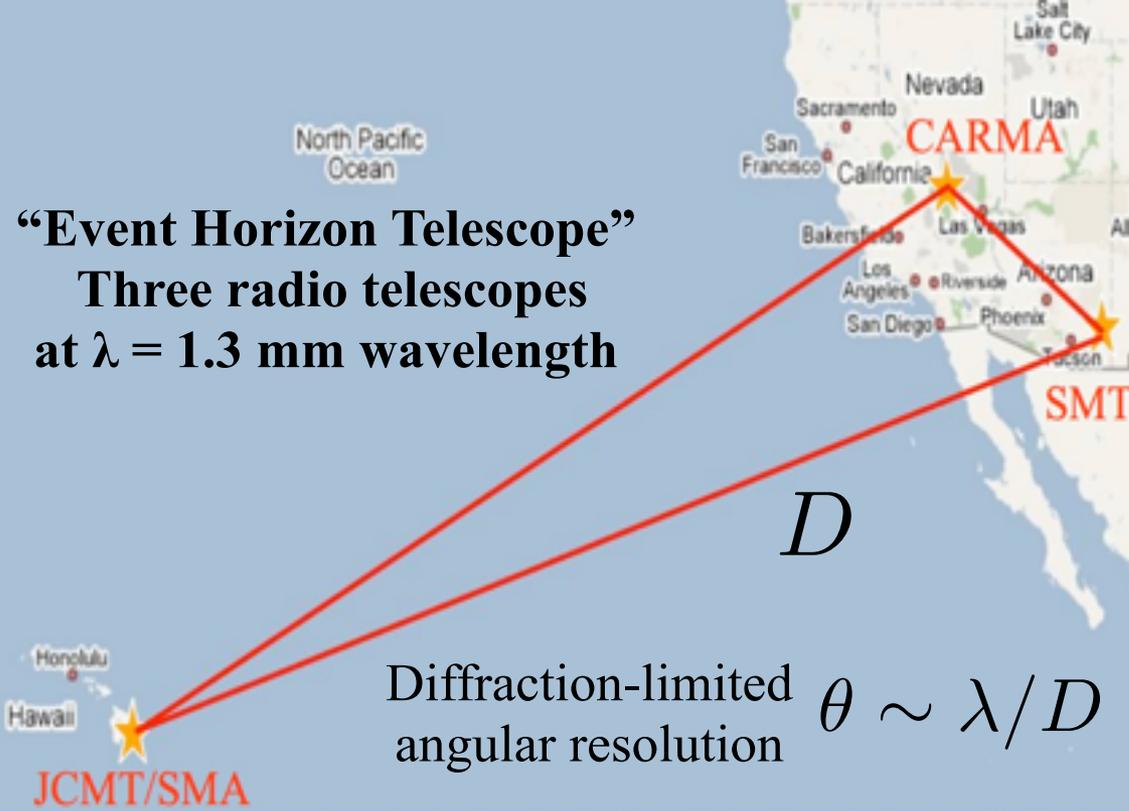
= The Supermassive Black Hole at  
the Galactic Center

$$M \approx 4 \times 10^6 M_{\odot}$$

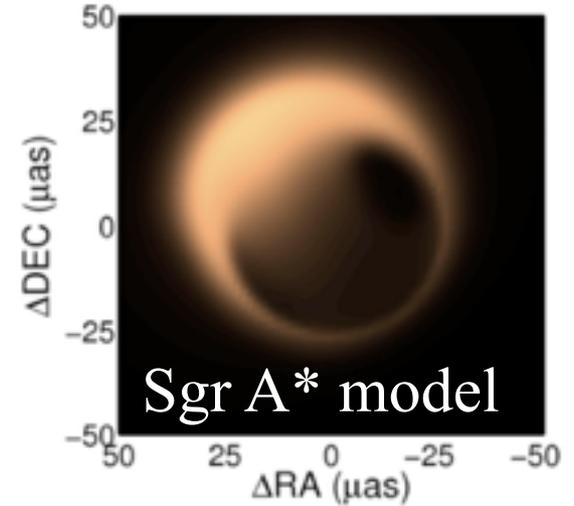
$$R_g \approx 0.04 \text{ AU}$$

$$q_{\text{star}} \approx 0.01 \text{ arcsec} \times 8 \text{ kpc} \approx 80 \text{ AU}$$

“Event Horizon Telescope”  
 Three radio telescopes  
 at  $\lambda = 1.3$  mm wavelength



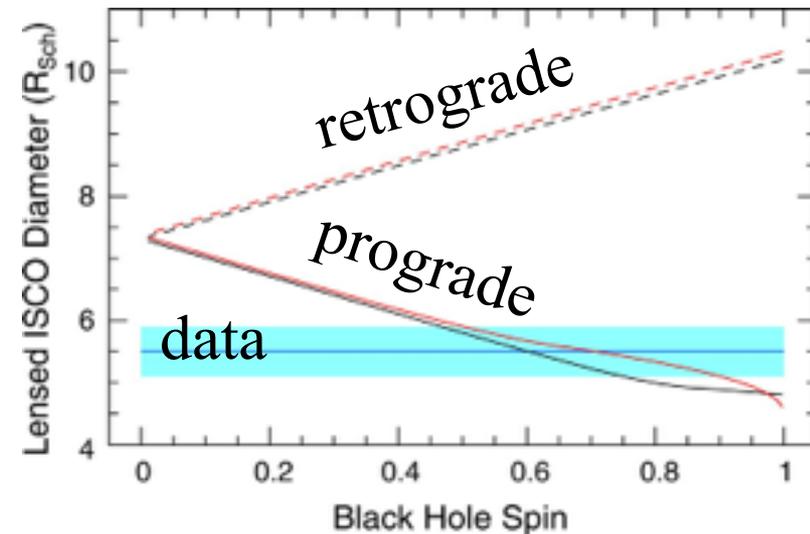
“The data presented here confirm structure in Sgr A\* on linear scales of  $4R_{Sch}$ , but the exact nature of this structure is not well determined.”

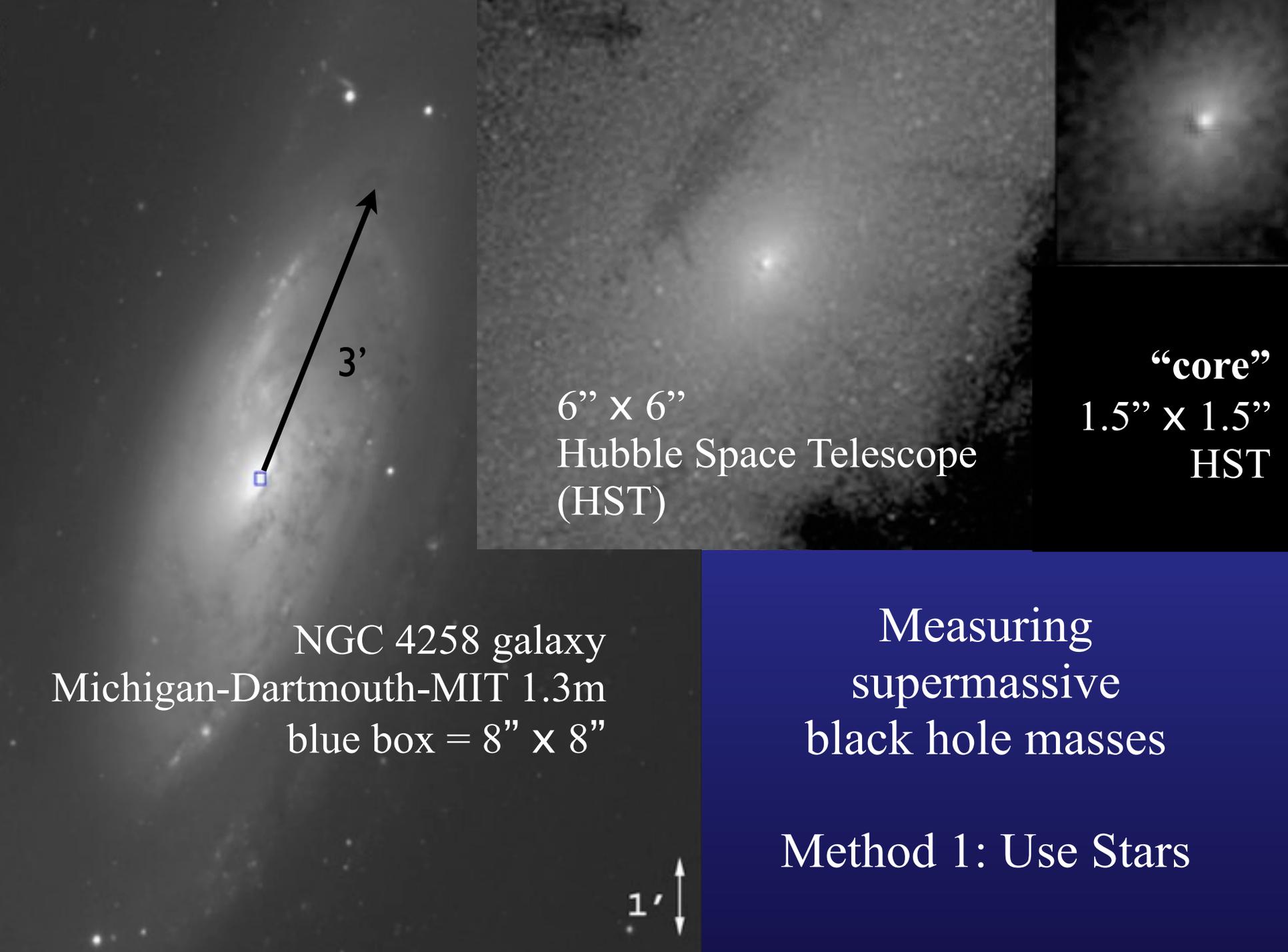


M87 galaxy  
 in optical



M87 with  
 “EHT”  
 (model)



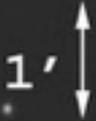


6'' x 6''

Hubble Space Telescope  
(HST)

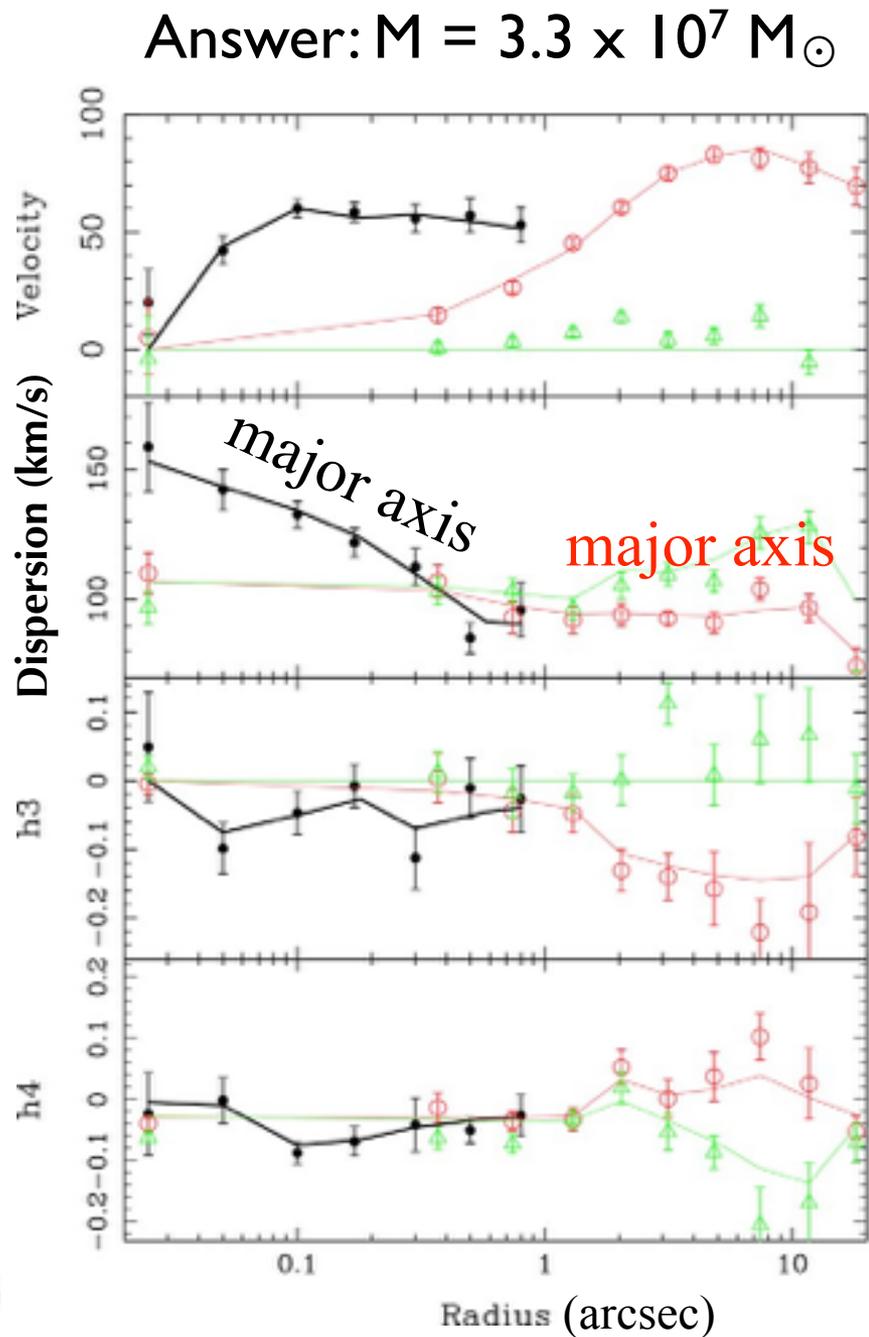
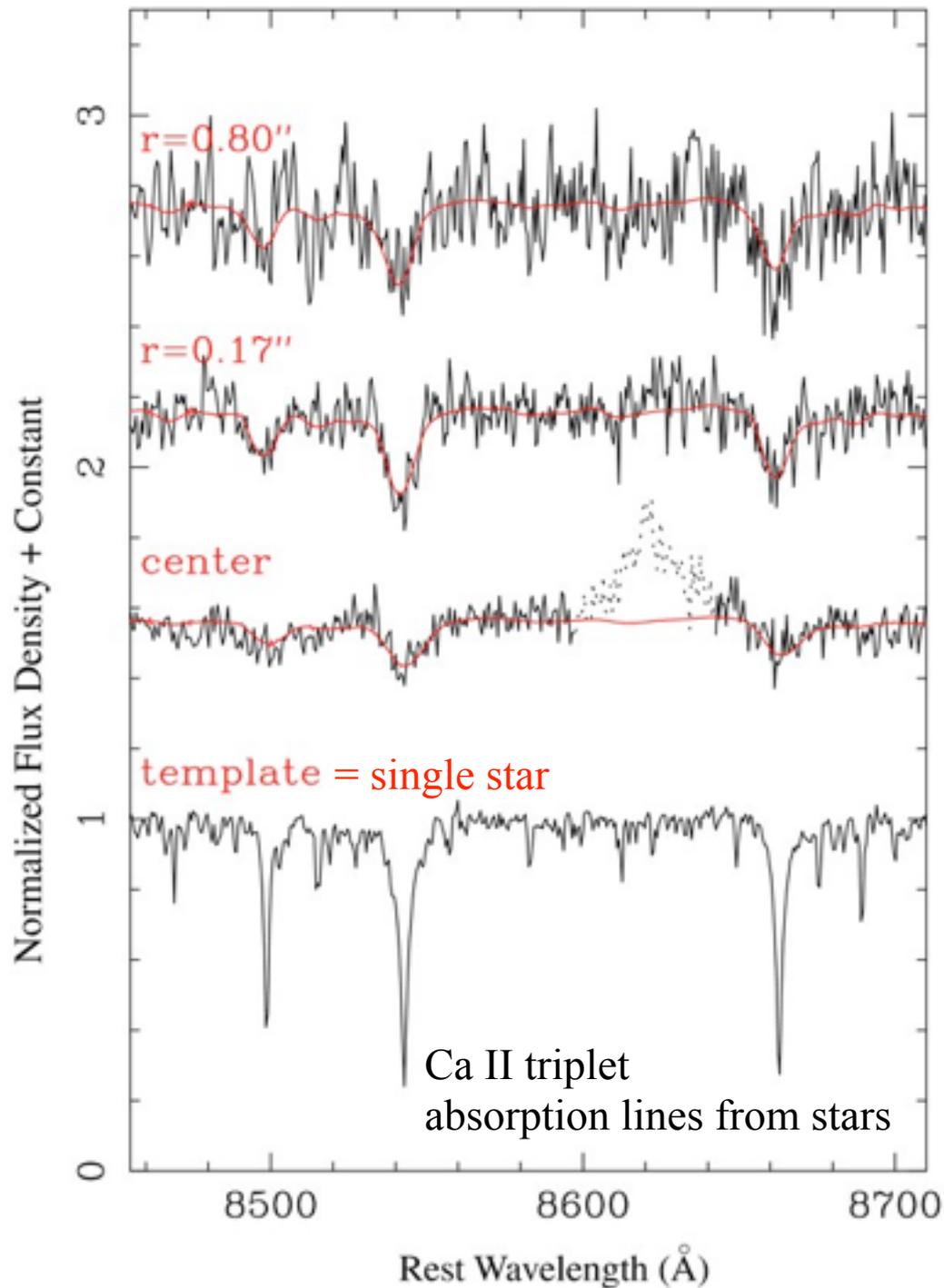
**“core”**  
1.5'' x 1.5''  
HST

NGC 4258 galaxy  
Michigan-Dartmouth-MIT 1.3m  
blue box = 8'' x 8''

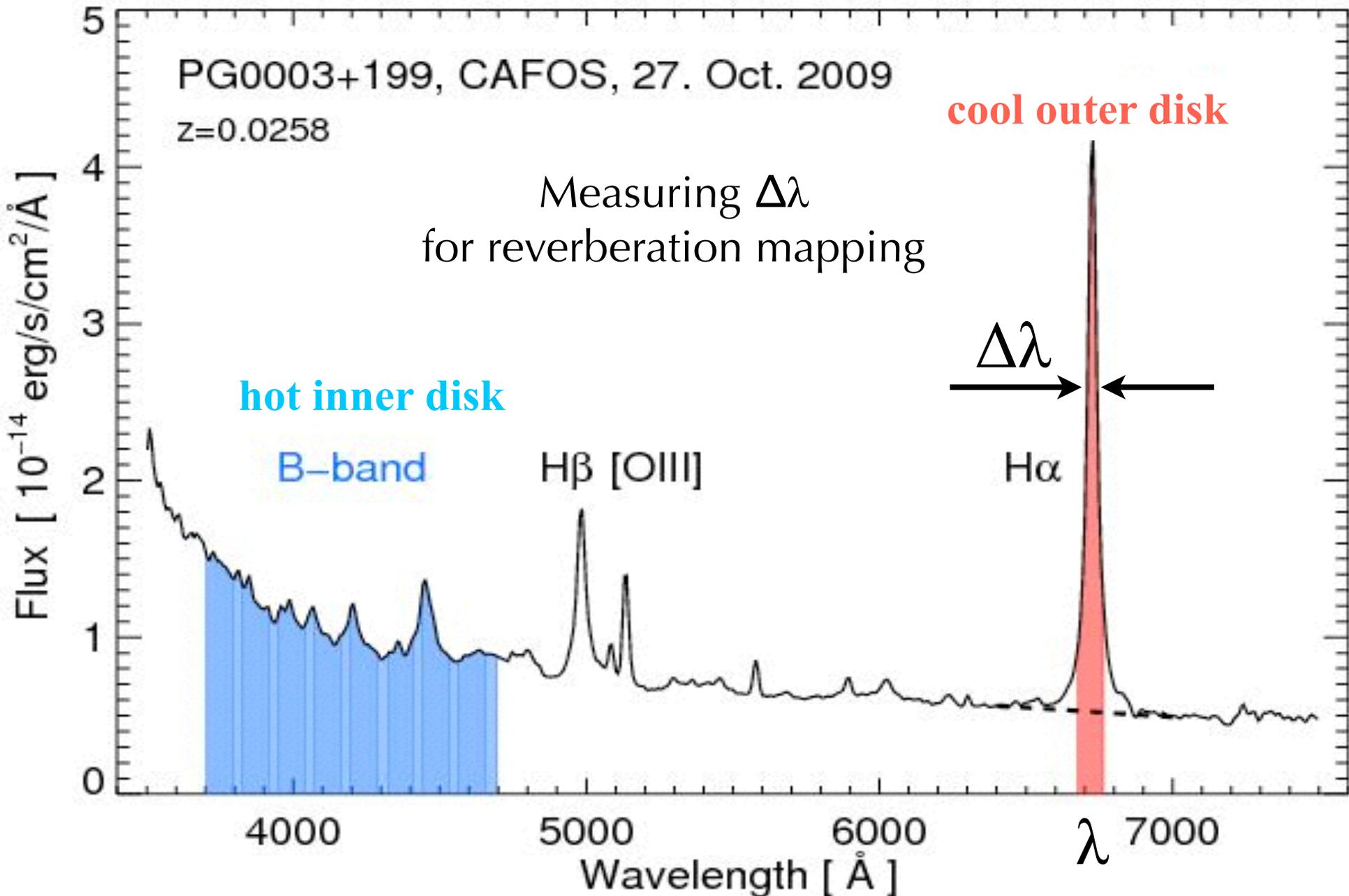


Measuring  
supermassive  
black hole masses

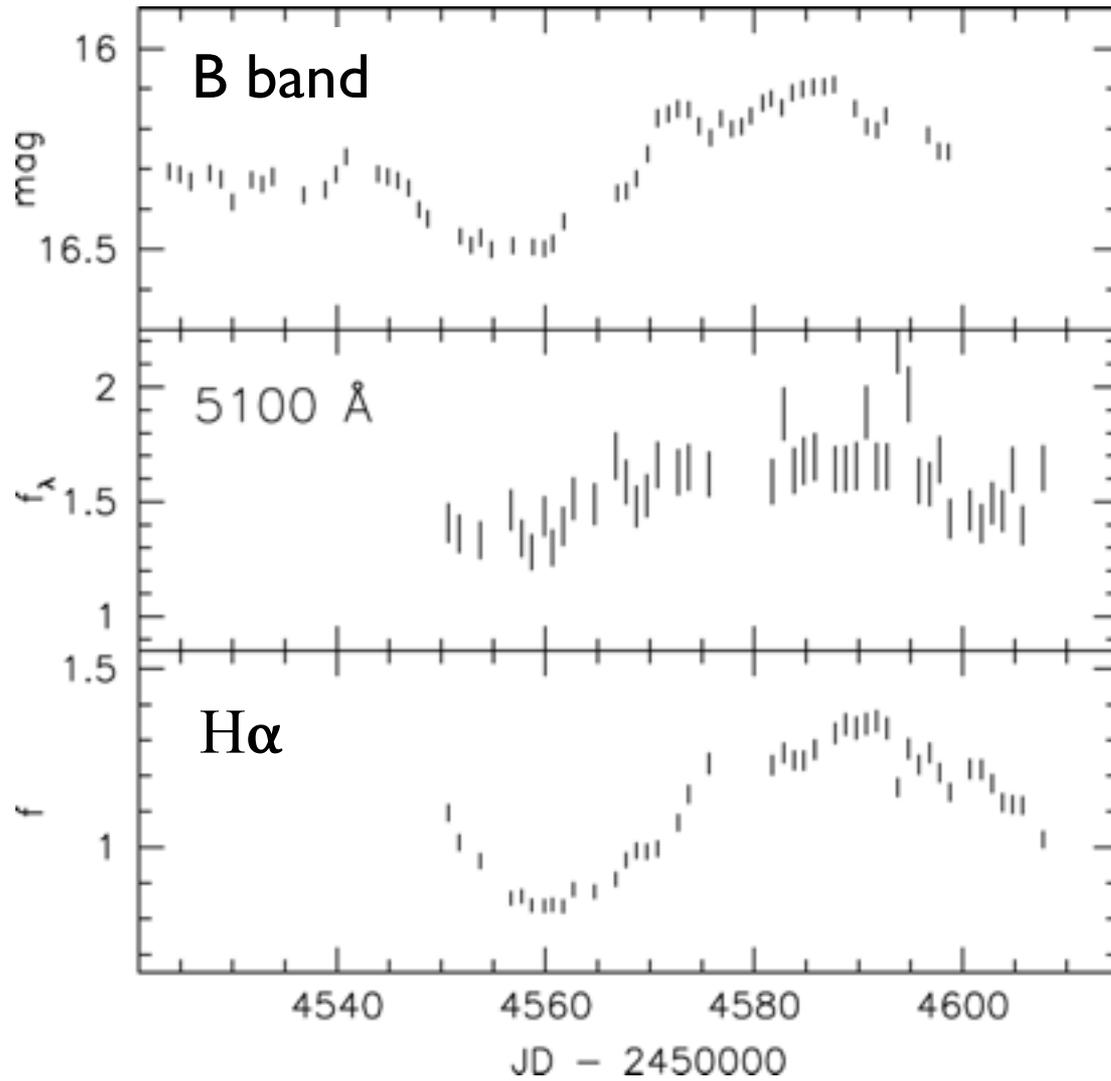
Method 1: Use Stars



# Method 2: Use Accretion Disk (Reverberation Mapping)

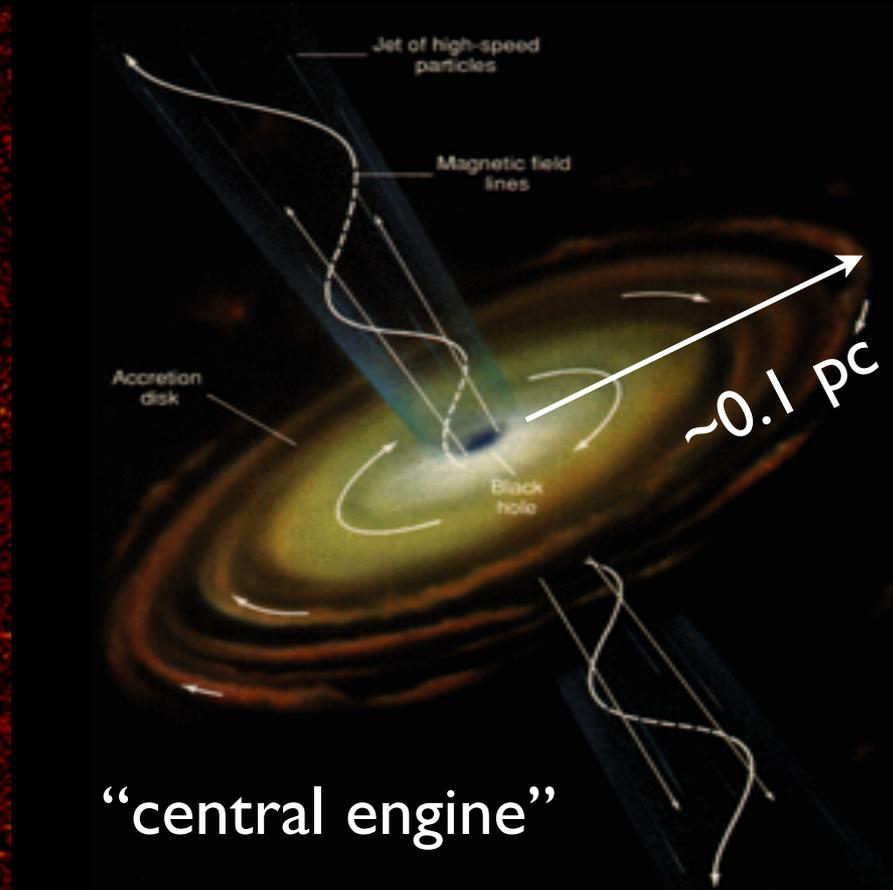
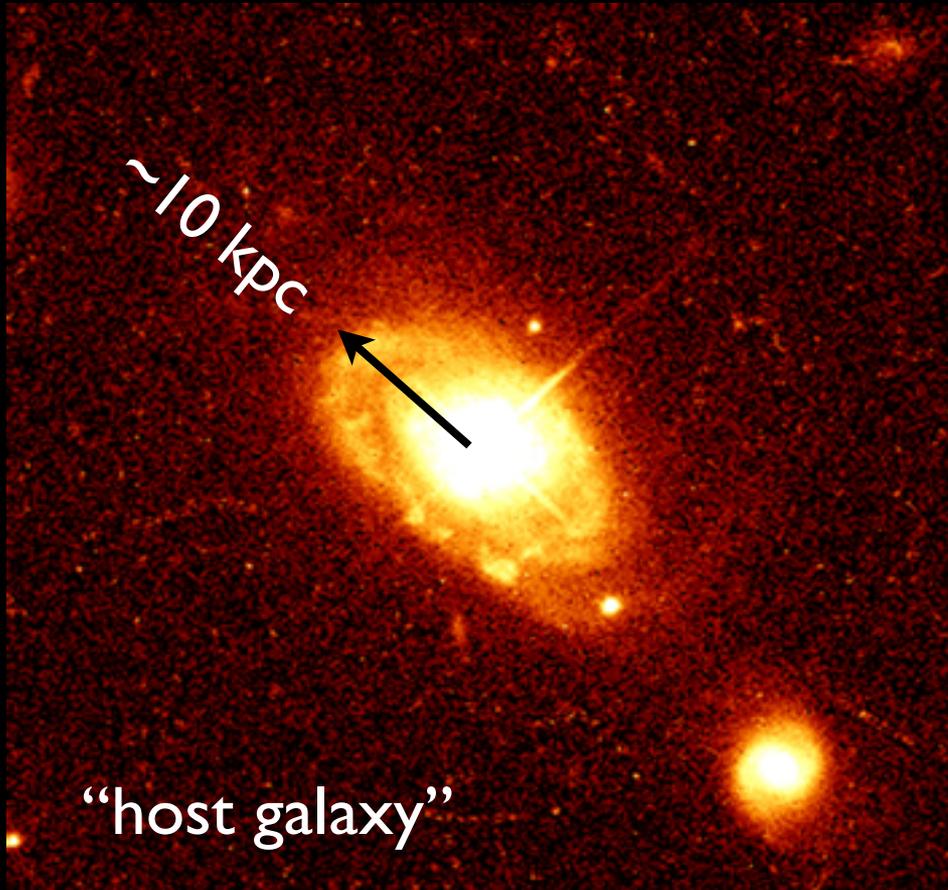


## Method 2: Use Accretion Disk (Reverberation Mapping)

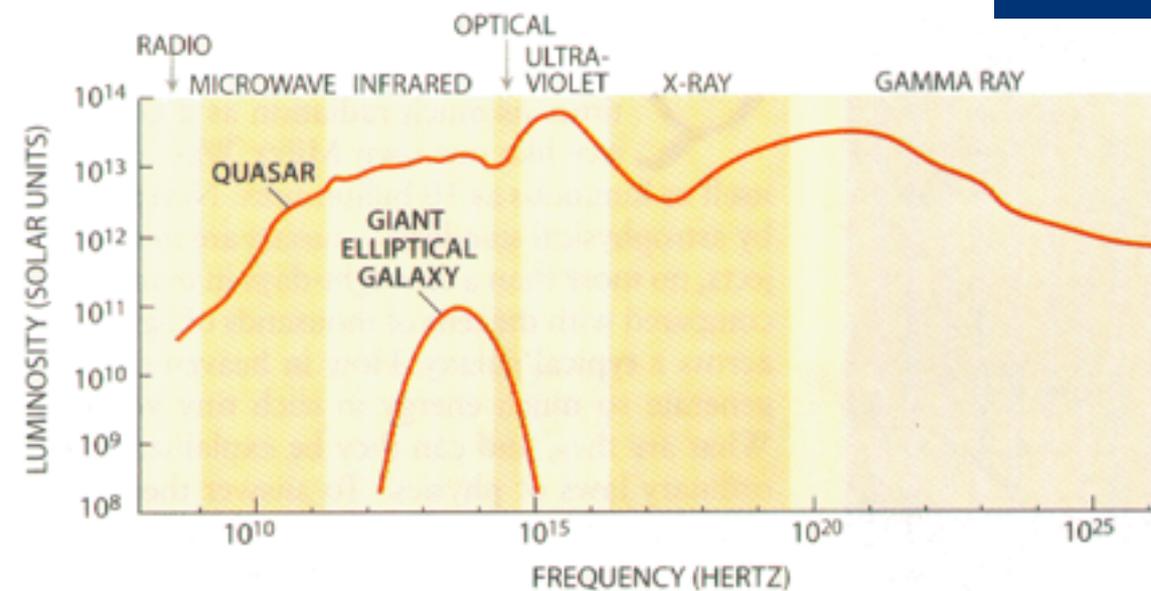
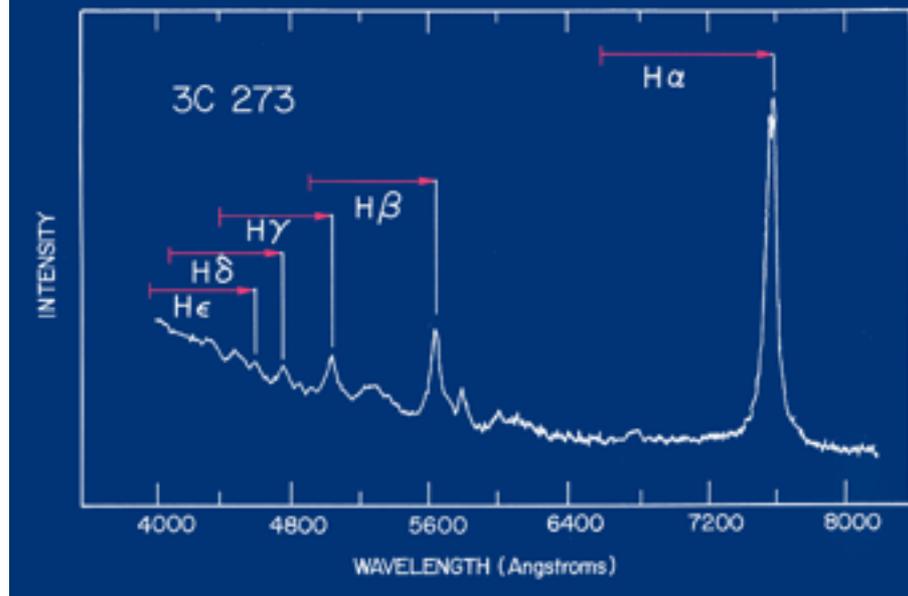
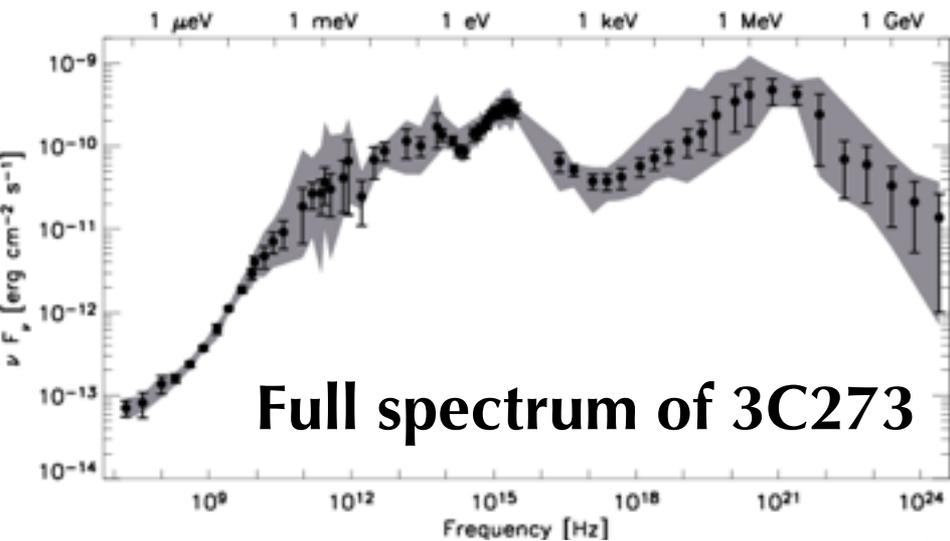


$$\Delta t_{\text{lag}} \sim 4 \text{ days}$$

# Quasars: the most luminous objects in the universe



Quasars = accreting supermassive black holes  
in the nuclei of galaxies

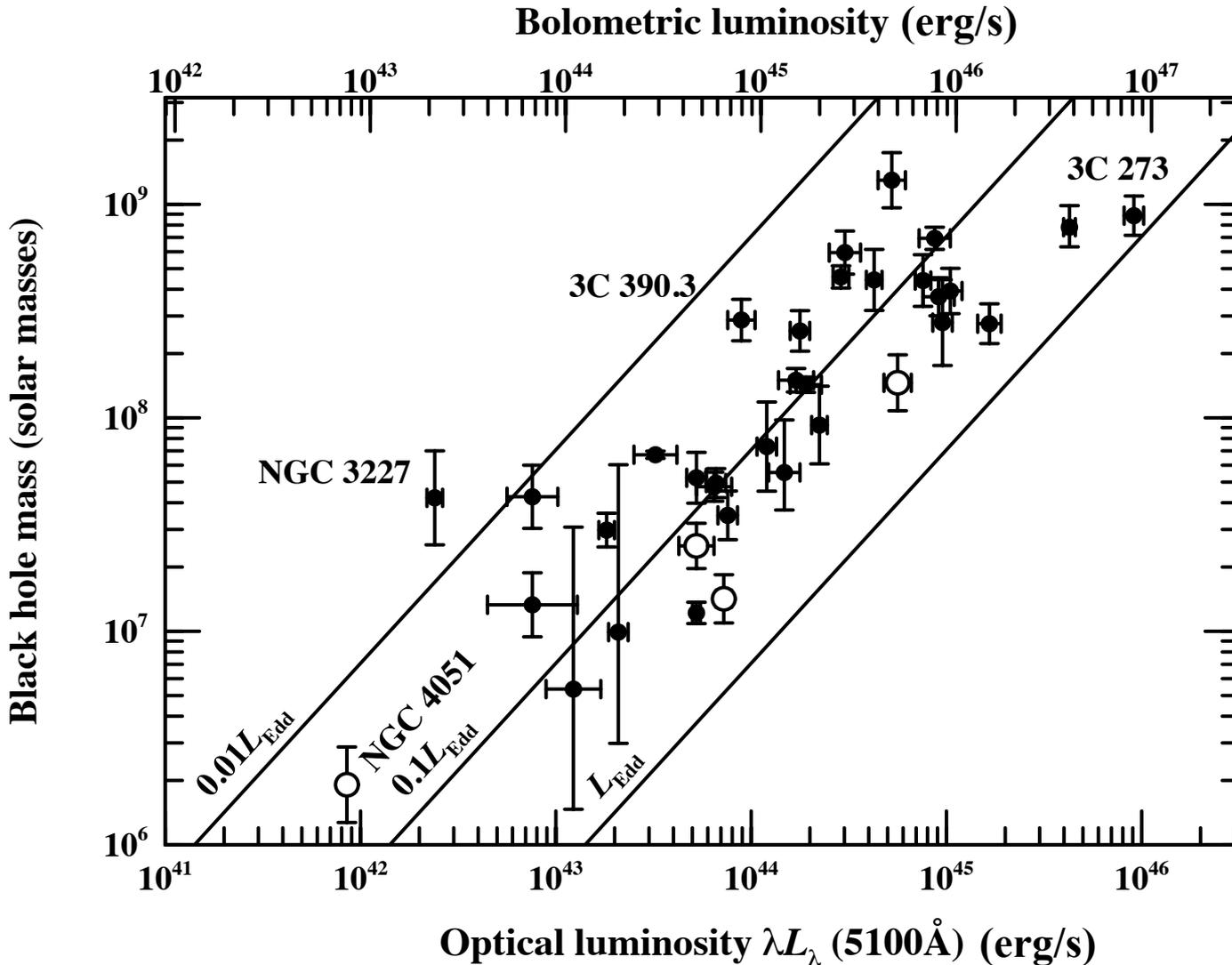


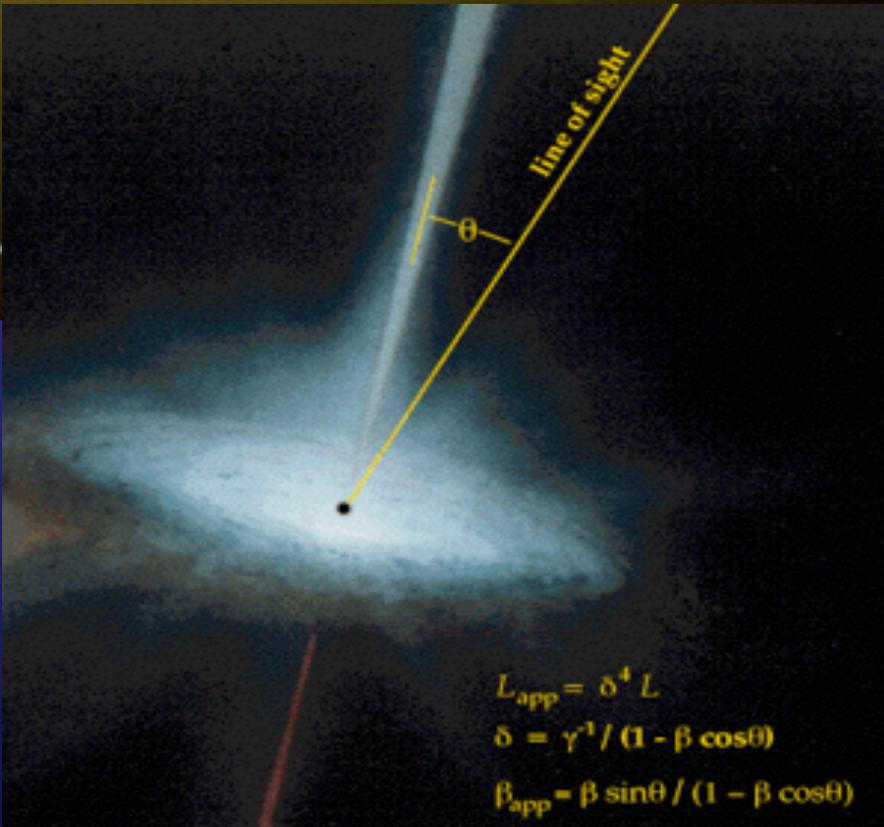
QUASAR SPECTRUM of 3C273—one of the brightest quasars and the first to be discovered—is far broader than the spectrum of a typical giant elliptical galaxy (*left*). In the optical range, the quasar is hundreds of times more luminous. Quasars were most numerous when the universe was two to four billion years old

**Luminosity of 3C273**  
 $\sim 10^{47}$  erg/s

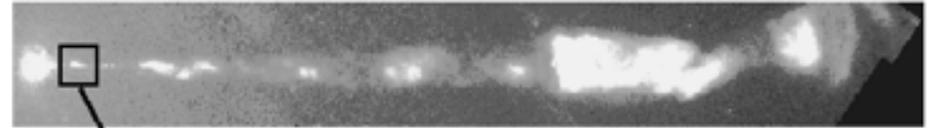
**Luminosity of host galaxy**  
 $\sim 10^{44}$  erg/s

# Quasar Masses from Reverberation Mapping

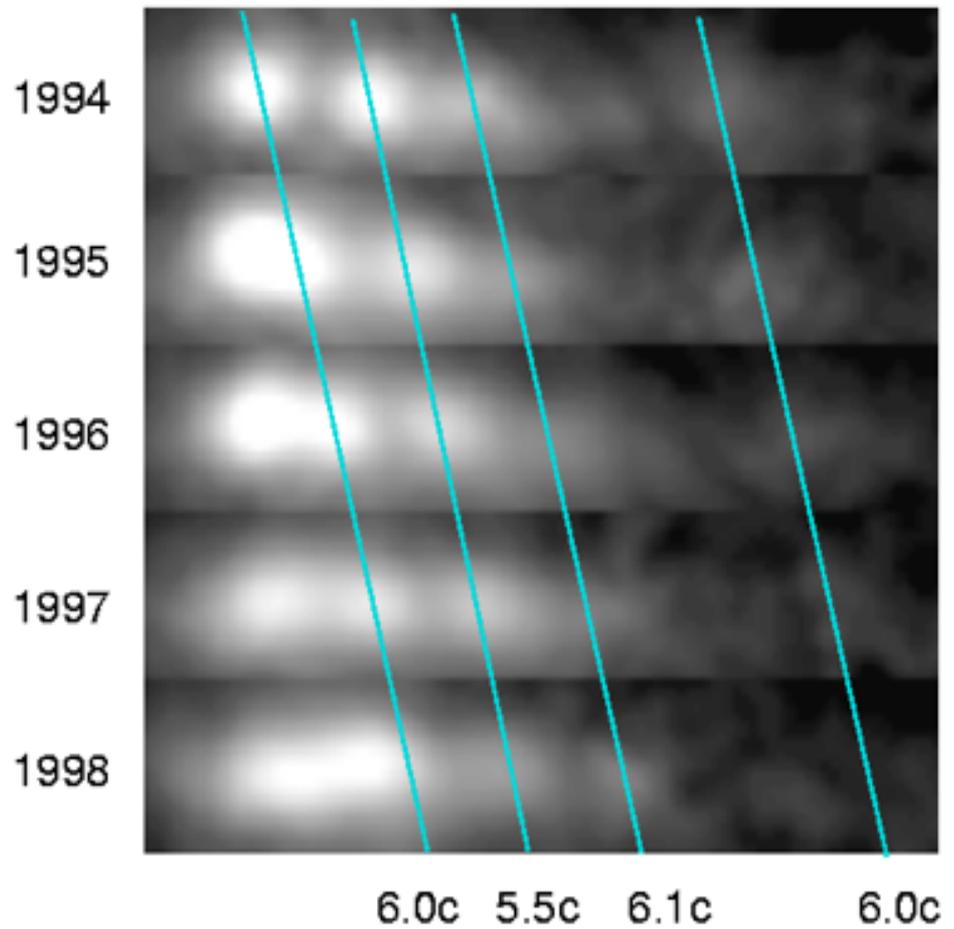




## Superluminal Motion in the M87 Jet

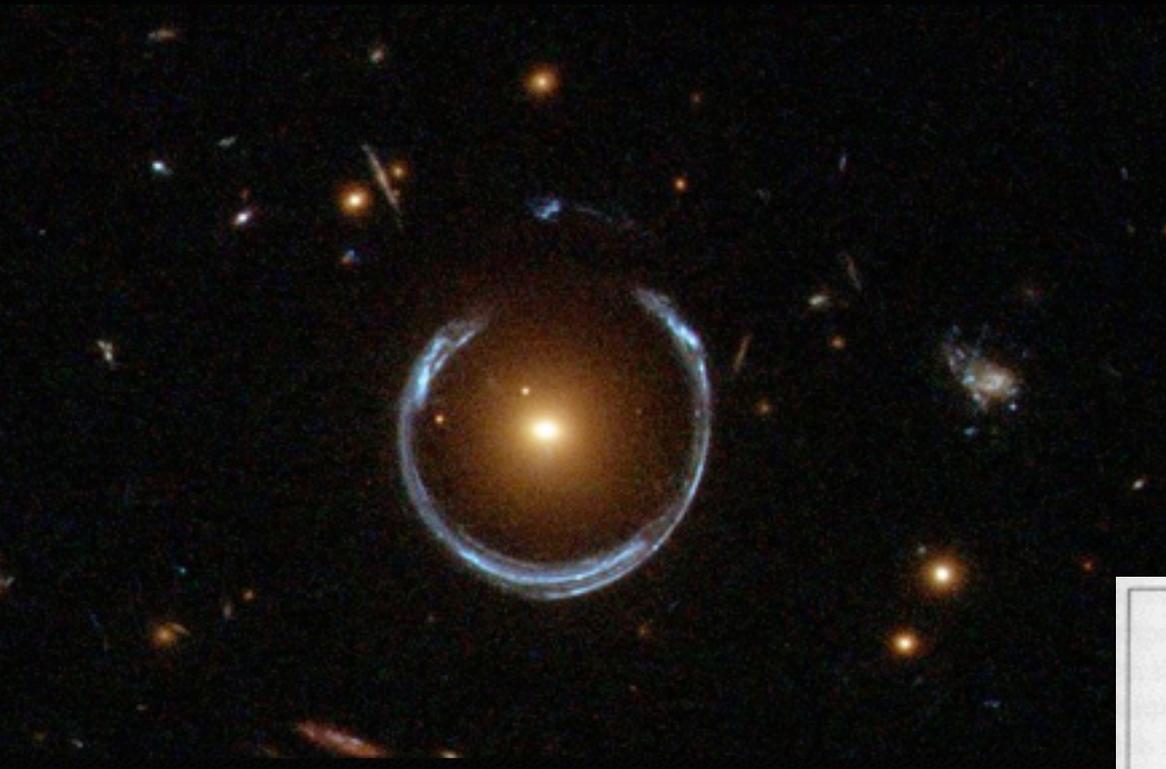


proper motion  $\mu \sim 25$  milliarcsecond/yr  
distance to M87 =  $d = 16$  Mpc  
apparent transverse velocity  $v = \mu d = 6c > c!$



# Gravitational Lensing

Resolved “strong”  
lensing



Gravitational Lens in Abell 2218

HST · WFPC2

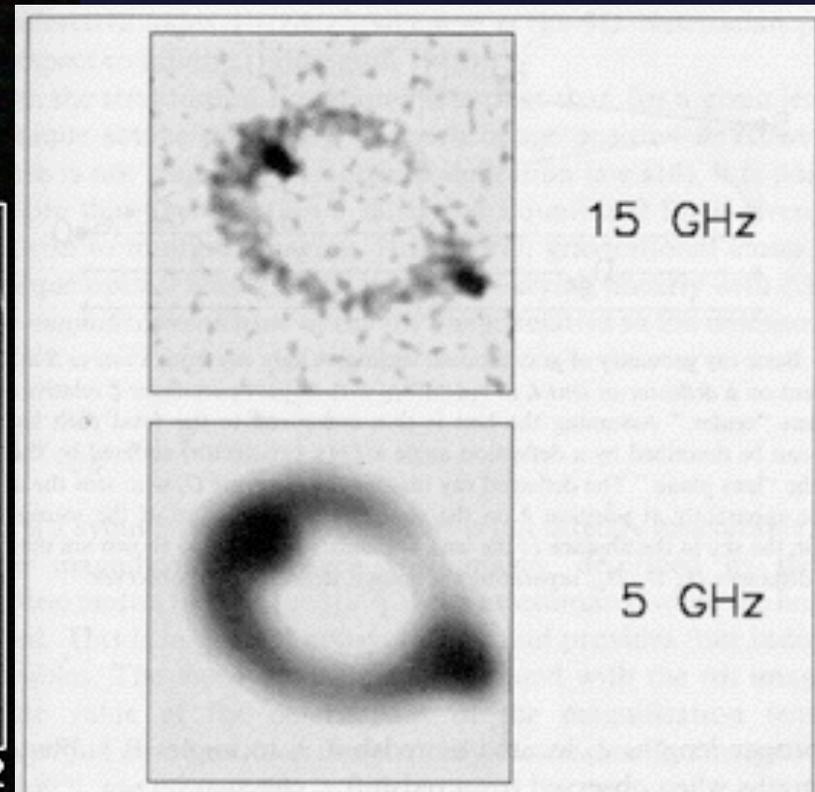
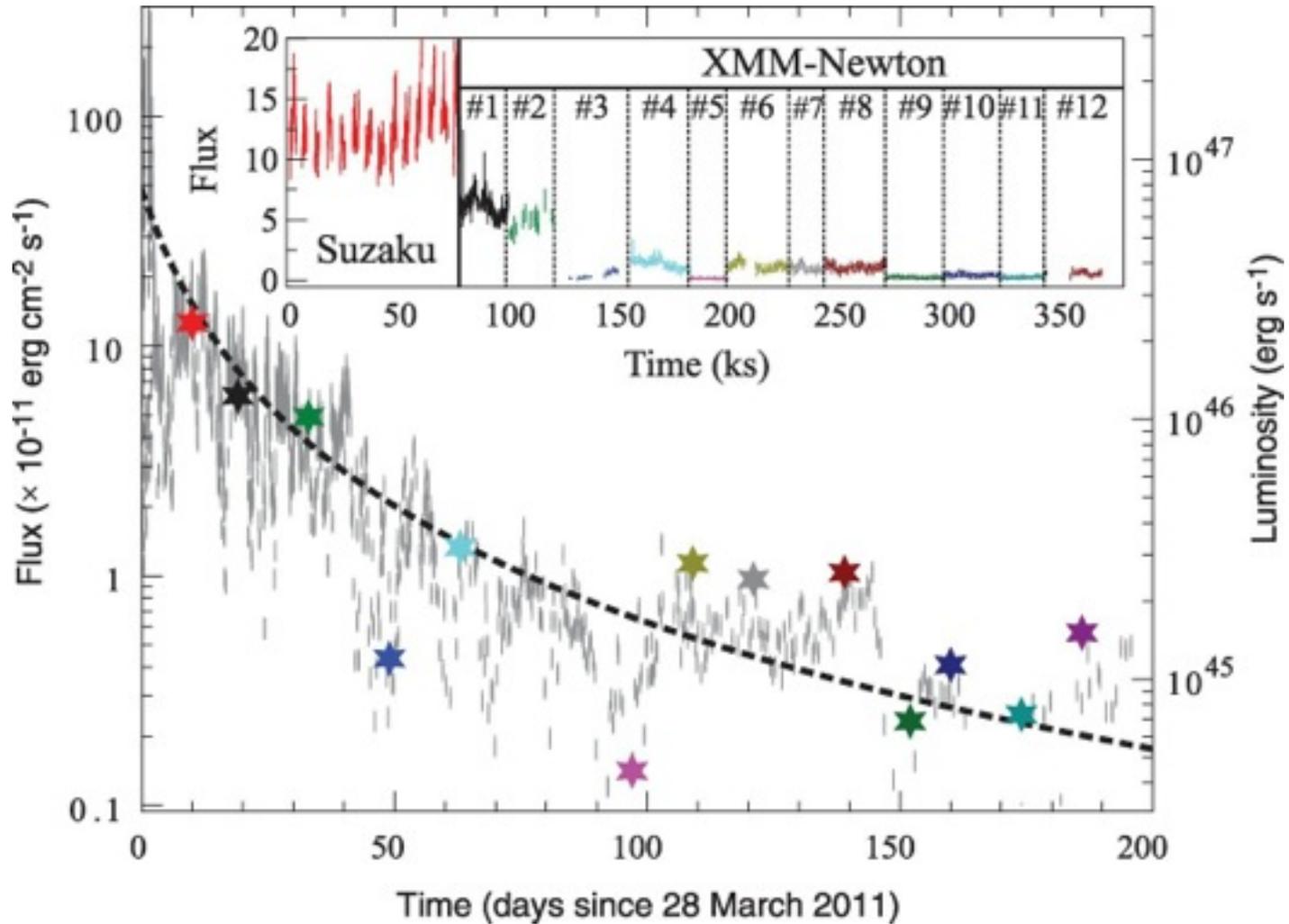


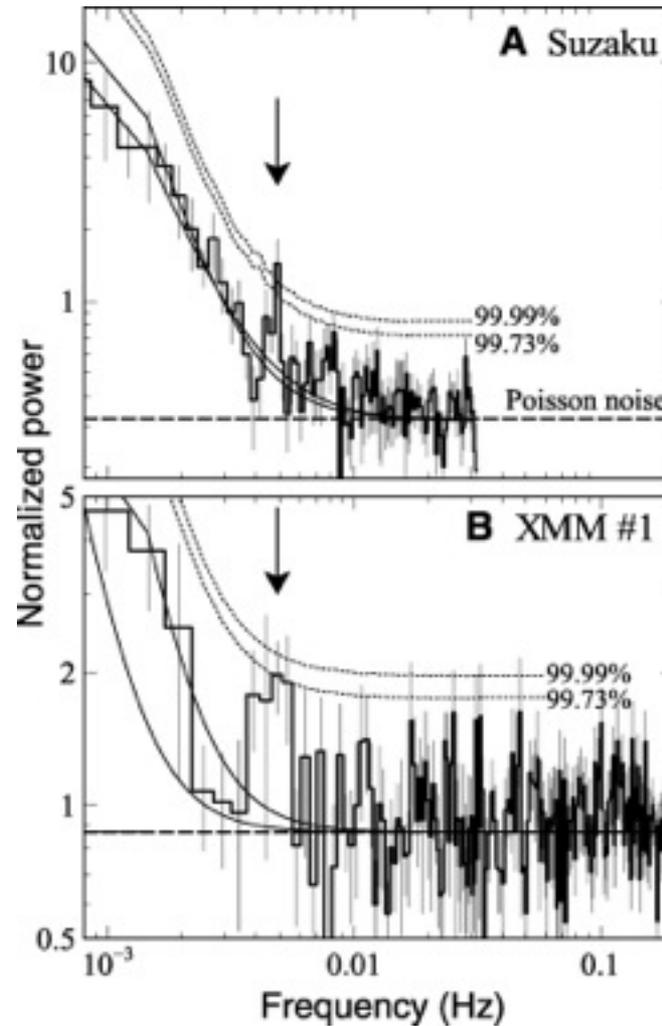
Fig. 1 XMM-Newton and Suzaku light curve of Sw J1644+57 together with the SWIFT-XRT 0.3-to-10-keV light curve for reference (gray) (37, 38).



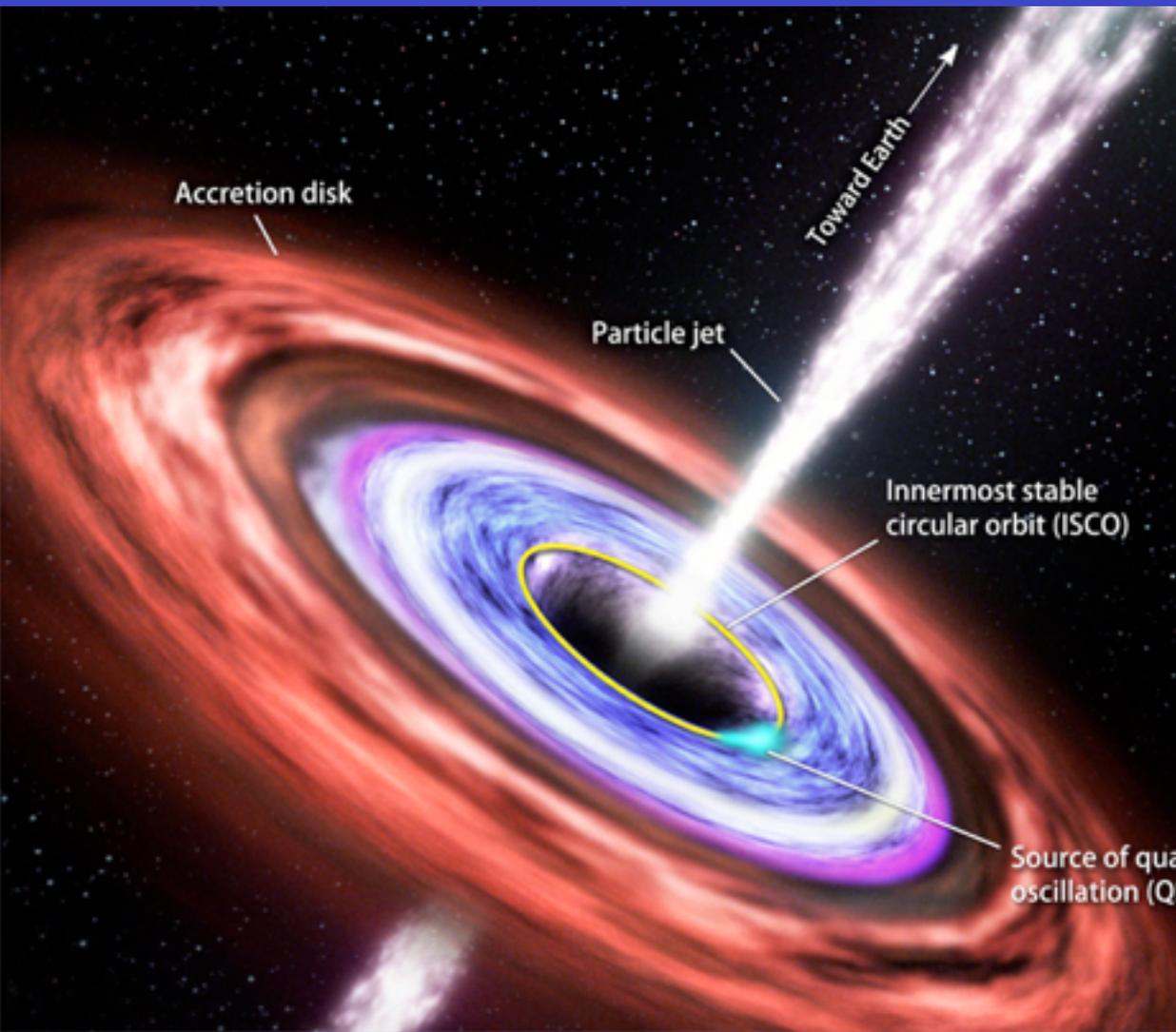
R C Reis et al. Science 2012;337:949-951



Fig. 2 (A) Power spectra for Suzaku and (B) for XMM #1, in the 2- to 10-keV energy range.



R C Reis et al. Science 2012;337:949-951



## Stats for Swift J1644+57

**Discovery** March 28, 2011, by NASA's Swift

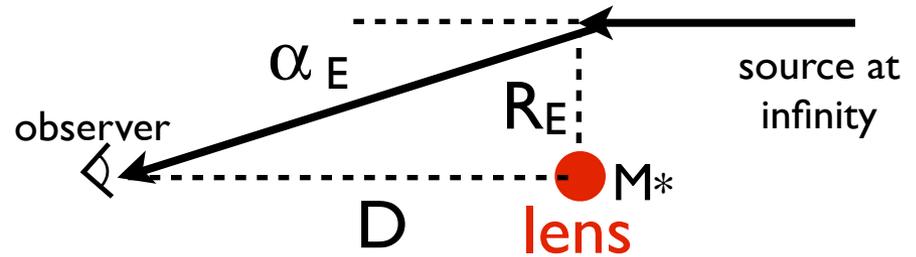
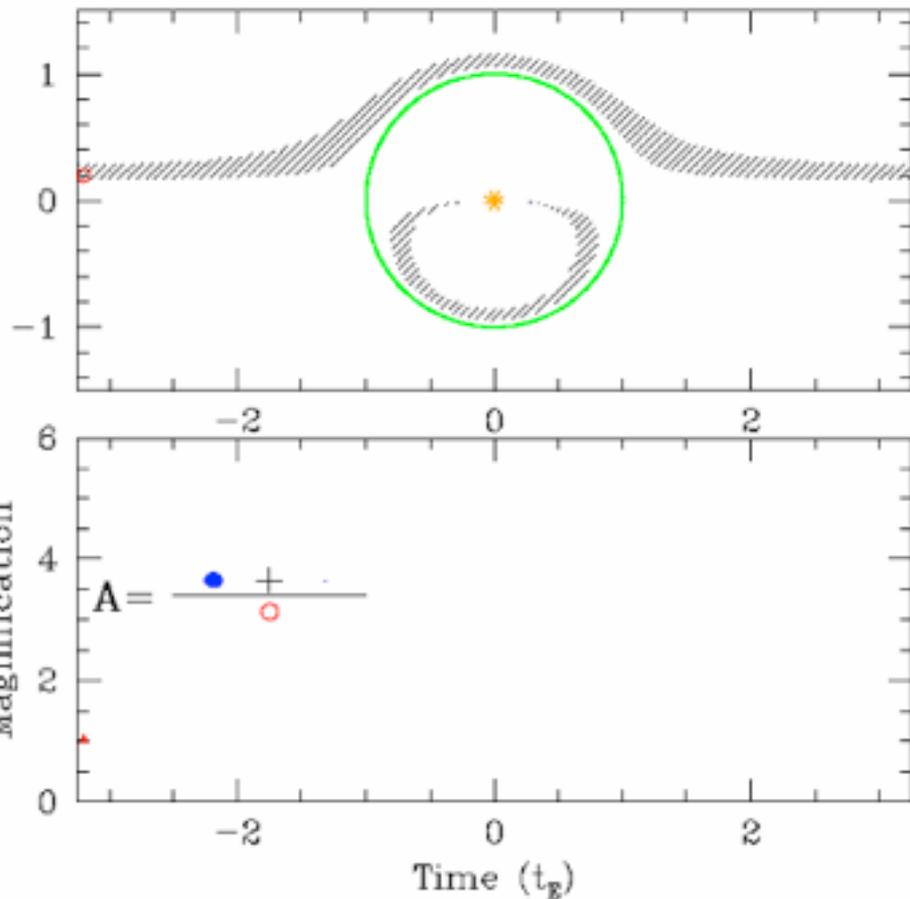
**Description** Supermassive black hole activated by tidal breakup of a passing star

**Mass** 450,000 to 5 million times the sun's mass

**Distance** 3.9 billion light-years

**QPO** Period of about 3.5 minutes, frequency of 4.8 millihertz. Most distant ever detected.

**ISCO** 2.5 to 5.8 million miles (4 to 9.3 million km) from the center of the black hole. This range is roughly 3 to 6 times the sun's diameter.

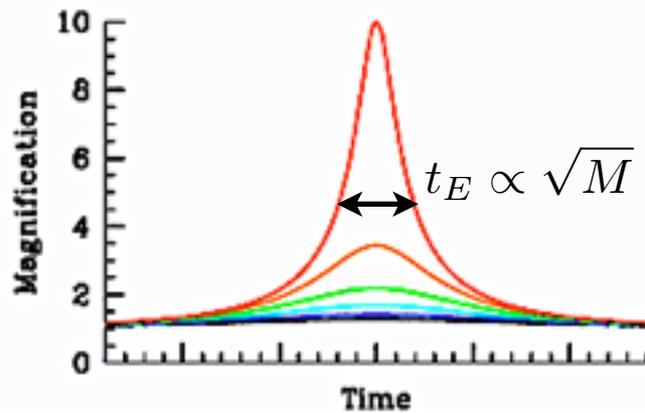
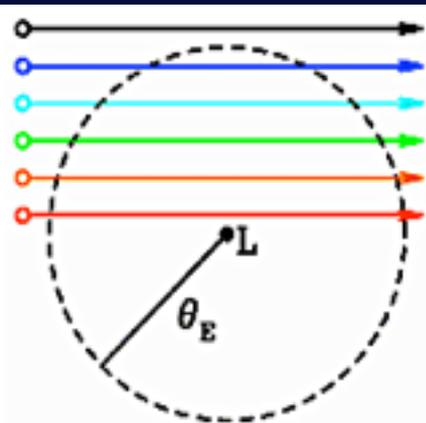


$$\Delta v \sim \frac{GM_*}{R_E^2} \times \frac{2R_E}{c} \quad \alpha_E \sim \frac{\Delta v}{c} \sim \frac{R_E}{D}$$

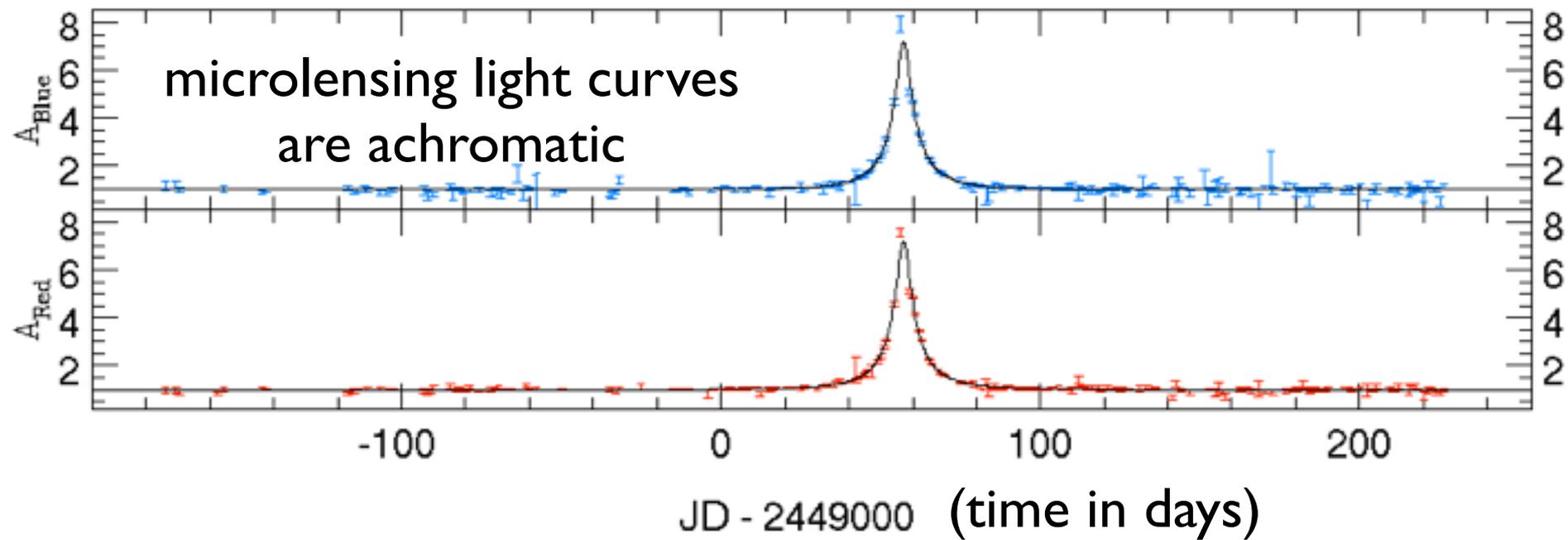
$$\alpha_E \sim 1 \text{ mas} \left( \frac{M_*}{M_\odot} \right)^{1/2} \left( \frac{\text{kpc}}{D} \right)^{1/2}$$

$$R_E \sim 1 \text{ AU} \left( \frac{M_*}{M_\odot} \right)^{1/2} \left( \frac{D}{\text{kpc}} \right)^{1/2}$$

$$t_E \sim \frac{R_E}{v_{\text{rel}}} \sim 1 \text{ month} \left( \frac{R_E}{\text{AU}} \right) \left( \frac{30 \text{ km/s}}{v_{\text{rel}}} \right)$$



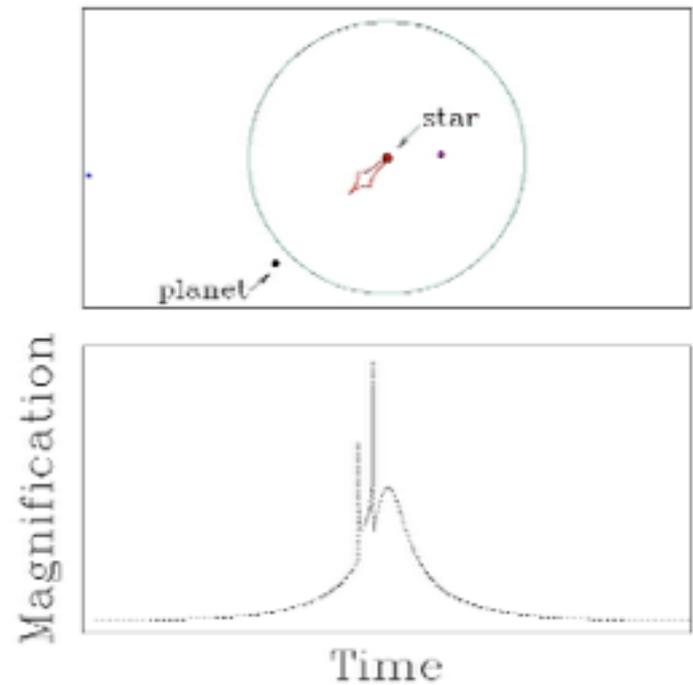
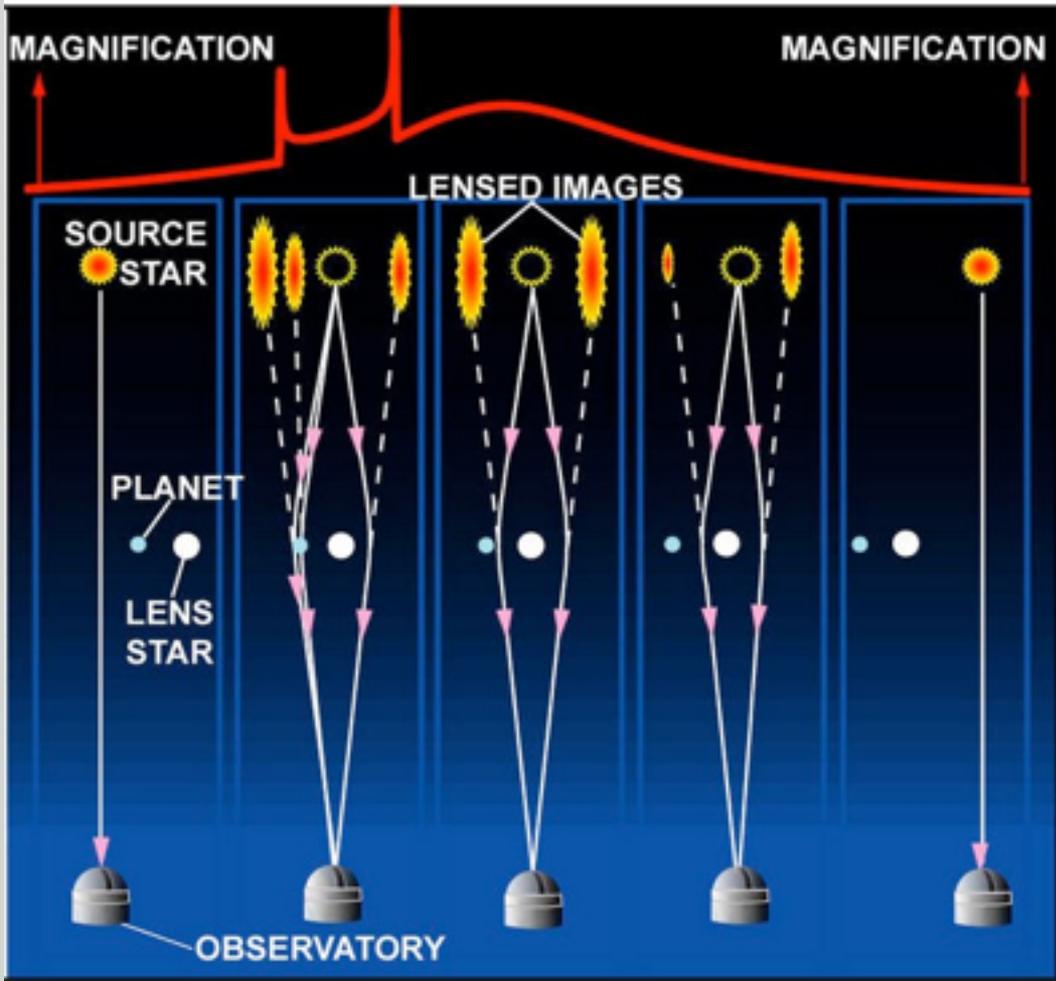
**“Microlensing”**  
 = unresolved strong  
 lensing  
 = measure light curve  
 only



Microlensing surveys stare at  
dense star fields

Large Magellanic Cloud  
(LMC = satellite galaxy of Milky Way)

Galactic Bulge  
(of our Milky Way)



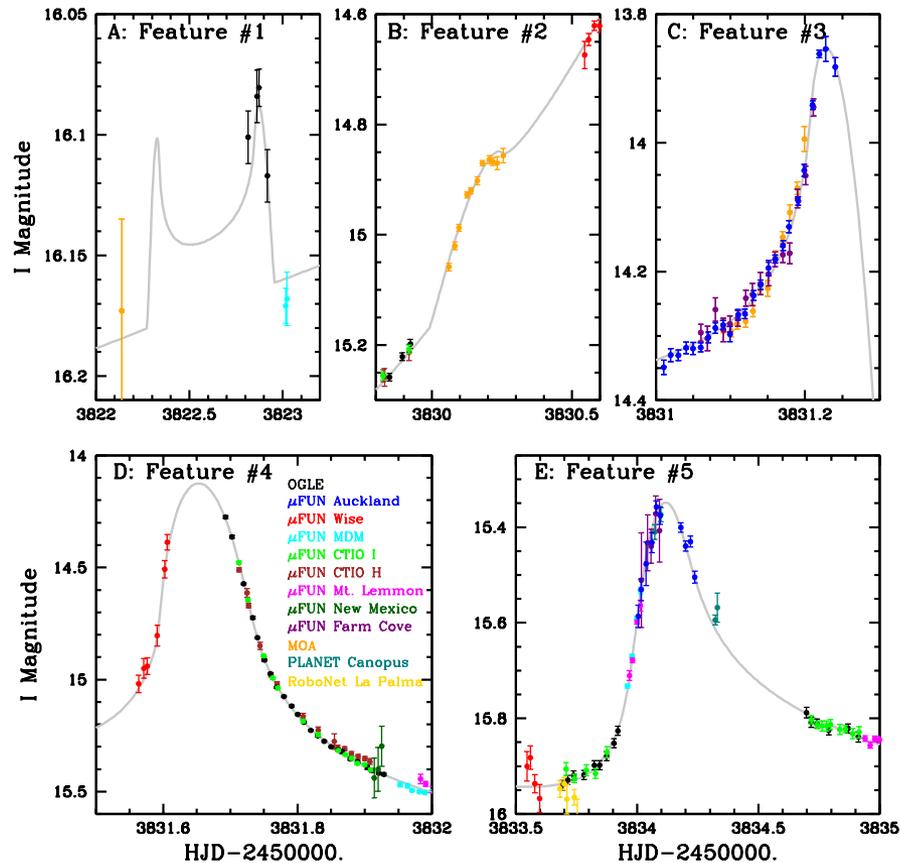
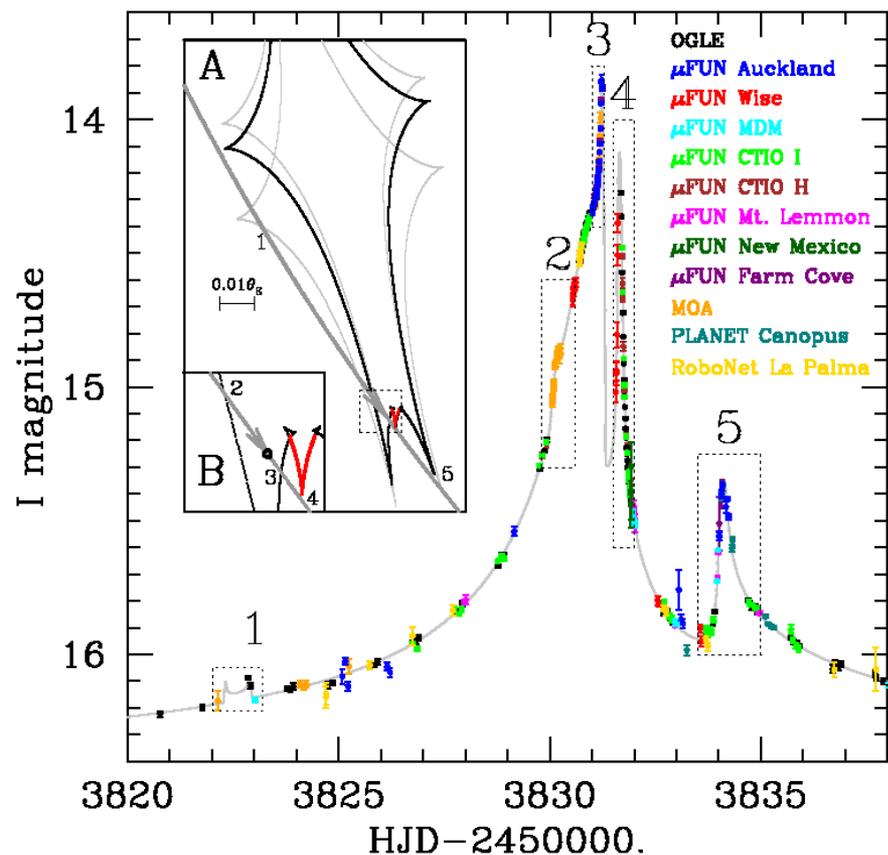
## Planetary Microlensing (Binary lens)

$$D \approx 5.2_{-2.9}^{+0.2} \text{ kpc}$$

$$a \approx 3.0_{-1.7}^{+0.1} \text{ AU}$$

$$M_* \approx 0.36_{-0.28}^{+0.03} M_{\odot}$$

$$M \approx 1.5_{-1.2}^{+0.1} M_J$$



$$M_* \approx 0.5M_\odot \quad D \approx 1.5 \text{ kpc}$$

$$M_1 \approx 0.71M_J \quad a_1 \approx 2.3 \text{ AU}$$

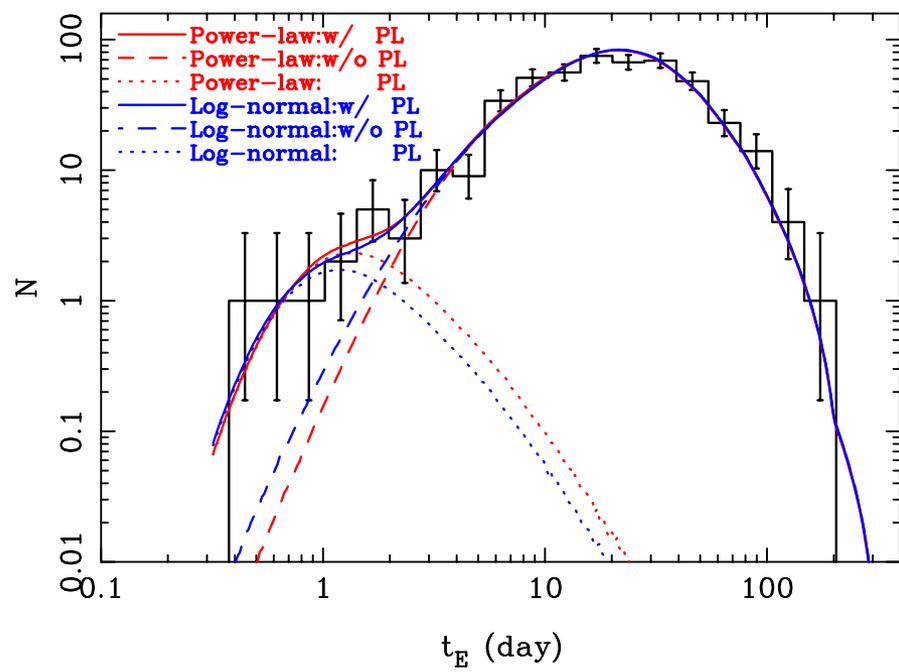
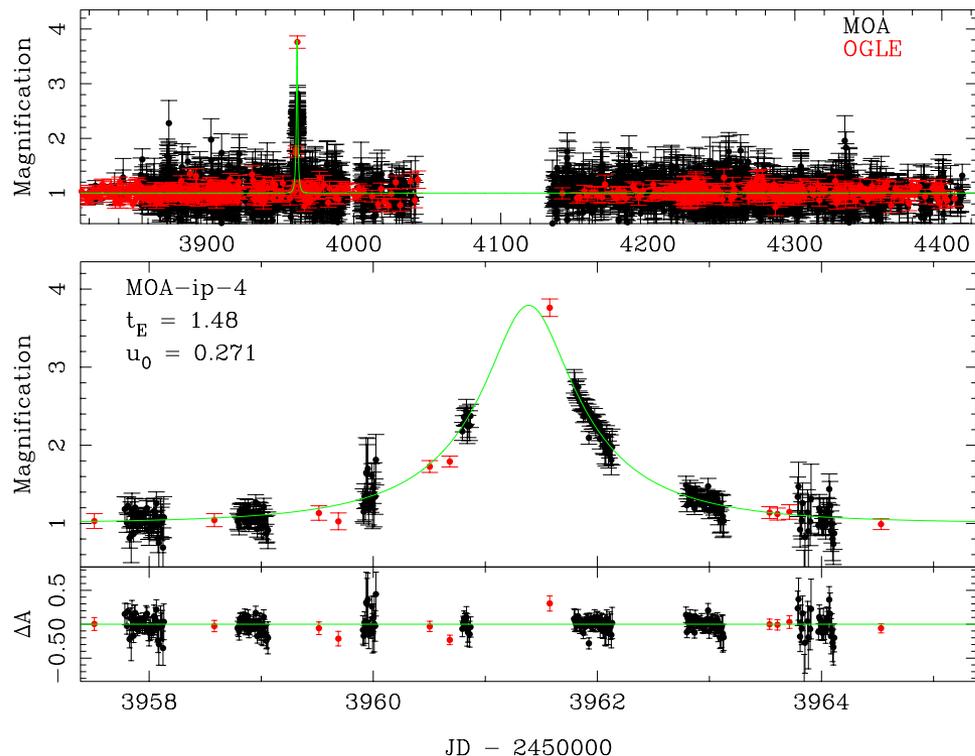
$$M_2 \approx 0.27M_J \quad a_2 \approx 4.6 \text{ AU}$$

$$e_2 \approx 0.15^{+0.17}_{-0.10}$$

from orbital  
motion!

# Statistics from microlensing

1. At fixed  $M_*$ ,  $dN/dM \propto M^{-1.7}$
2. 20 +/- 10% of stars have 5-15  $M_{\oplus}$  at a  $\sim 1.6$ -4.3 AU
3. For every star  $0.08 M_{\odot} < M_* < 1 M_{\odot}$ , there are  $\sim 2$  *free-floating* Jupiter-mass objects!



A Long Time Ago  
In a Disk Galaxy  
Far Far Away

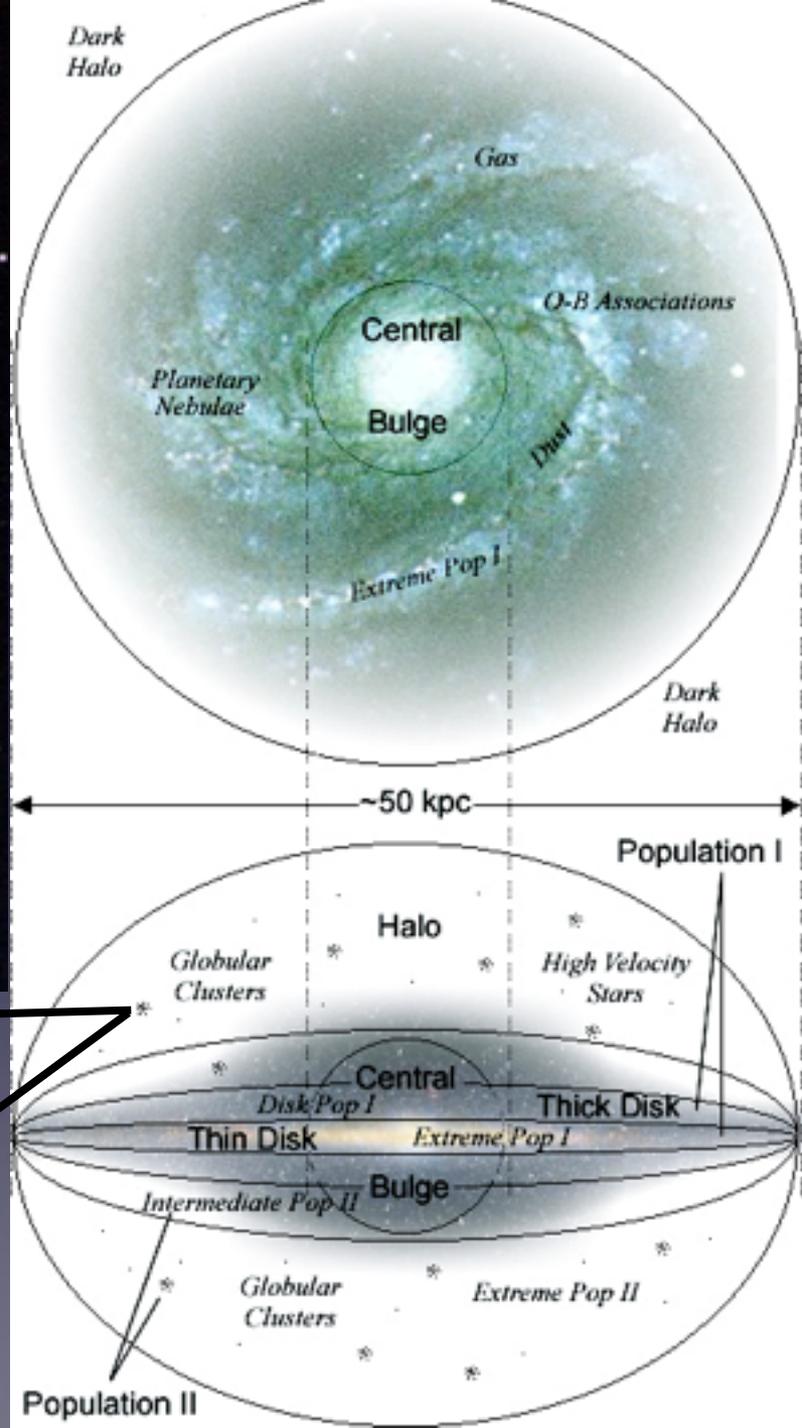




**Stars form in  
interstellar clouds  
(~1 - 100 pc in size)**



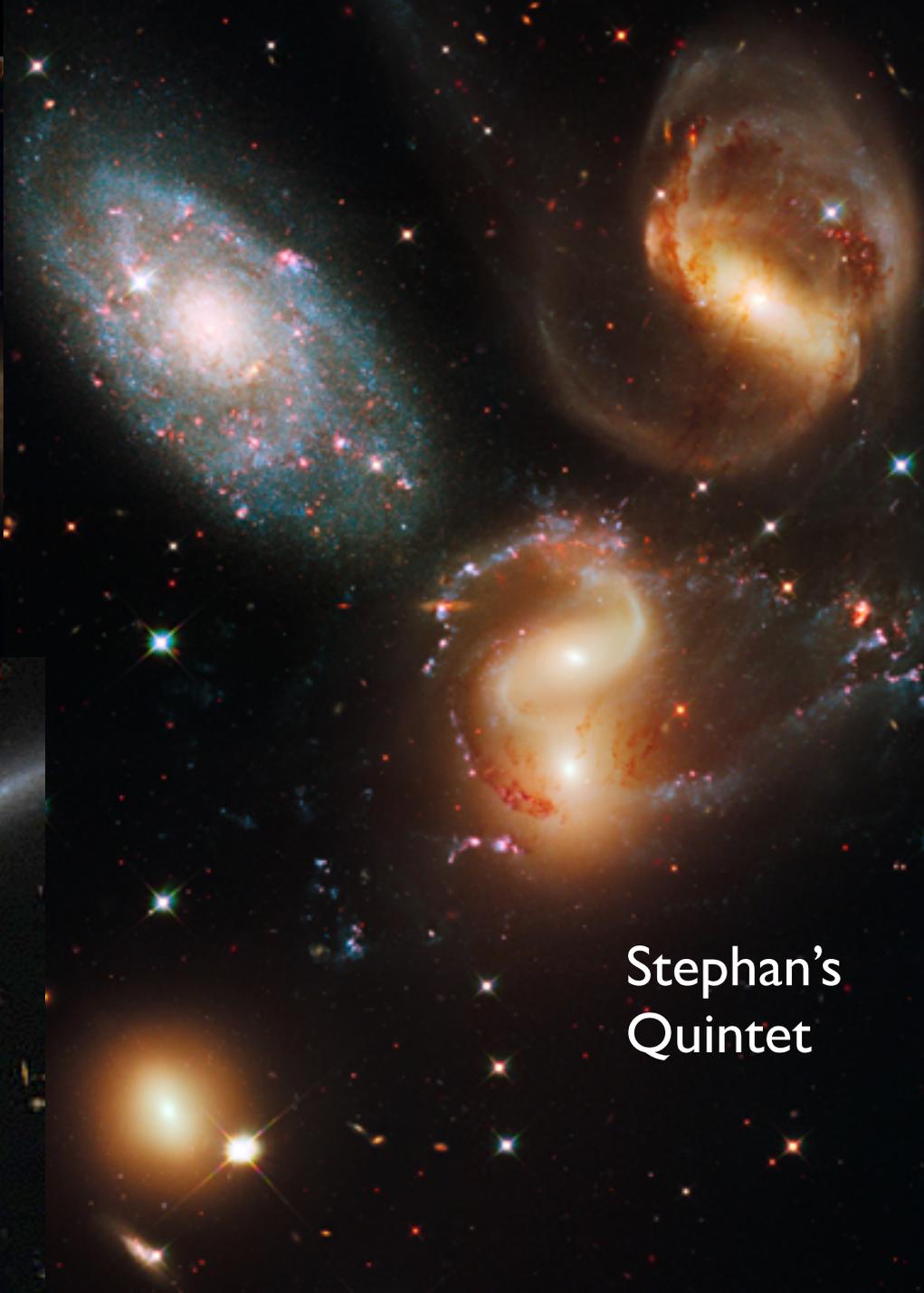
Globular cluster



NGC 2207 and IC 2163



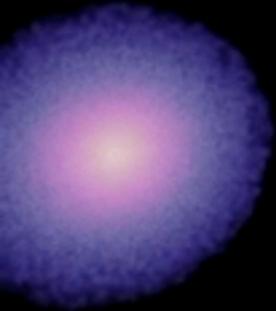
The Mice



Stephan's  
Quintet

T = 0 Myr

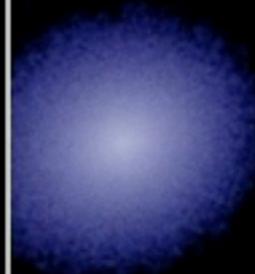
Gas



face on view

T = 0 Myr

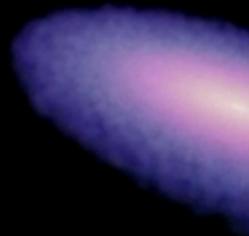
Stars



face on view

T = 0 Myr

Gas



edge on view

T = 0 Myr

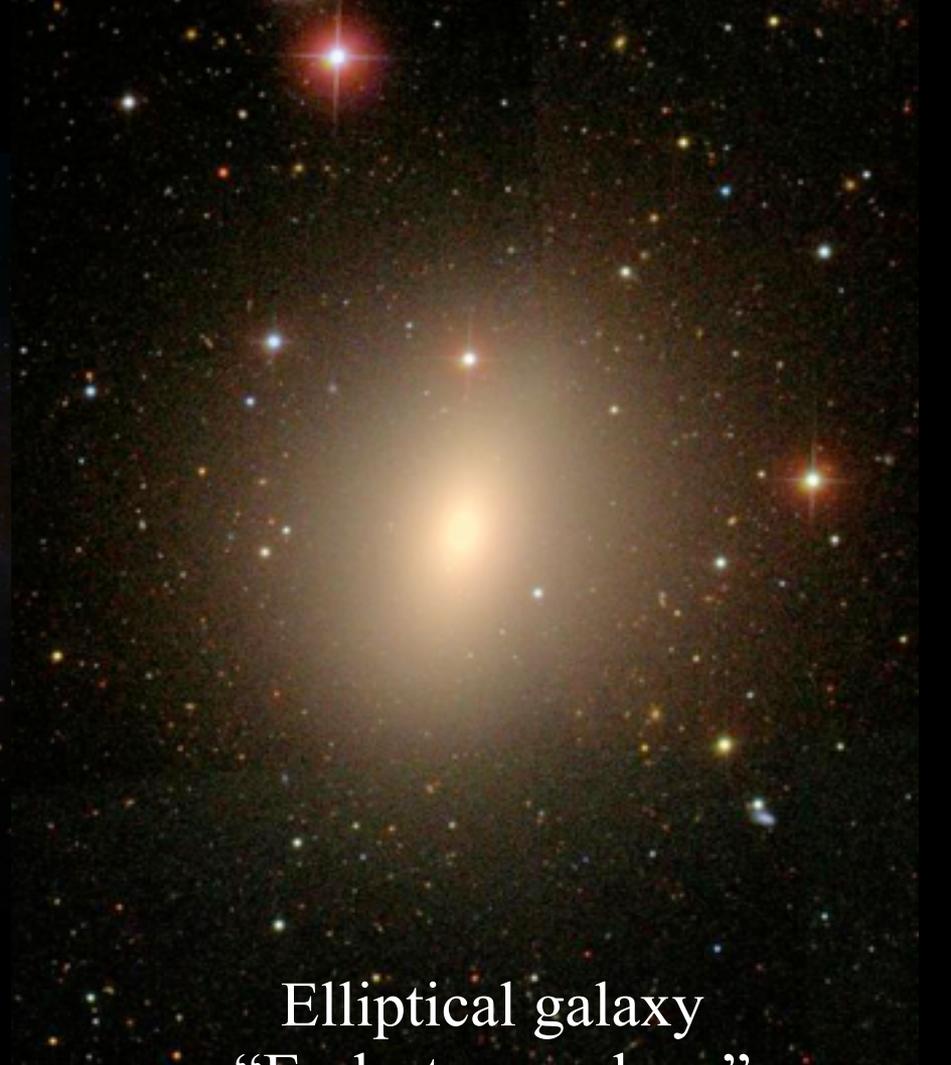
Stars



edge on view



Disk galaxy  
“Late-type galaxy”  
“Blue galaxy”  
“Star-forming galaxy”

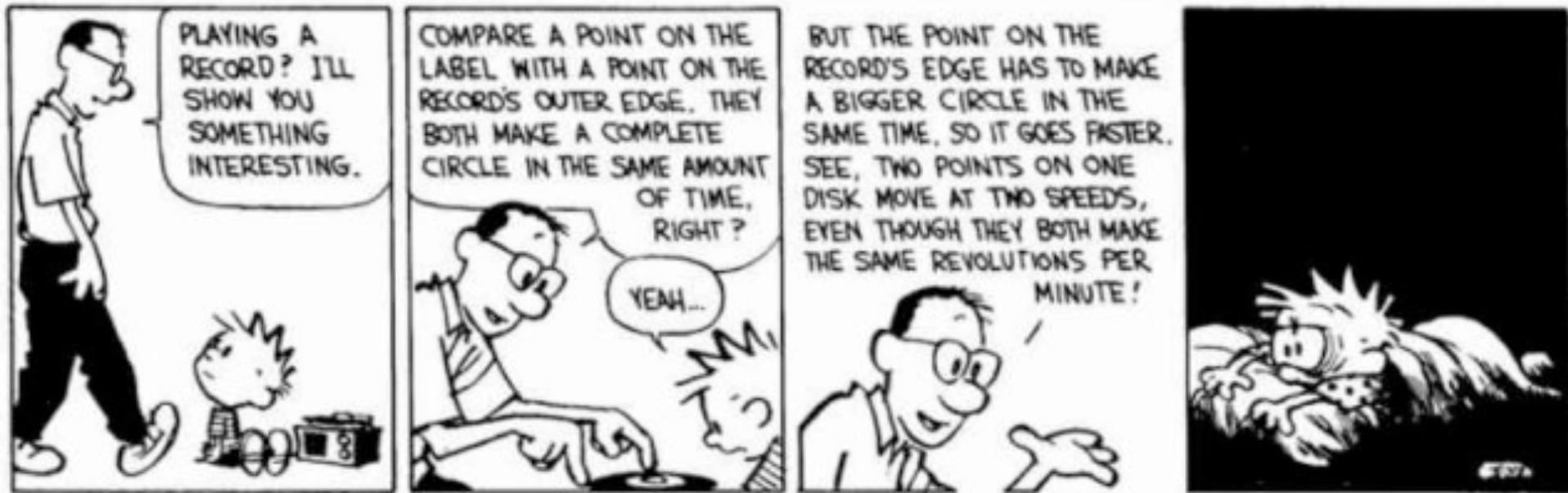


Elliptical galaxy  
“Early-type galaxy”  
“Red galaxy”  
“Red and dead”



Galaxy Cluster  
Massive Ellipticals Inside  
Less Massive Disks Outside

# Differential Rotation



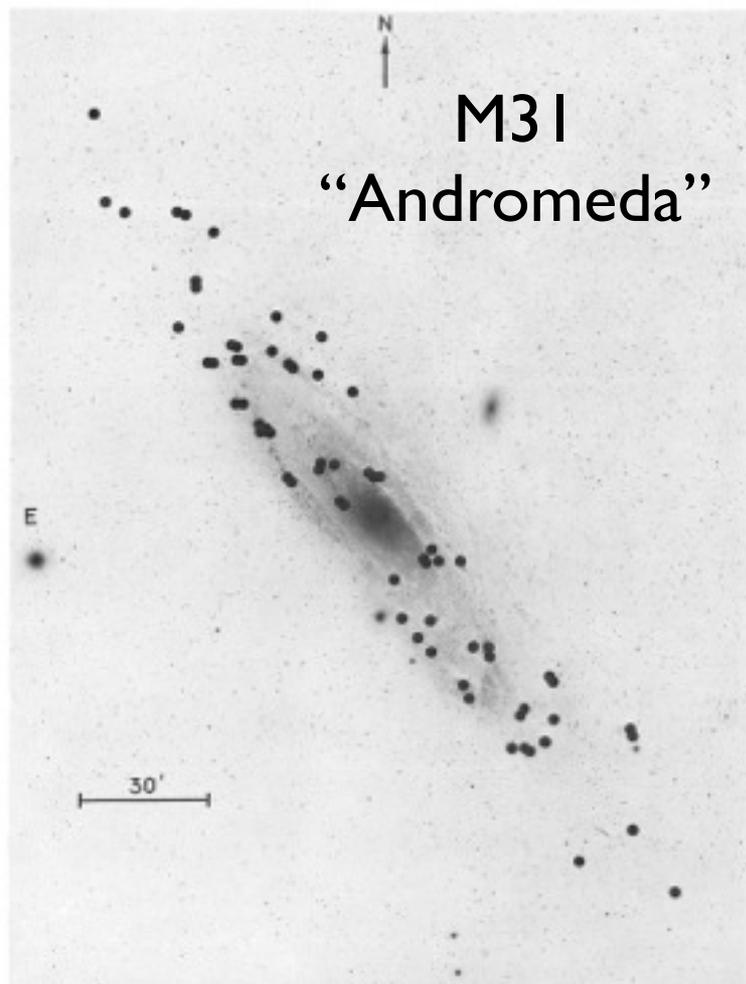


FIG. 1.—Identification chart for emission regions in M31 for which velocities have been obtained. Palomar 48-inch Schmidt ultraviolet photograph, 103a0 plate + UG 1 filter, courtesy of Dr. S. van den Bergh.

RUBIN AND FORD (see page 380)

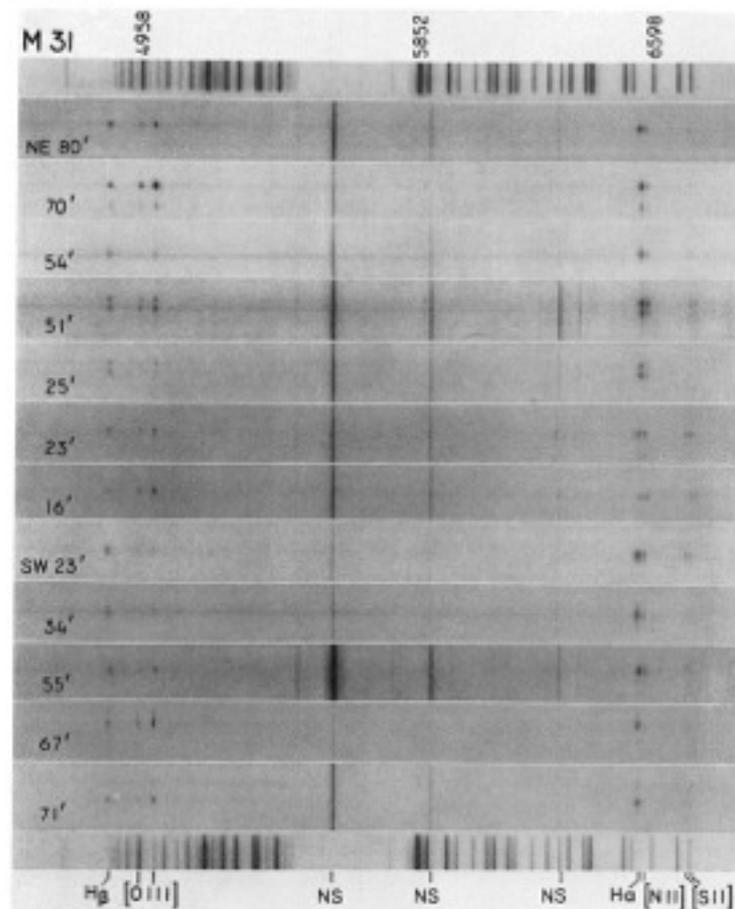


FIG. 2.—Representative spectra of emission regions in M31, arranged according to distance from center. Lines crossing the spectra are night sky lines. Comparison spectrum is Ne 4-Fe, original dispersion  $135 \text{ \AA mm}^{-1} = 6 \text{ km sec}^{-1} \mu^{-1}$  at H $\alpha$ . Principal emission lines are indicated.

RUBIN AND FORD (see page 381)

Many of the brighter emission regions in M31 are large, complex associations of emission regions and OB stars. However, there are also completely isolated Strömgren spheres. In an effort to obtain spectra from all parts of M31, we have generally observed no more than two regions within a single association, and have observed some of the isolated regions. In Table 1, column (9), we have classified the observed regions according to their membership in the M31 OB associations as defined by van den Bergh (1964). Fourteen of the regions we have observed are not in OB associations; the remaining fifty-three regions are distributed in thirty-four associations. Seventeen associations have velocities from two or more regions; for all association members the average deviation from the mean association velocity is  $10 \text{ km sec}^{-1}$ , or just equal to the average internal error in the radial velocity of an emission region. We conclude that the line-of-sight component of the velocity dispersion within an association is less than  $\pm 10 \text{ km sec}^{-1}$ , or less than  $15 \text{ km sec}^{-1}$  projected to a circular-velocity component in the plane of M31.

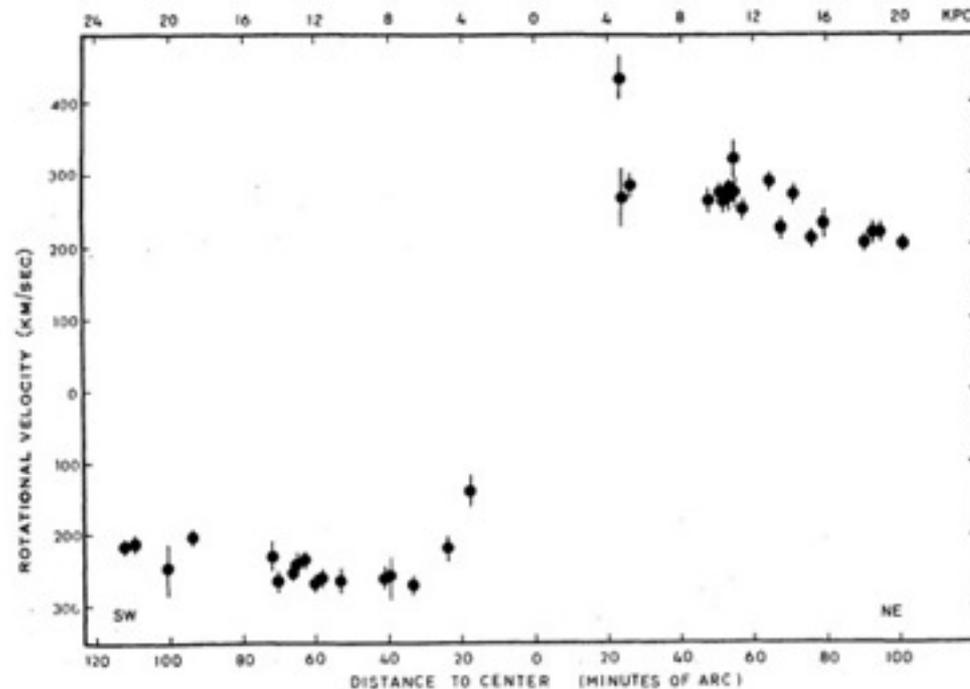


FIG. 4.—Mean rotation velocities for thirty-four OB associations plus three outer regions in M31, as a function of distance from the center. Error bars indicate average error of the velocity of the association, calculated from the velocities of individual emission regions in the association, or from the observational errors, whichever are larger.

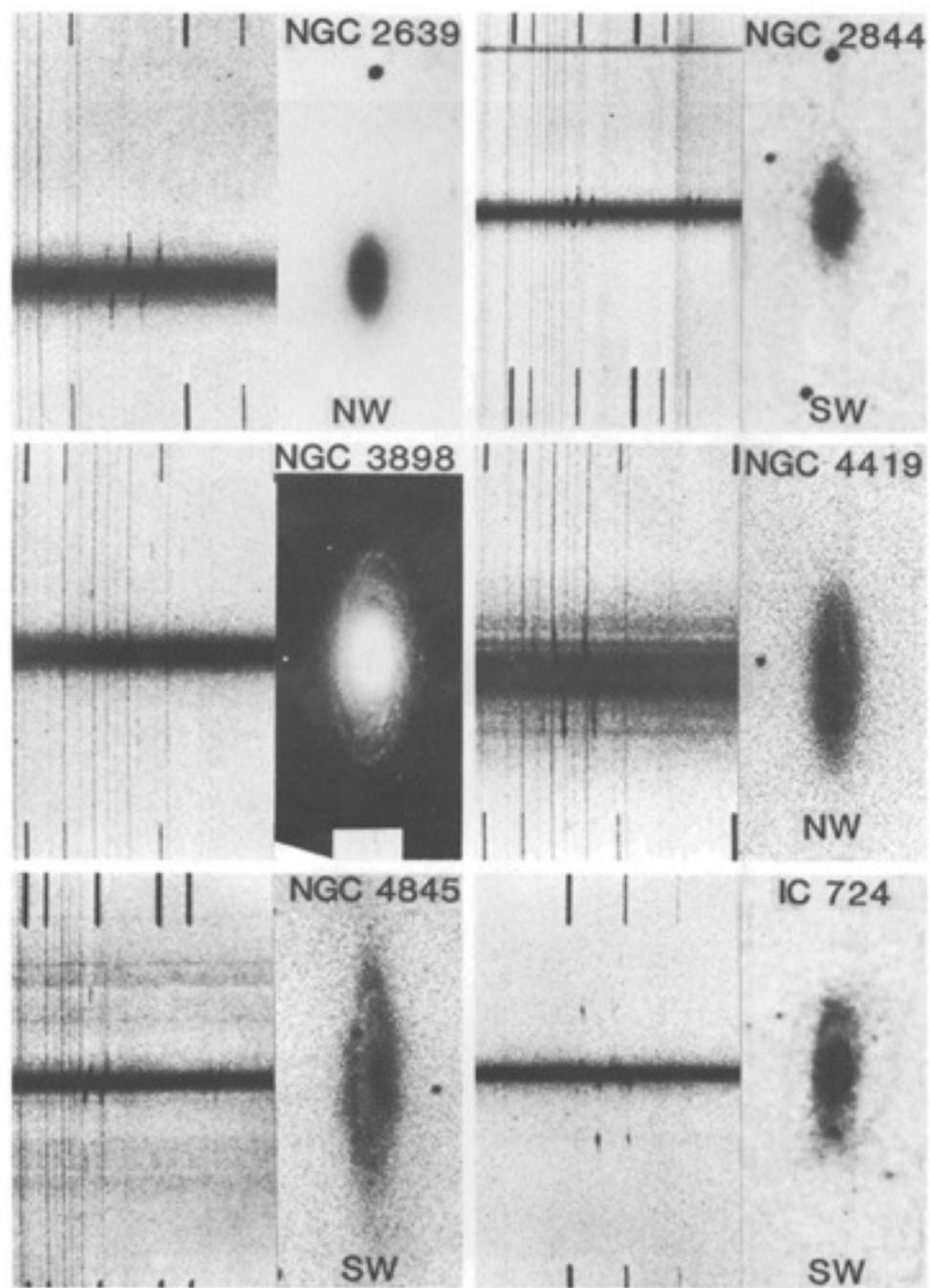


FIG. 1a

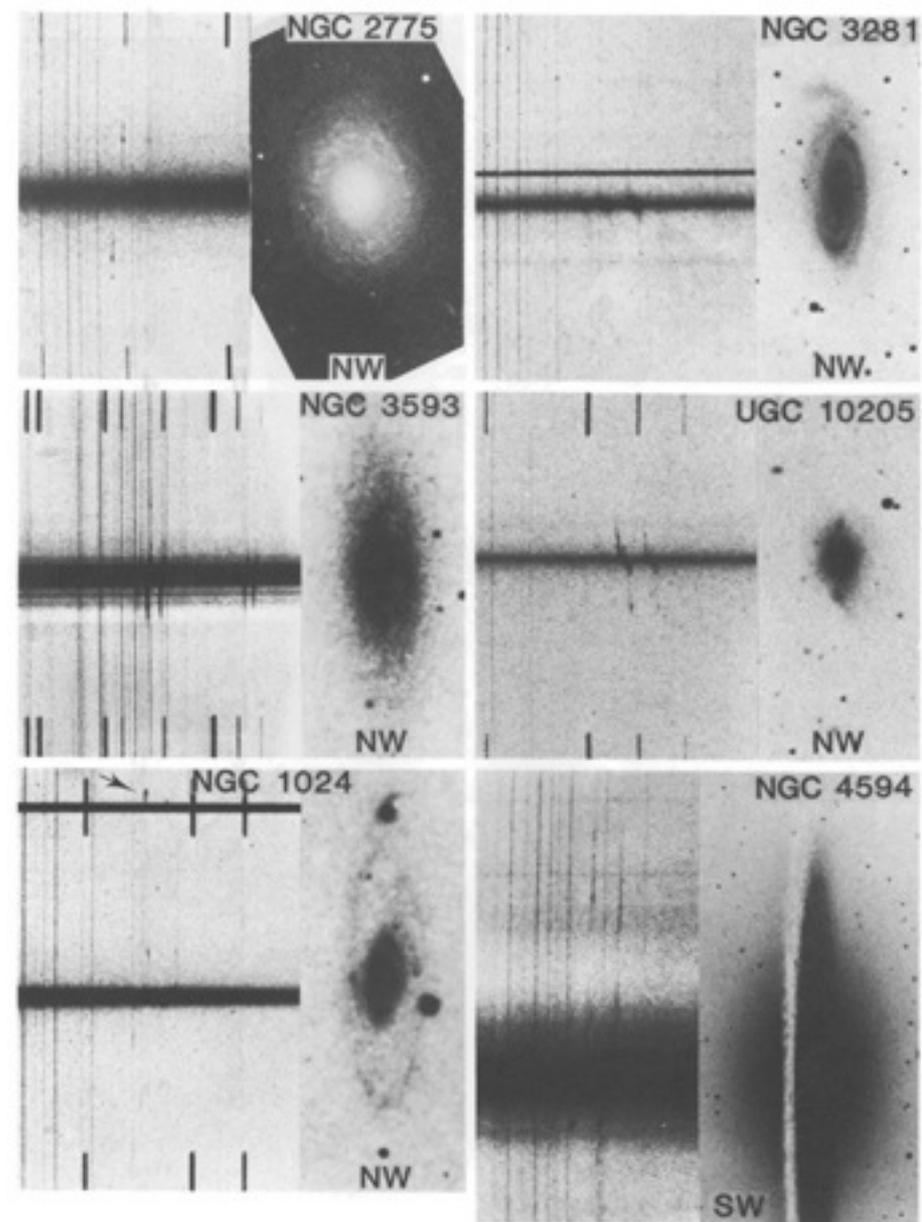


FIG. 1b

# Rotation Curves of Disk Galaxies

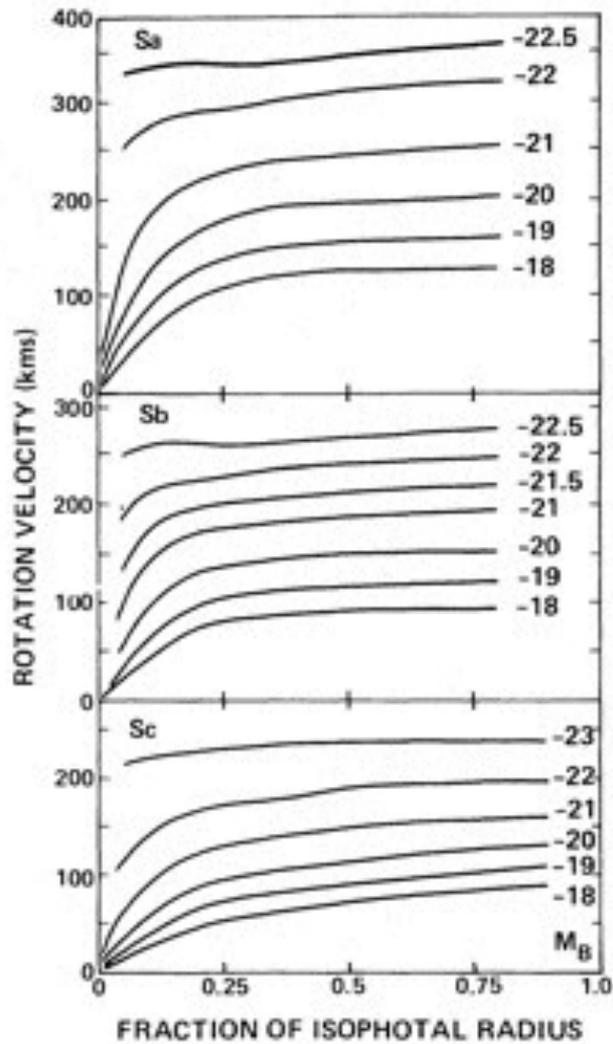
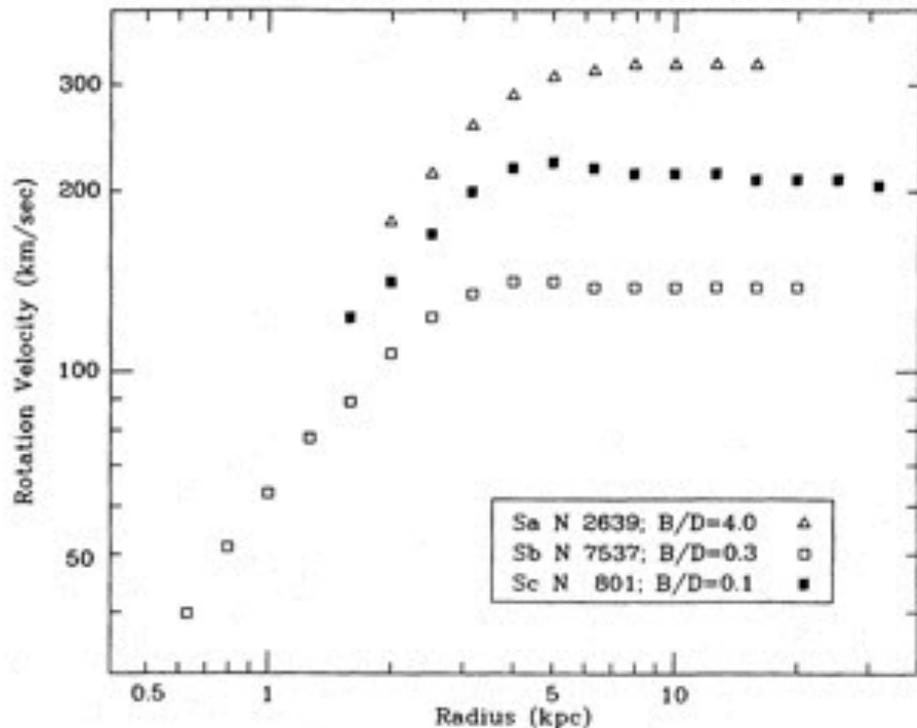


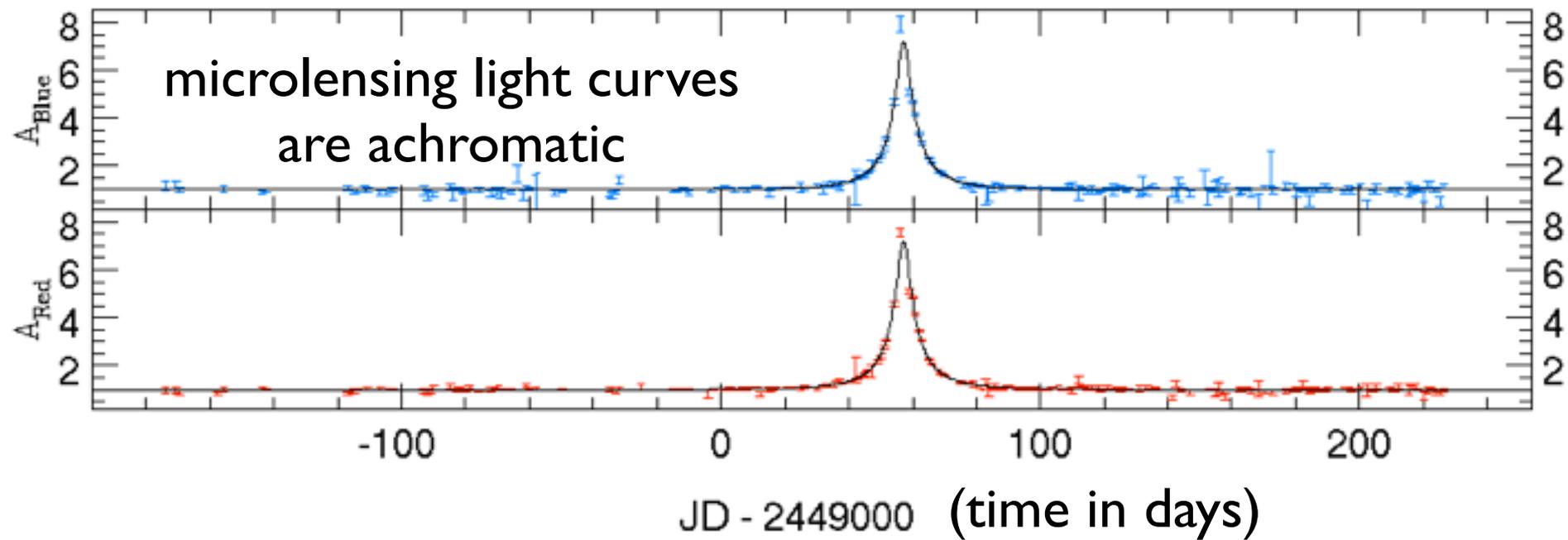
FIG. 4.—Synthetic rotation curves showing average smoothed rotation velocity as a function of fraction of isophotal radius,  $R_{25}$ , for (top) Sa galaxies of successive luminosities, (middle) Sb galaxies, and (bottom) Sc galaxies. The procedure for synthesizing these curves is described in Appendix B.

optical galaxy. As we discuss below in connection with  $M/L$  ratios, the nonluminous component must contribute to the overall gravitational potential at all radii. These observations suggest that the forms of the distributions of mass within the more extended halos are similar except for radial and/or density scale factors, independent of the morphology of the optical galaxy.

## IV. COMPARATIVE INTEGRATED PROPERTIES OF THE Sa, Sb, AND Sc PROGRAM GALAXIES

In this section, we compare integral properties of Sa, Sb, and Sc galaxies which we have studied. Rotational velocities for the Sa, Sb (Paper II), and Sc (Papers I and III) galaxies are tabulated in Appendix A.





Microlensing surveys stare at  
dense star fields

Large Magellanic Cloud  
(LMC = satellite galaxy of Milky Way)

Galactic Bulge  
(of our Milky Way)

# Dark Matter Candidates

Alternative to WIMP → MACHO

MACHO = MAssive Compact Halo Object

Try to find MACHOs using gravitational microlensing

- I. MACHO survey of Large and Small Magellanic Clouds
- II. OGLE = Optical Gravitational Lensing Experiment
- III. EROS = Experience de Recherche des Objets Sombres
- IV. AGAPE = Andromeda Gravitational Amplification Pixel Experiment

# Dark Matter Candidates

WIMP = Weakly Interacting Massive Particle



WIMPs feel gravity

WIMPs might also feel the weak nuclear force

e.g., radioactive decay like  $^{14}\text{C}$  into  $^{14}\text{N}$  = conversion of neutron into proton

e.g., pp chain in Sun = conversion of proton to neutron

range  $\sim 10^{-15}$  cm (incredibly small!)

Are WIMPS “hot” or “cold”?

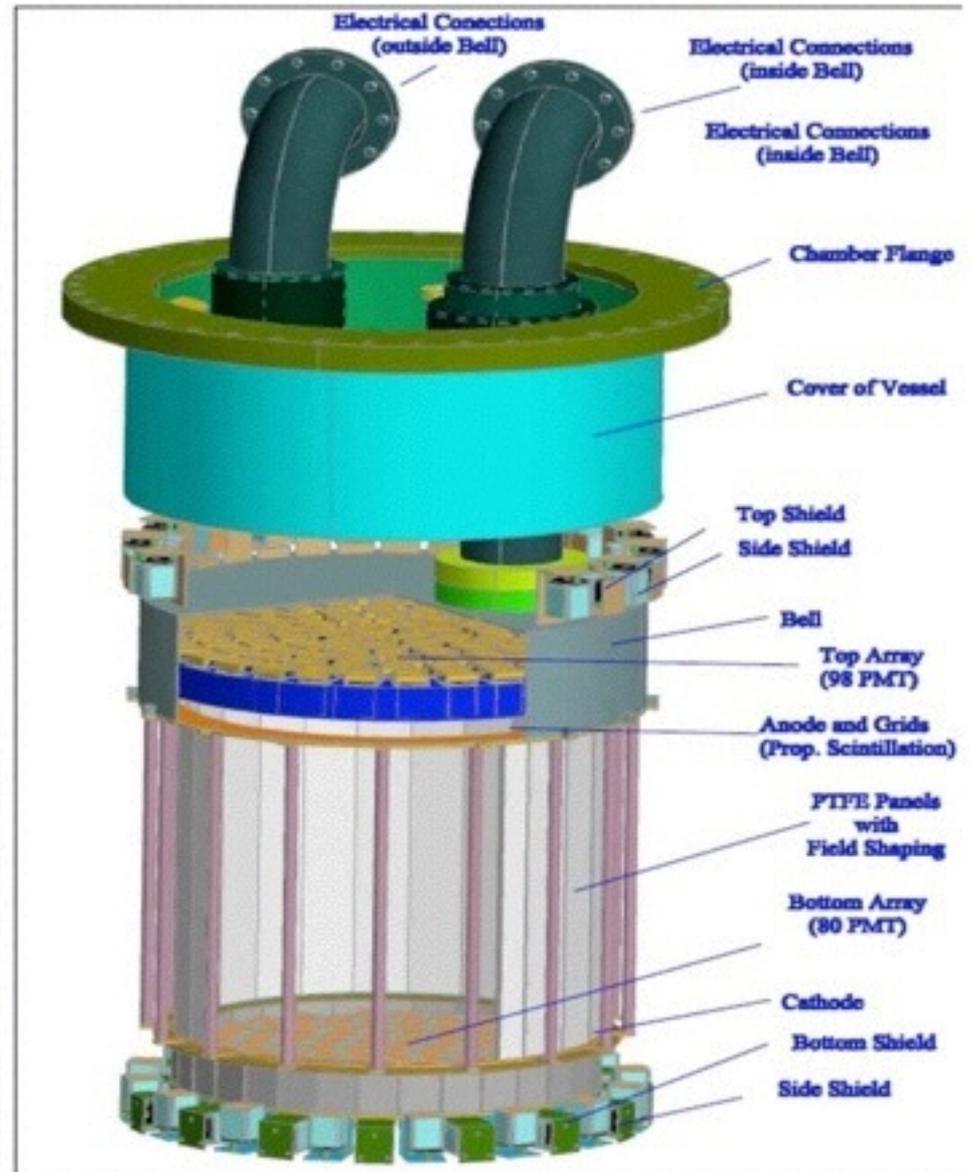


- “hot” = near light speed = “relativistic”
- low mass
- hot dark matter is known to exist  
e.g., neutrino: non-zero mass  $< 1$  eV
- but hot dark matter doesn't clump to form galaxies

- “cold” ( $v \ll c$ ) = “non-relativistic”
- high mass (maybe  $\sim 100$  GeV)
- many cold dark matter particles predicted, none (yet) detected
- nicely clumps to form galaxies

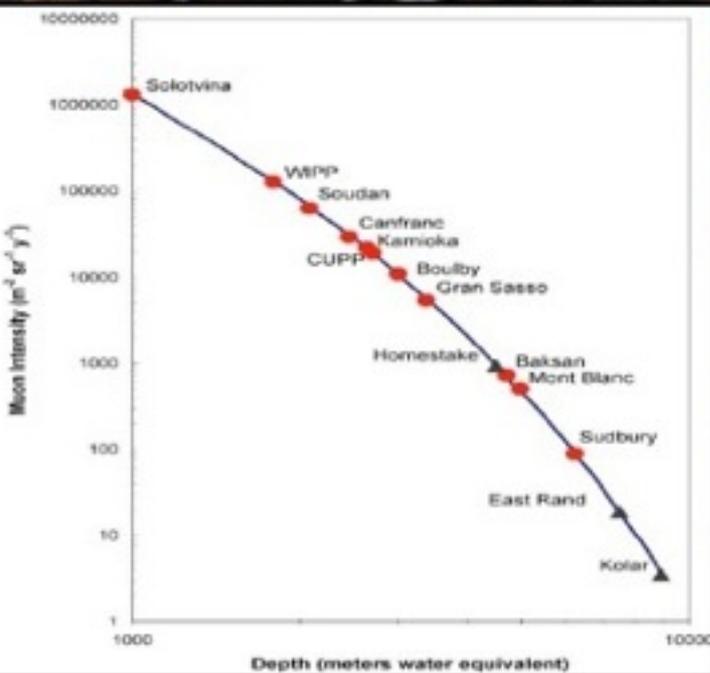
# XENON100: The TPC Assembly

- *x10 fiducial mass of XENON10*
- *x100 less back than XENON10*



# XENON100 @ LNGS

1.4 km of rock (3100 mwe) muon intensity is approximately  $1/(m^2 h)$



Installed underground in 2008

# World Wide Experiments for Dark Matter Direct Detection

YangYang  
KIMS

Homestake  
LUX

SNOLAB  
DEAP/CLEAN  
PICASSO

Boulby  
ZEPLIN  
DRIFT

Kamioka  
XMASS

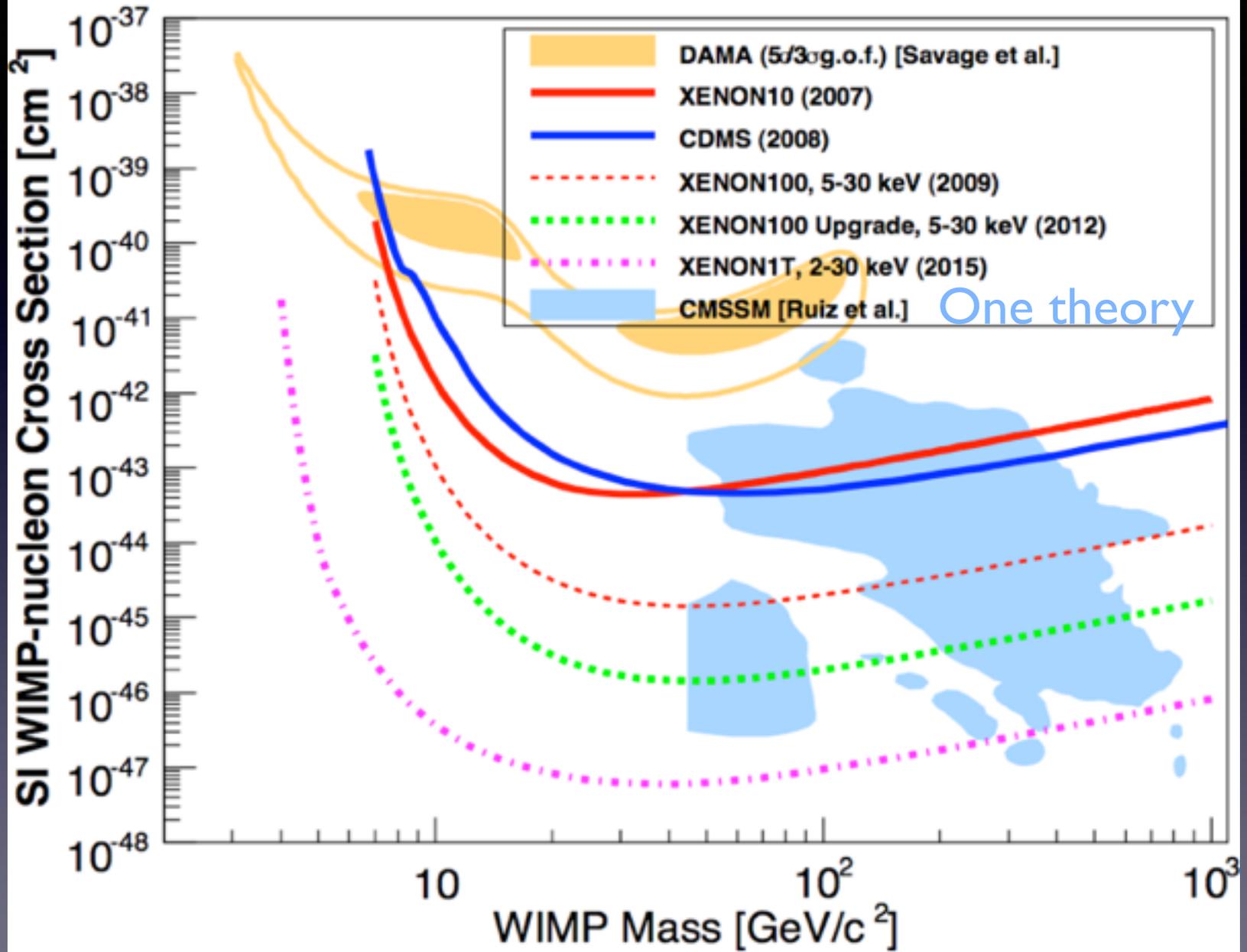
Soudan  
CDMS

Frejus/  
Modane  
EDELWEISS

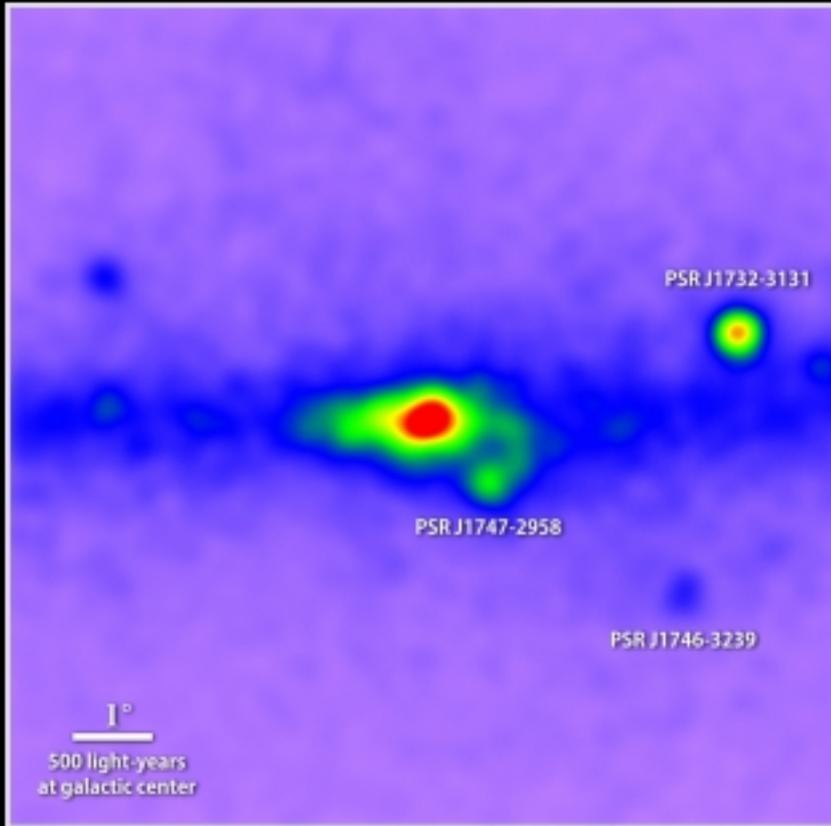
Canfranc  
ArDM  
ANAIS

Gran Sasso  
CRESST  
DAMA/LIBRA  
WARP  
XENON

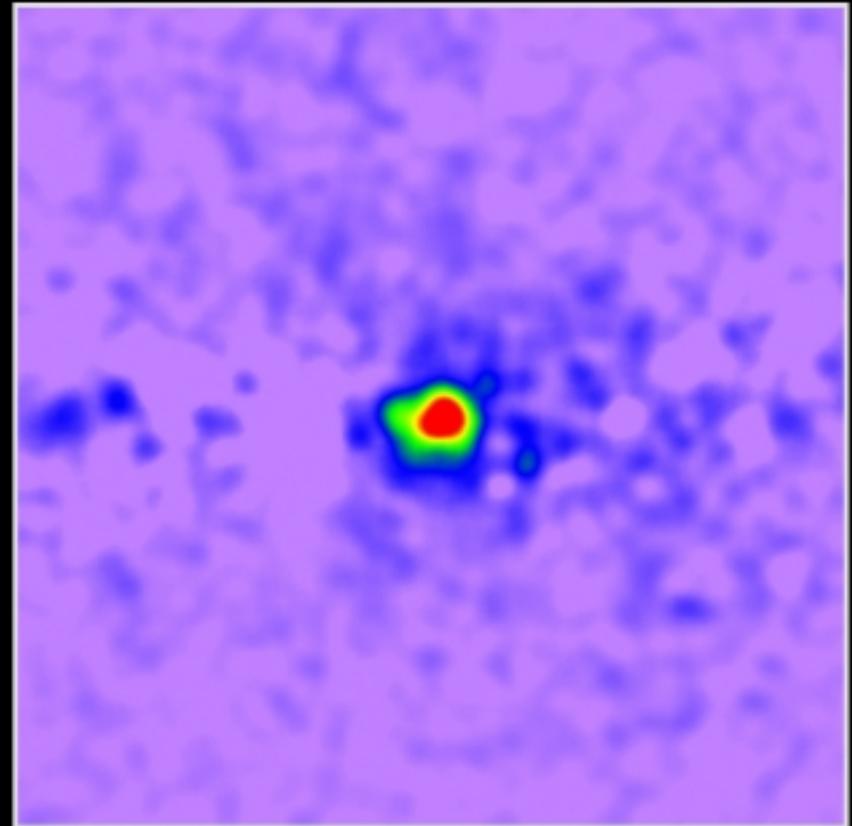




## Uncovering a gamma-ray excess at the galactic center



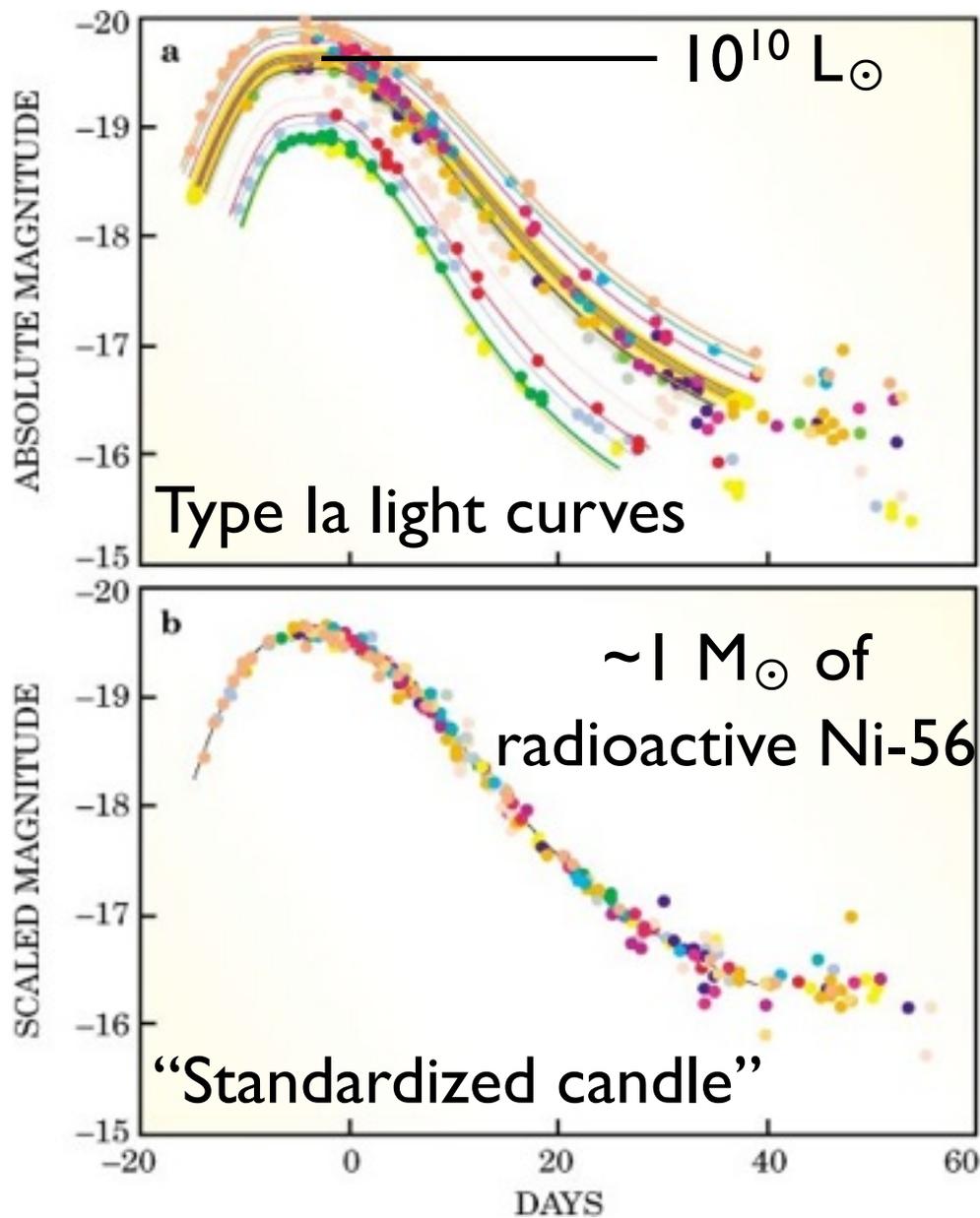
Unprocessed map of 1.0 to 3.16 GeV gamma rays



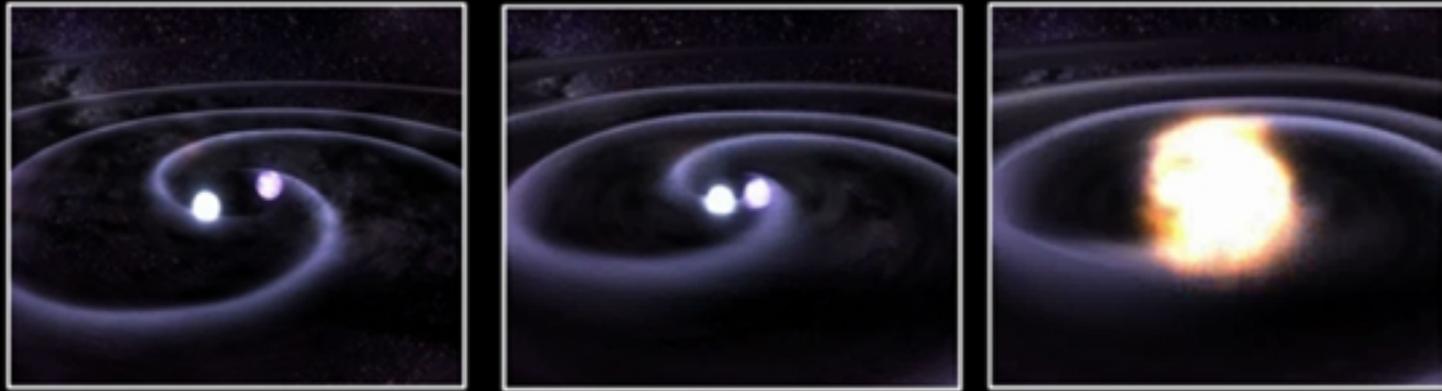
Known sources removed

Possible “dark matter self-annihilation signal”

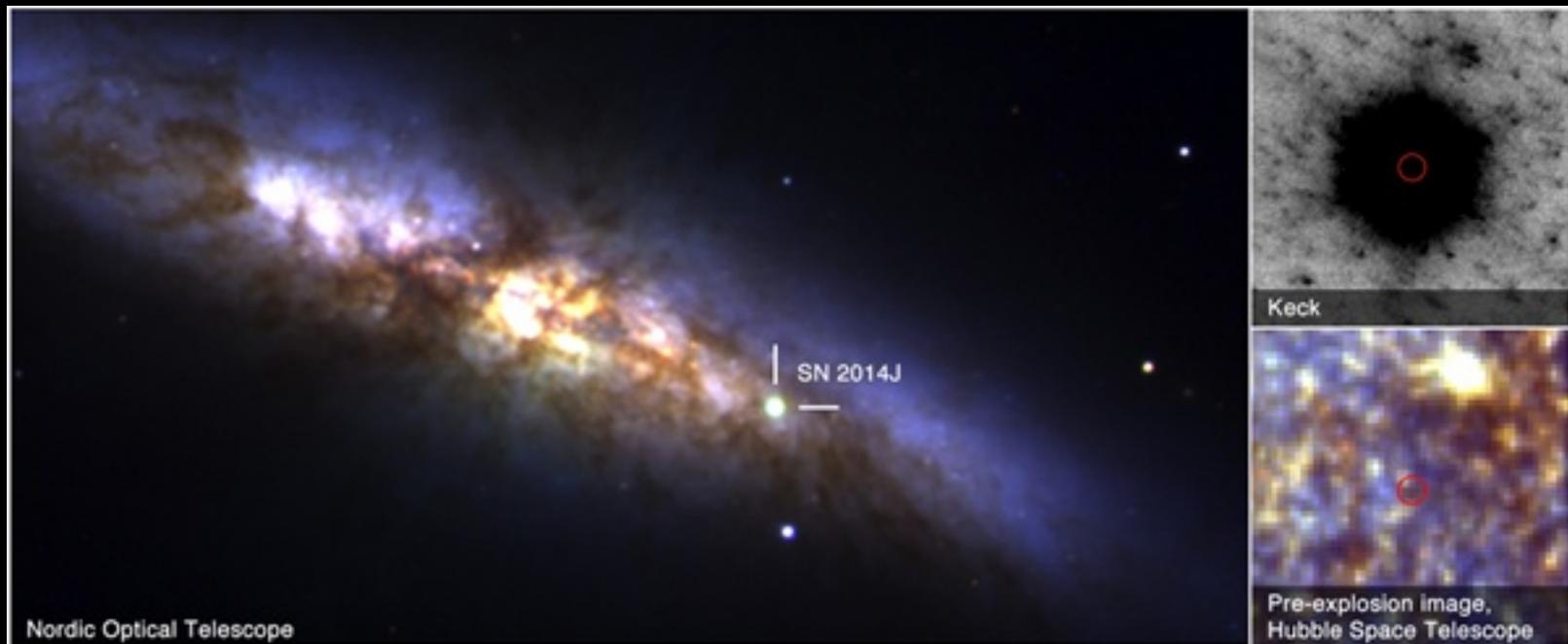
# Type Ia supernova as standard candle

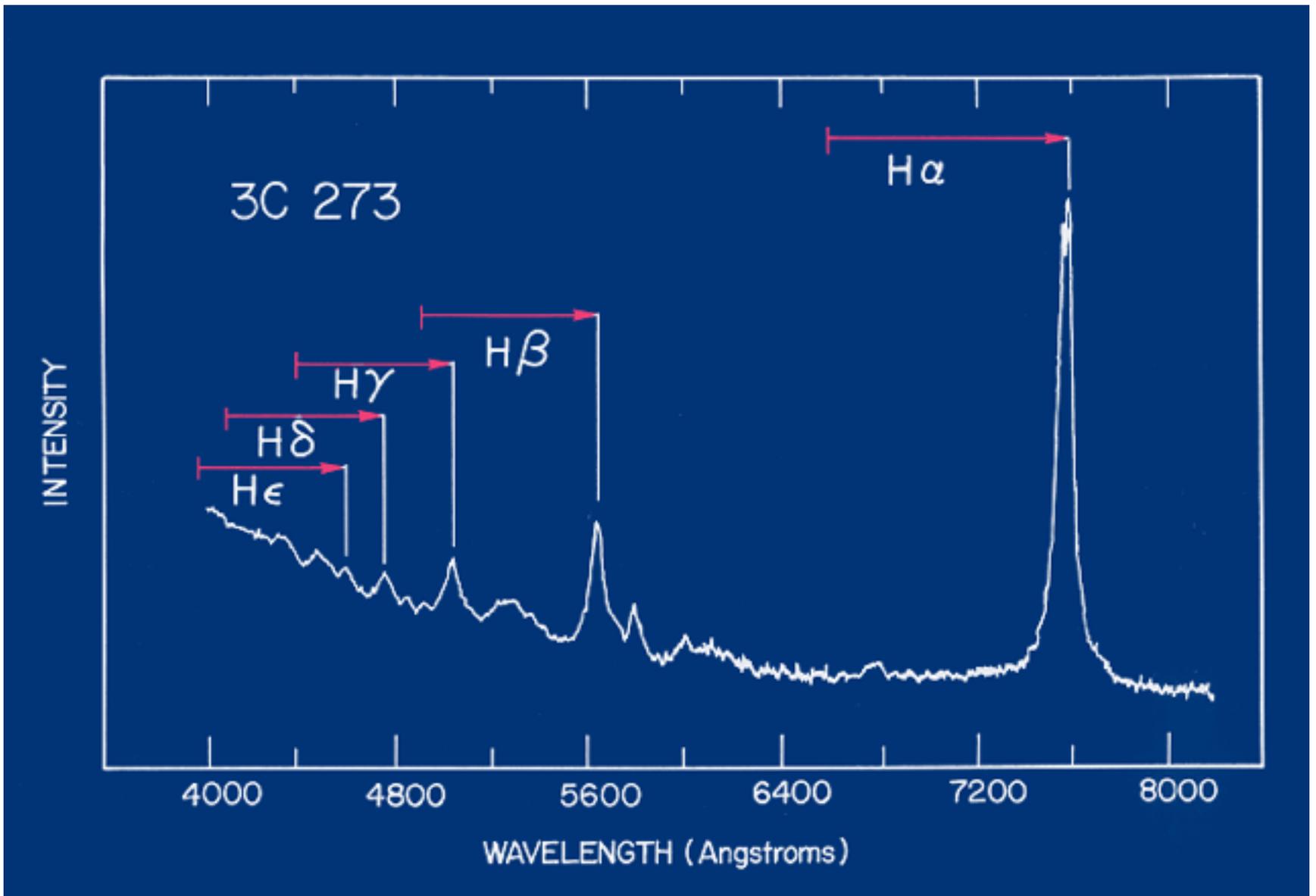


# Type Ia Supernova: “Double Degenerate” Progenitor



Two white dwarfs merging: C, O fuses to Si fuses to Ni

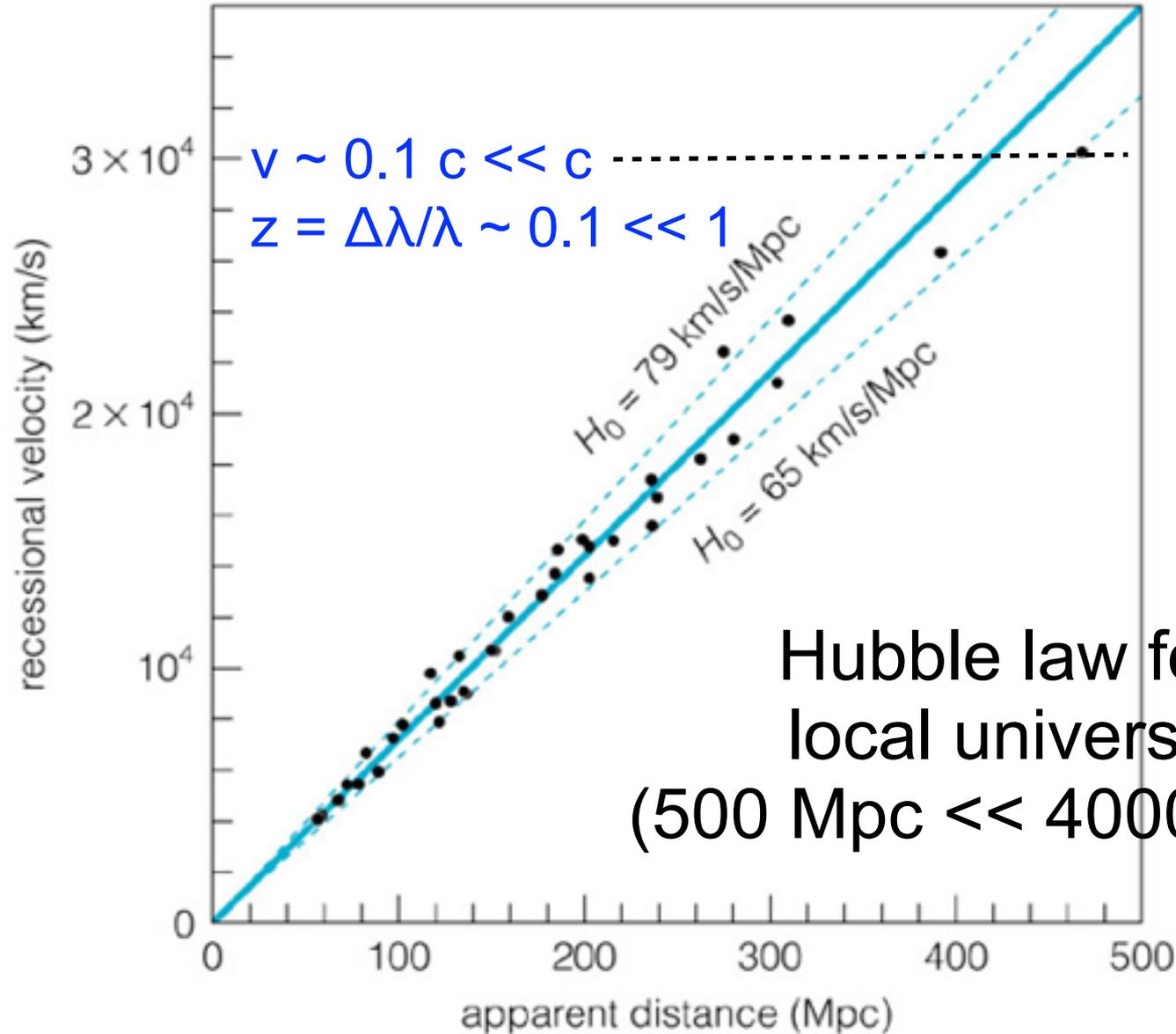




Redshift  $\Rightarrow$  galaxy is receding from us

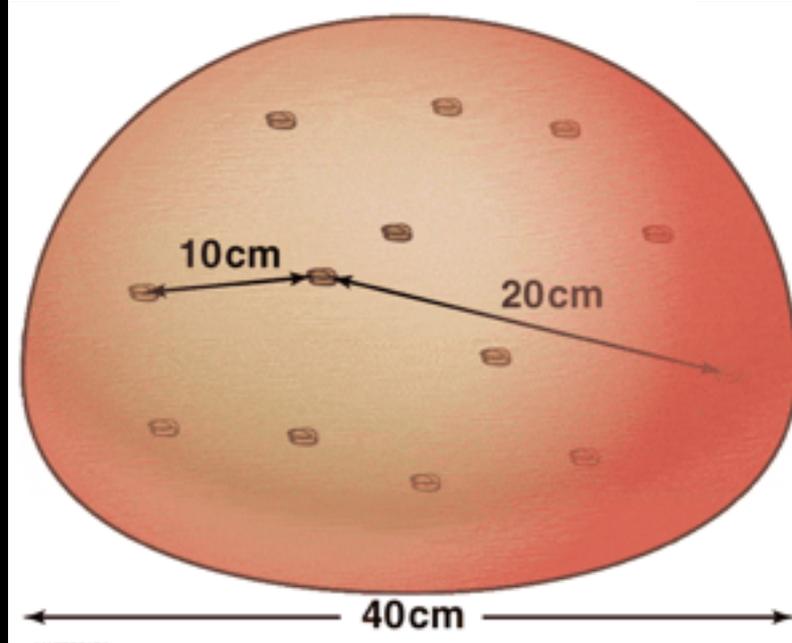
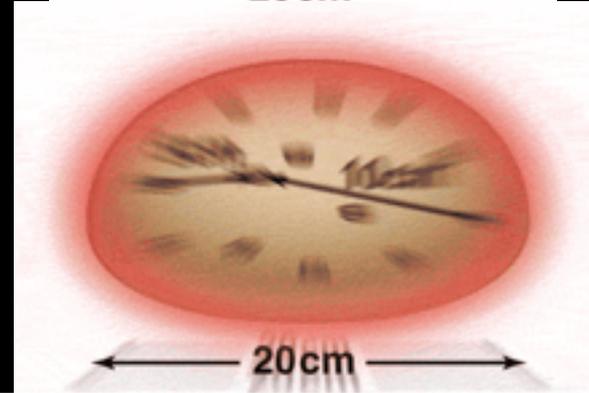
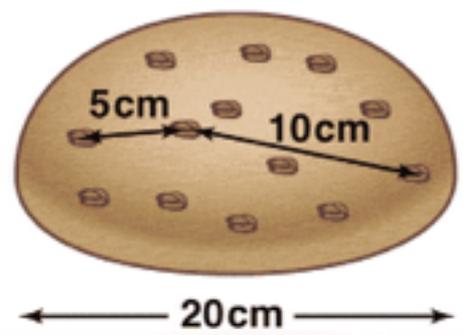
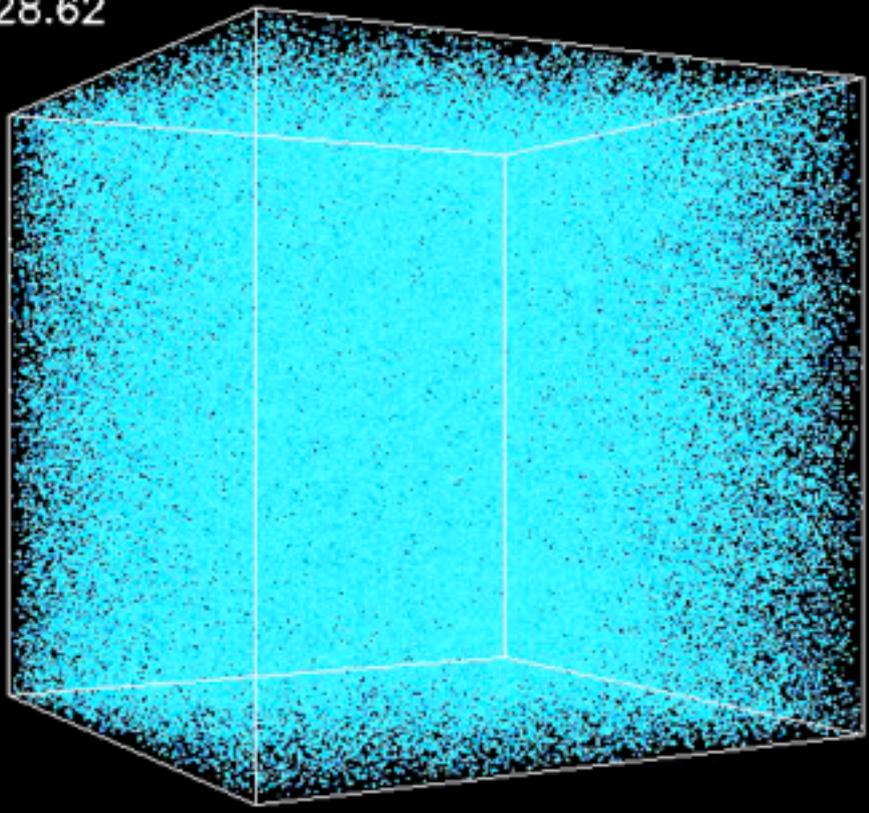


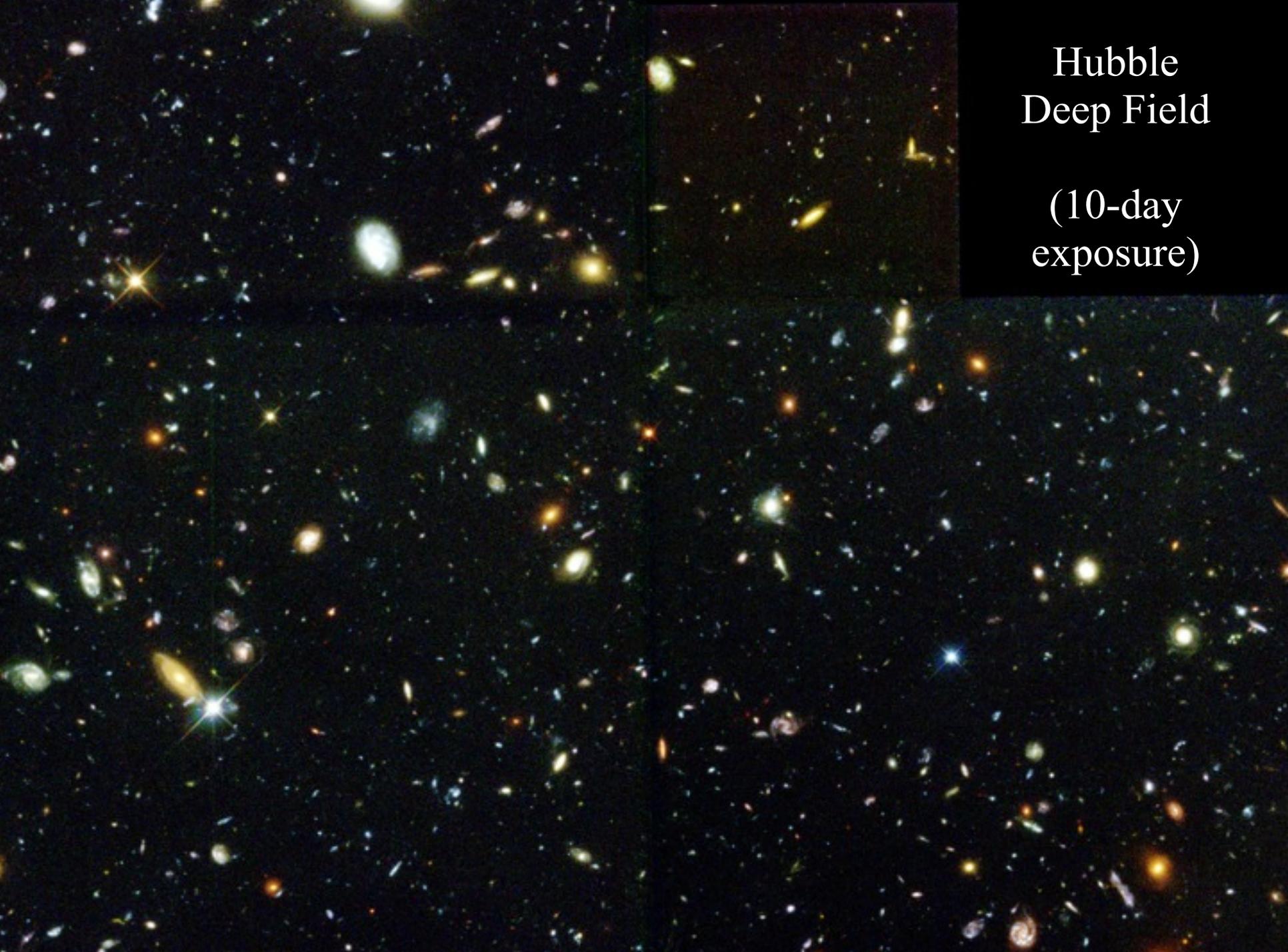
local universe = “low-redshift universe” = “low-z”





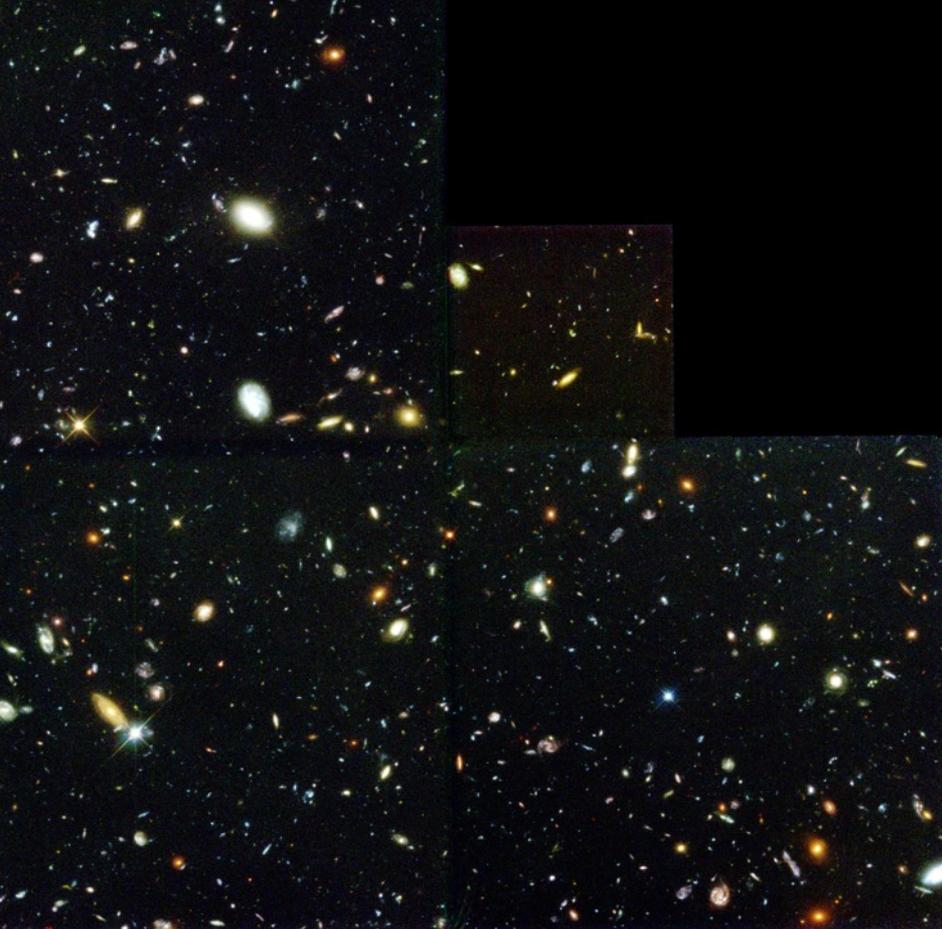
Z=28.62





Hubble  
Deep Field

(10-day  
exposure)



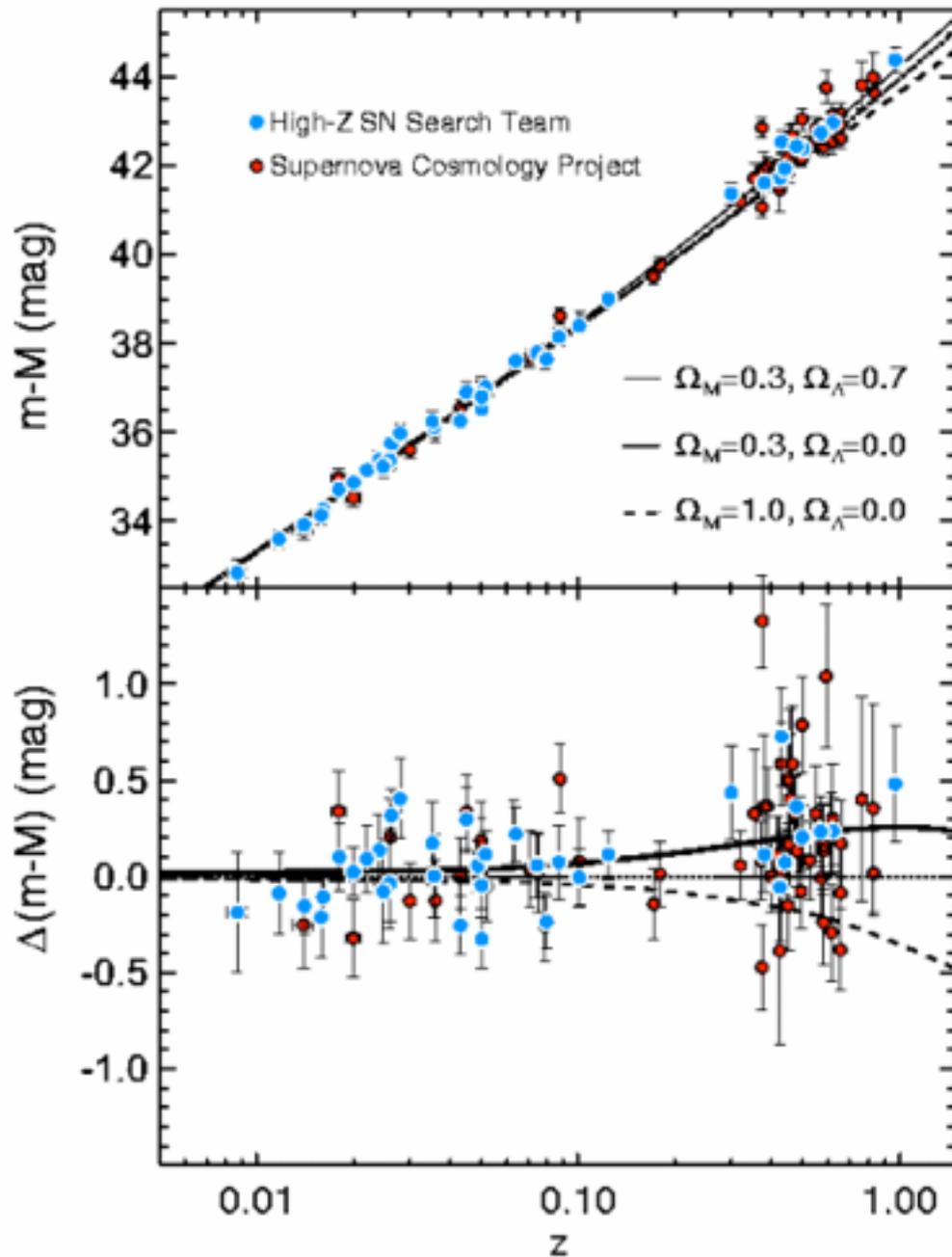
Hubble  
Deep Field

(10-day  
exposure)

Hubble  
eXtreme  
Deep Field

(23-day  
exposure)



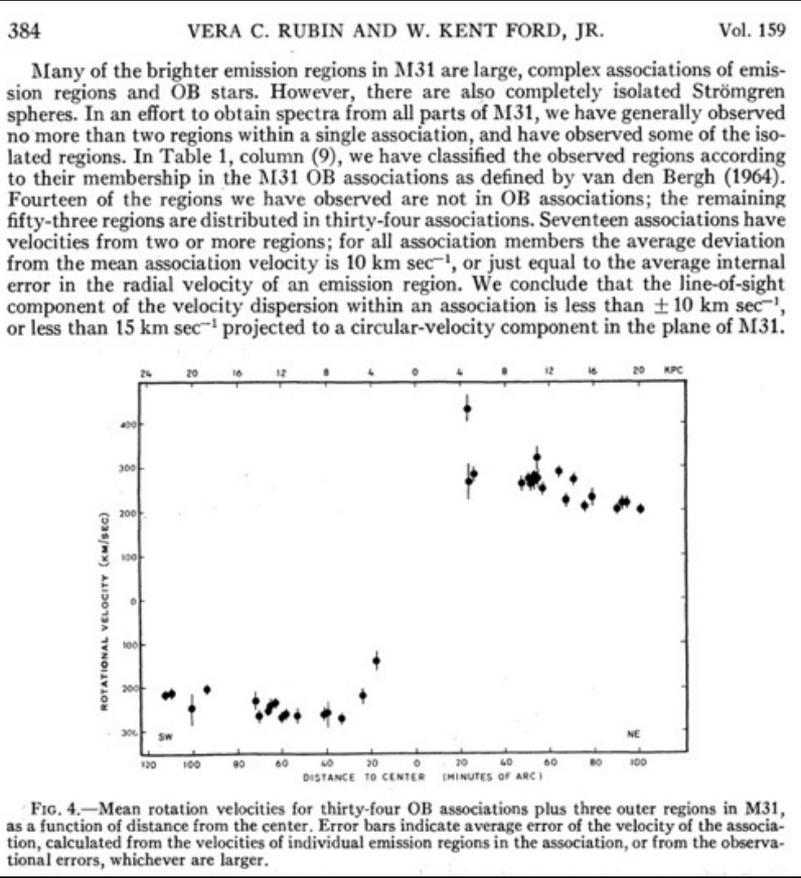


magnitude-redshift  
relation  
(a.k.a. generalized  
Hubble law)  
shows evidence  
for  
positively  
accelerating  
universe  
(old models were  
typically  
decelerating)

# Evidence for Dark Matter

## $\Omega_m = 0.26$

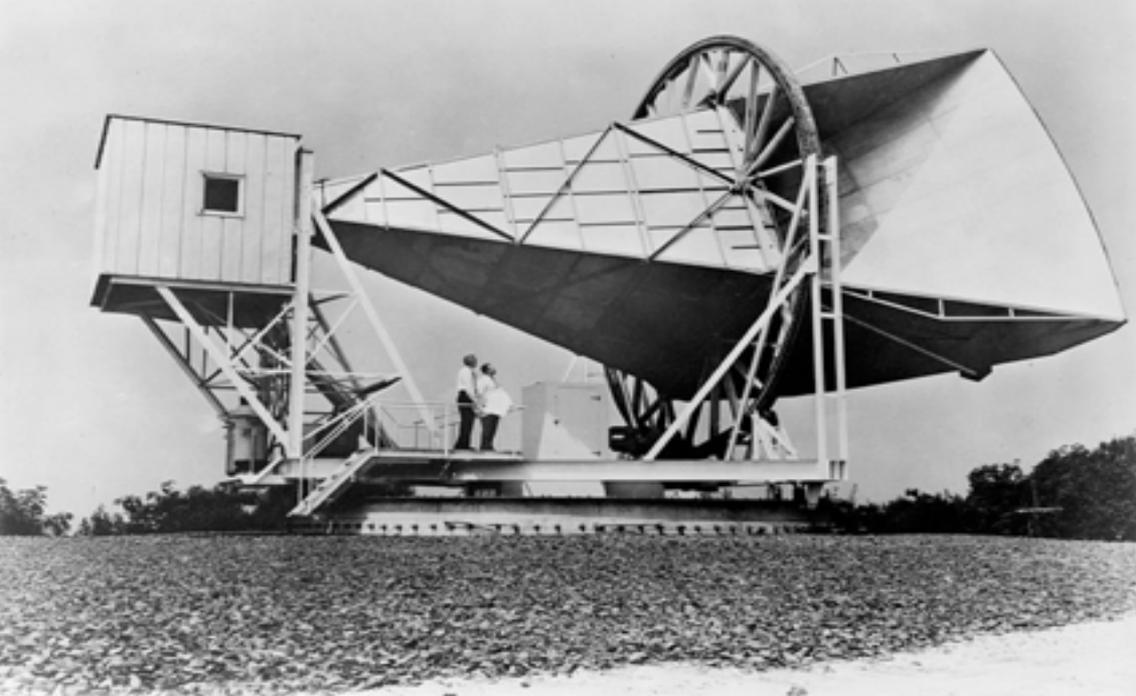
# Gravitational Lensing + Fast Galaxies (1000 km/s) in Galaxy Clusters



Hot Gas  
in Galaxy  
Haloes  
and  
Galaxy  
Clusters

Flat Rotation Curves  
of Individual Galaxies



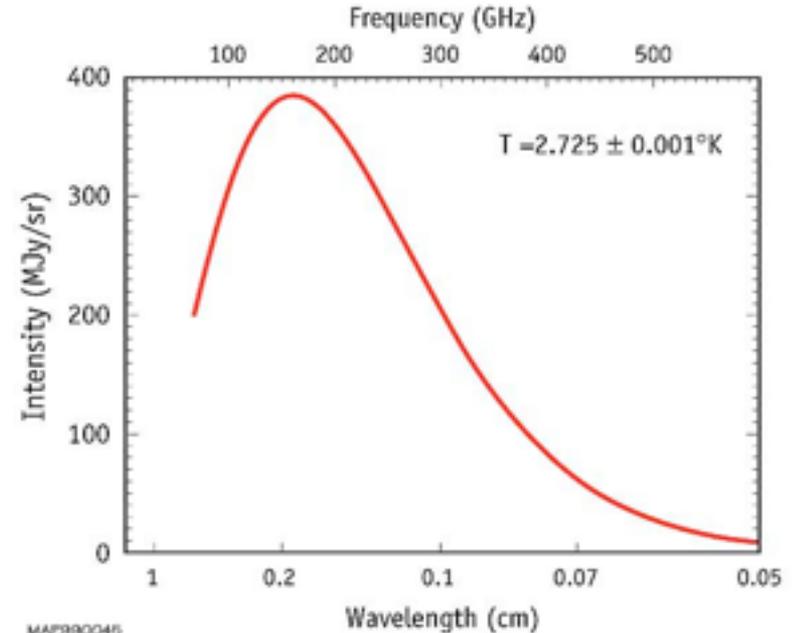


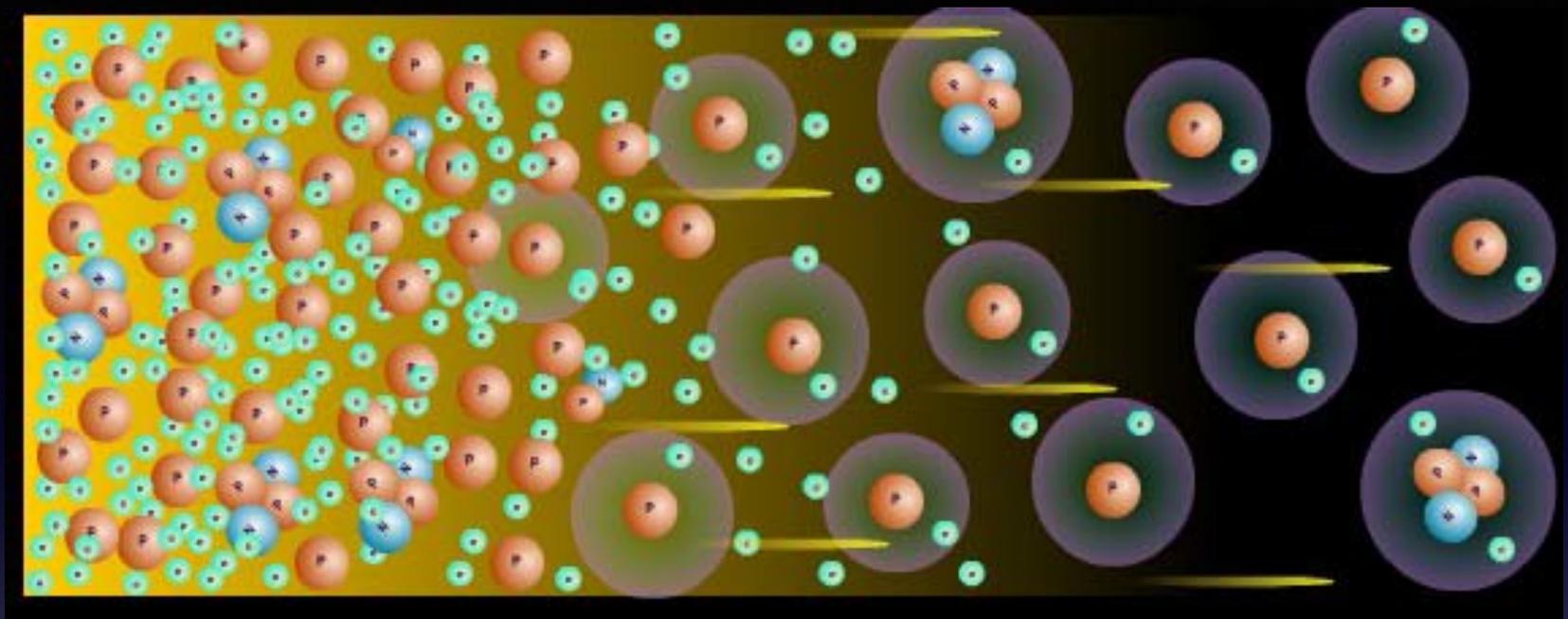
Holmdel  
Horn Antenna  
(Bell Labs)  
7.35 cm wavelength  
Penzias & Wilson  
1960s

### ISOTROPY OF THE COSMIC MICROWAVE BACKGROUND



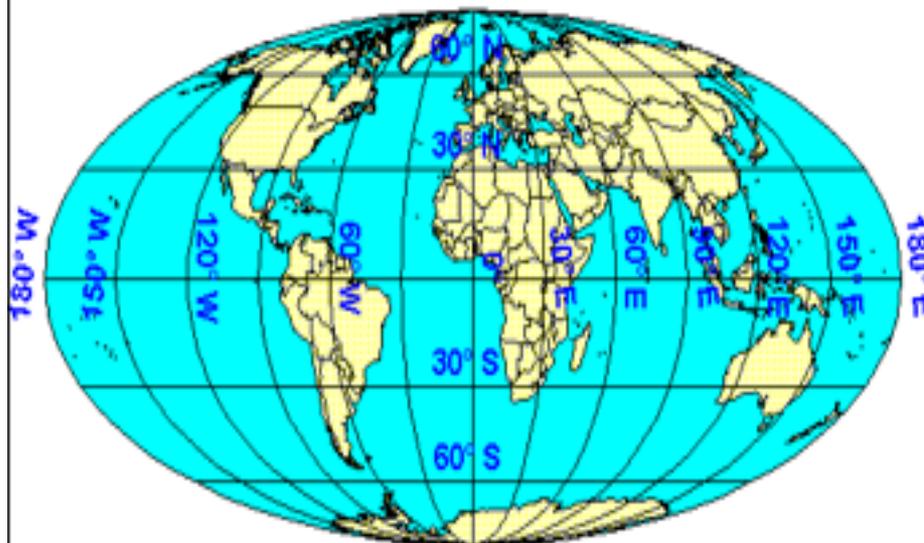
### SPECTRUM OF THE COSMIC MICROWAVE BACKGROUND





Peter H. Dana 9/20/94

## ISOTROPY OF THE COSMIC MICROWAVE BACKGROUND



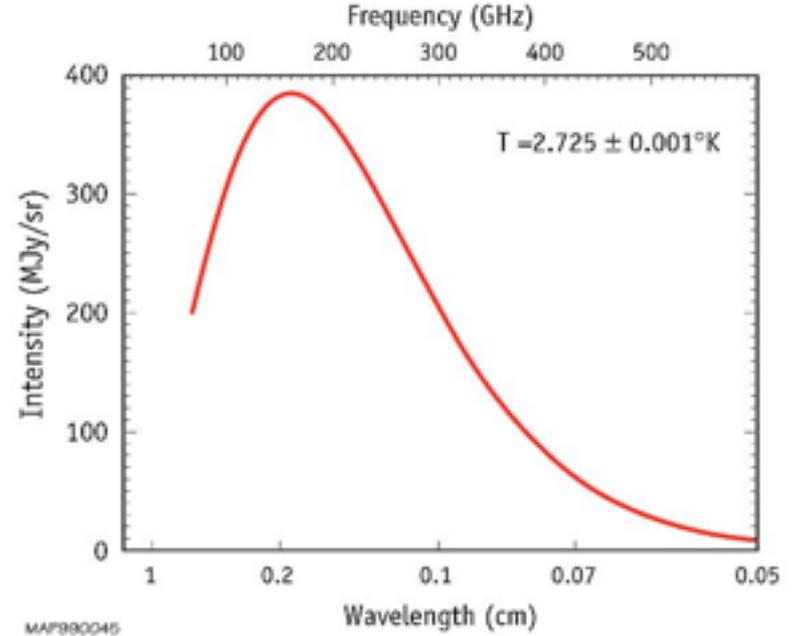
Mollweide Equal-Area

# ISOTROPY OF THE COSMIC MICROWAVE BACKGROUND



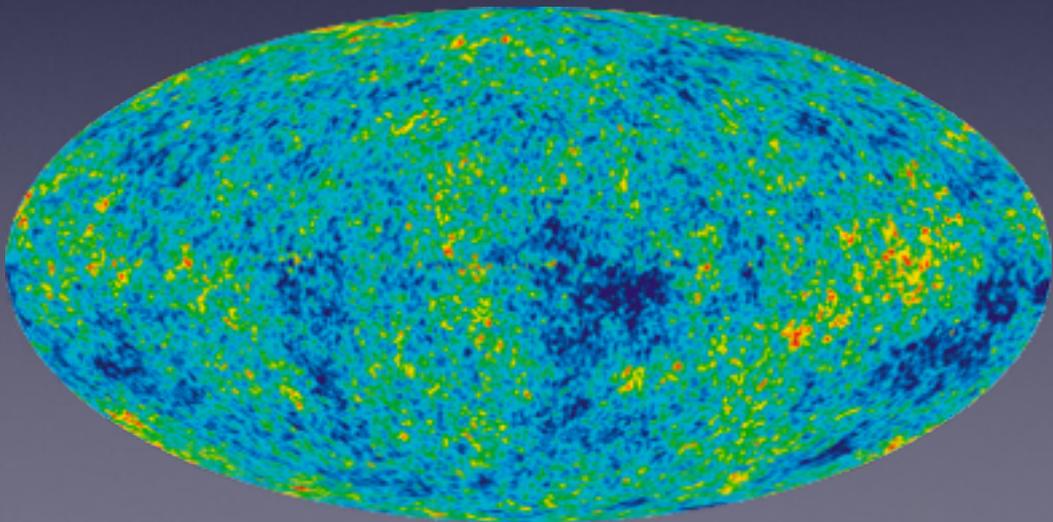
MAP990004

# SPECTRUM OF THE COSMIC MICROWAVE BACKGROUND



MAP990046

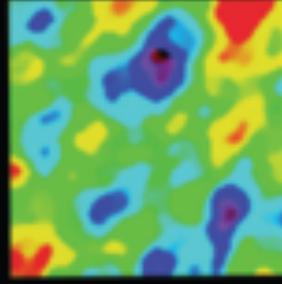
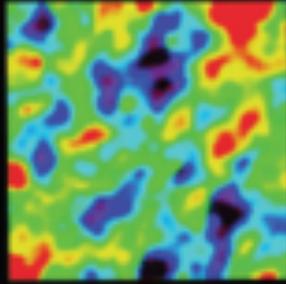
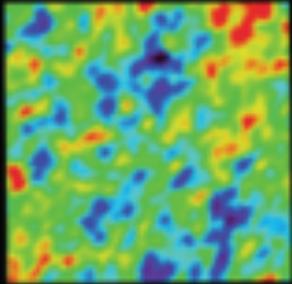
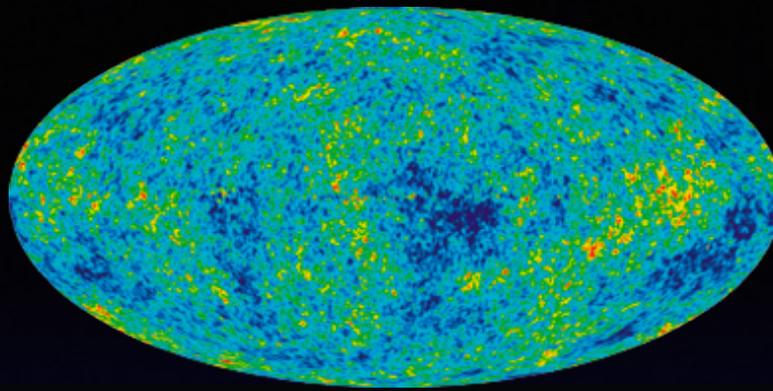
red = 2.726 K    blue = 2.724 K



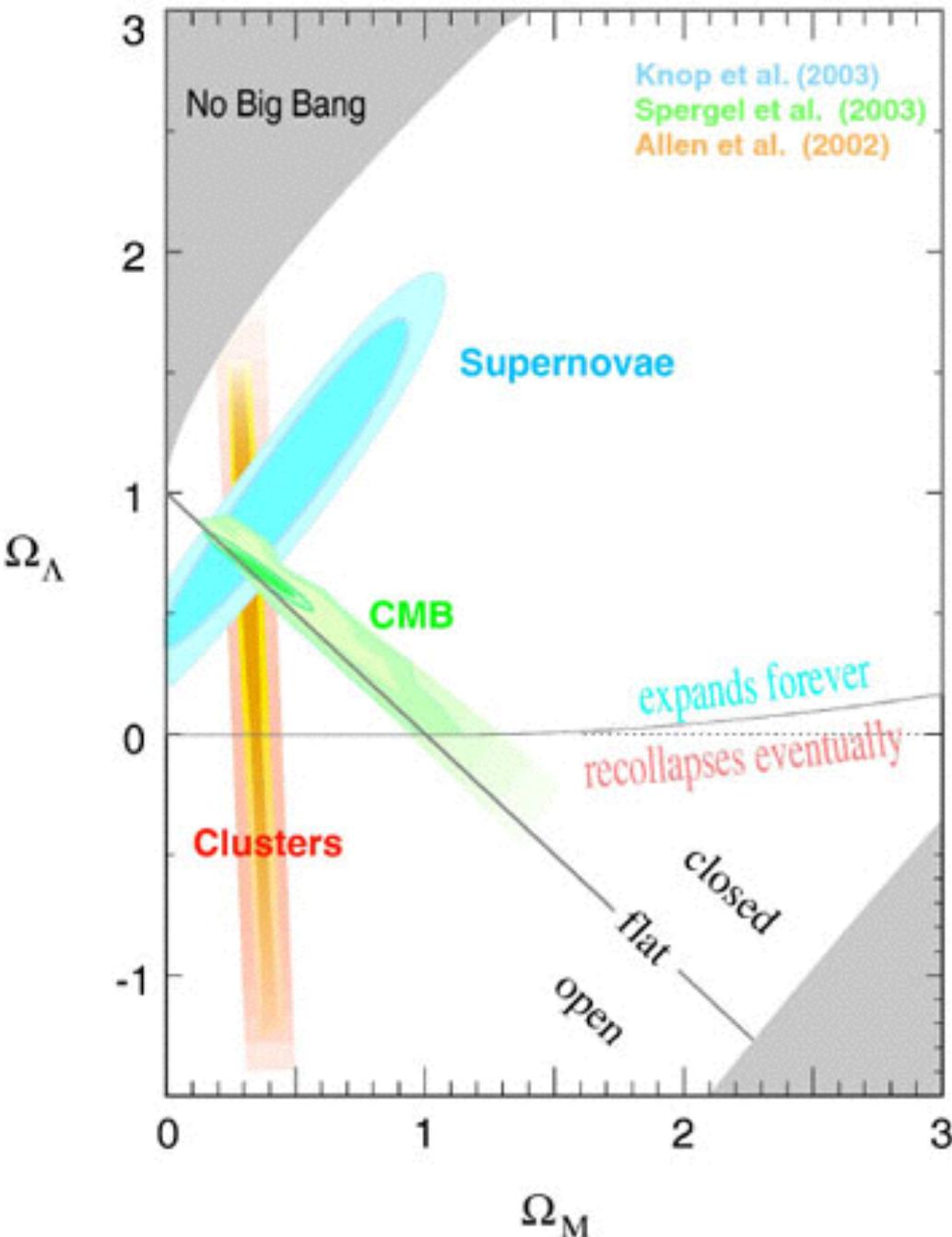
CMB temperature fluctuations  
= photons redshifted and blueshifted  
by **dark matter clumps and voids**  
at time of recombination

Angular size of dark matter clumps  $\sim 1^\circ$   
= **horizon size at time of recombination**

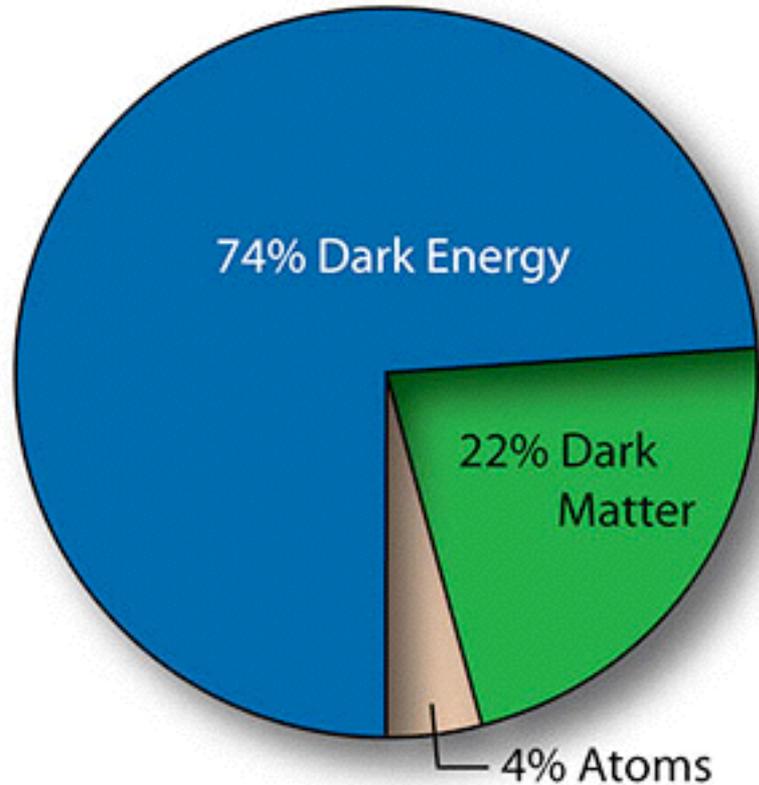
Clumps originate from **quantum fluctuations** in the  
inflation energy field,  
**grown to cosmic proportions by inflation**



*The physical size of the fluctuations is the horizon size at the last scattering surface.*



# The Consensus Model



REALIZATION I  
 The kind of supercooling can explain why the universe today is so incredibly flat and therefore why results the fine-tuning paradox pointed out by Bob Dicke in his Quinlan 200 lectures.

Let us first rederive the Dicke paradox. He relies on the empirical fact that the deceleration parameter today  $q_0$  is of order 1.

$$q_0 = -\ddot{R} \frac{R}{\dot{R}^2}$$

Use the age of the universe

$$\begin{aligned} \dot{R} &= -4\pi G \rho R^2 \\ \dot{R}^2 + k &= \frac{8\pi G}{3} \rho R^2 \end{aligned}$$

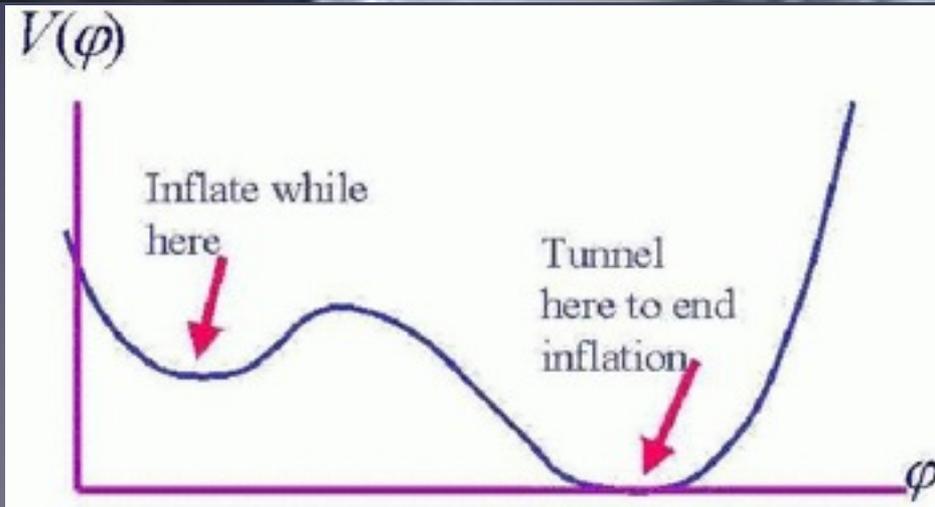
~~$$q_0 = \frac{1}{2} \left( 1 - \frac{2\pi G \rho R^2}{\dot{R}^2} \right)$$

$$= \frac{1}{2} \left( 1 - \frac{3k}{2\dot{R}^2} \right)$$~~

## Permanent collection of the Adler Museum (Chicago)

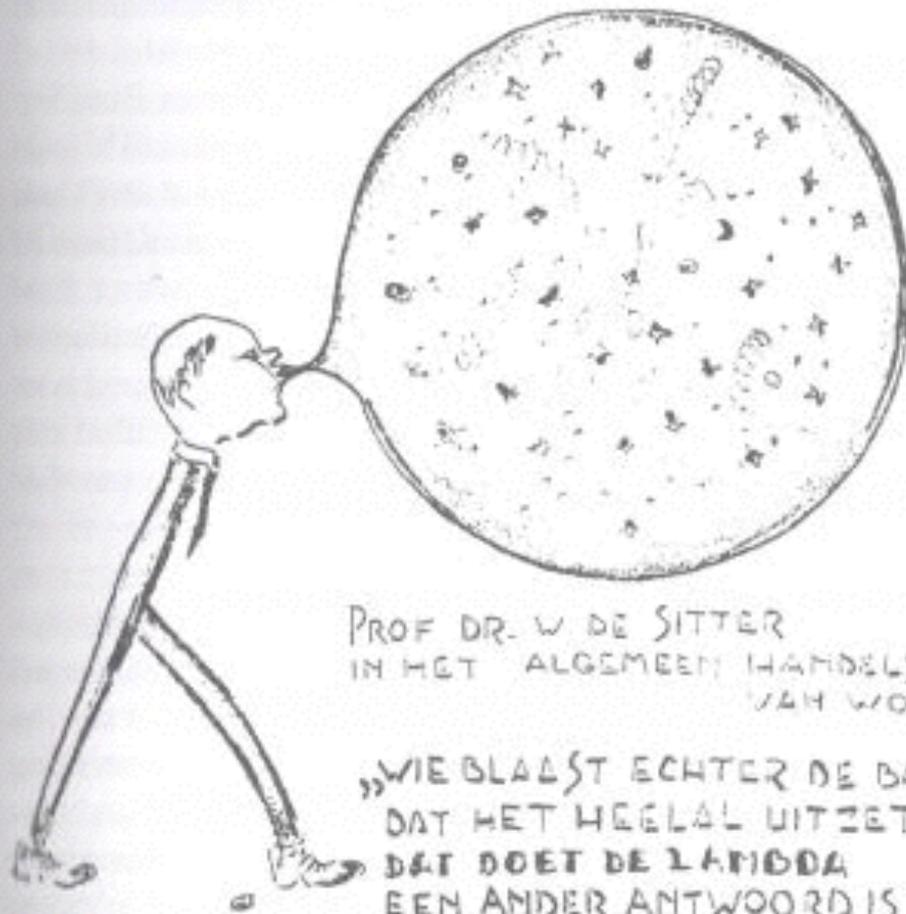
"So, after a few of the most productive hours I had ever spent at my desk, I had learned something remarkable. Would the supercooled phase transition affect the expansion rate of the universe? By 1:00 a.m. I knew the answer: Yes, more than I could have ever imagined."

# Supercooled water

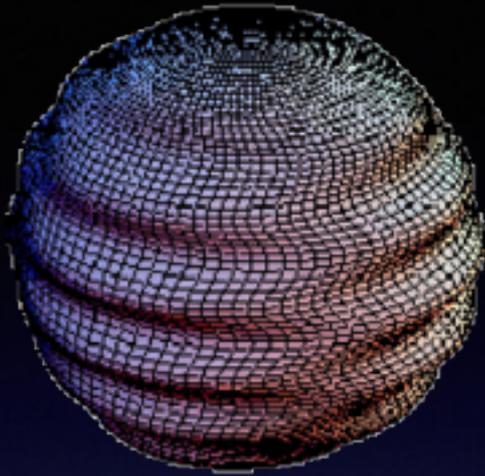


# “World’s Messiest Office” Boston Globe Winner 2005



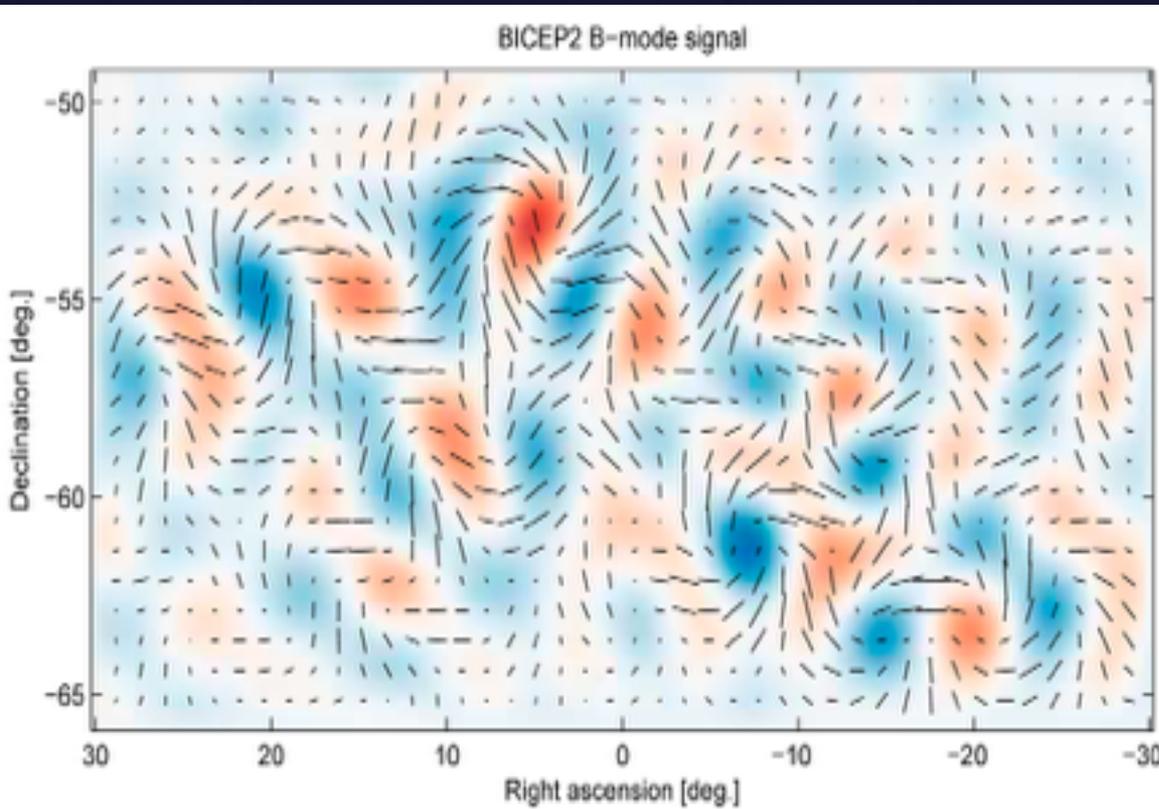


The quote is translated by van der Laan as: "What, however, blows up the ball? What makes the universe expand or swell up? That is done by the Lambda. Another answer cannot be given."



Quantum fluctuations in the gravitational field  
grow to cosmic proportions during inflation  
⇒ **Cosmic-scale gravitational waves**

**Gravitational waves polarize CMB photons**



Strength of polarization  
directly measures

$$U_{\text{inflation}}$$

and by extension

$$E_{\text{inflation}} \propto (U_{\text{inflation}})^{1/4}$$

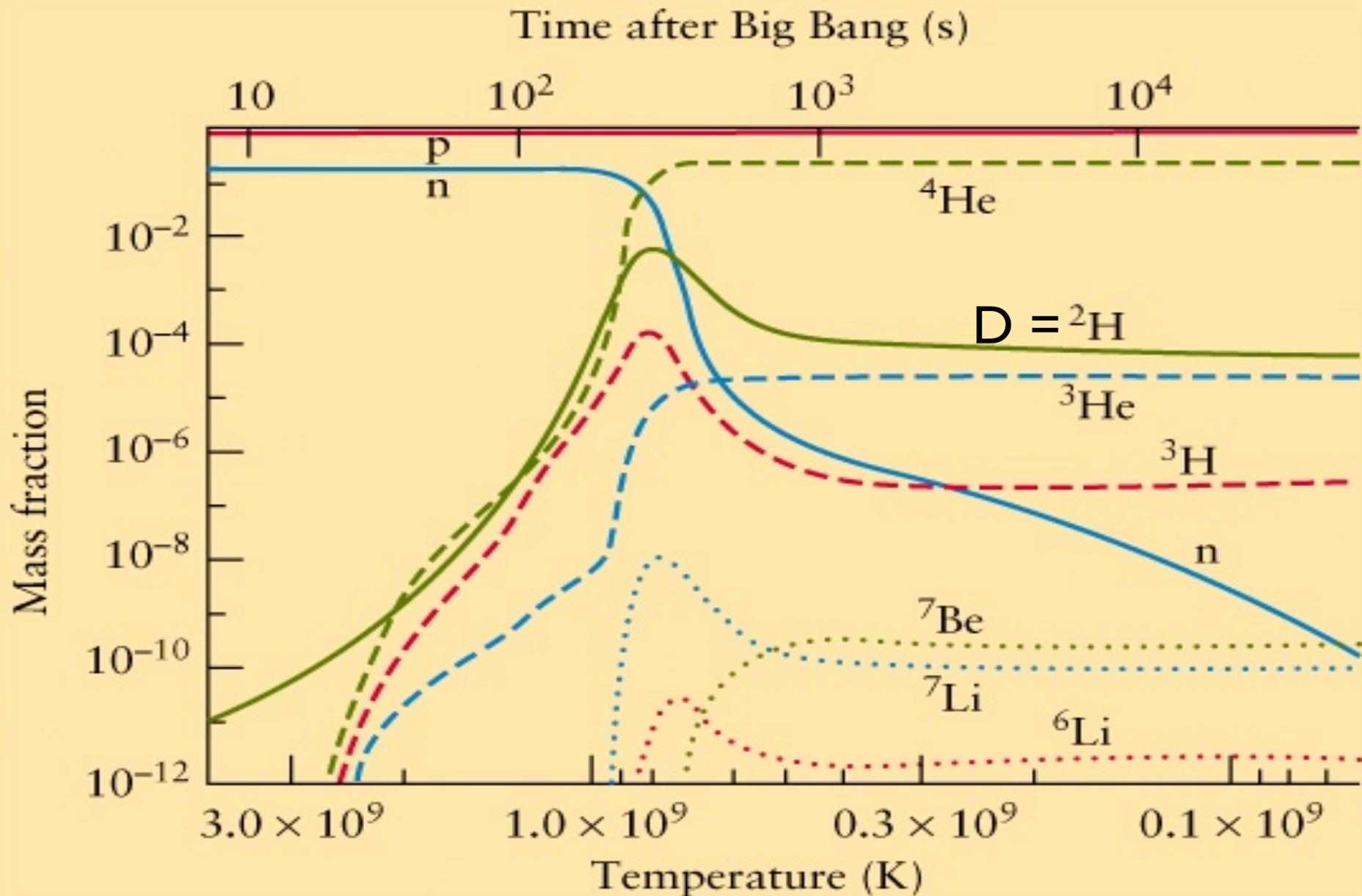
Observed polarization is  
consistent with

$$E_{\text{inflation}} \sim E_{\text{GUT}} \sim 10^{12} \text{ TeV}$$

and

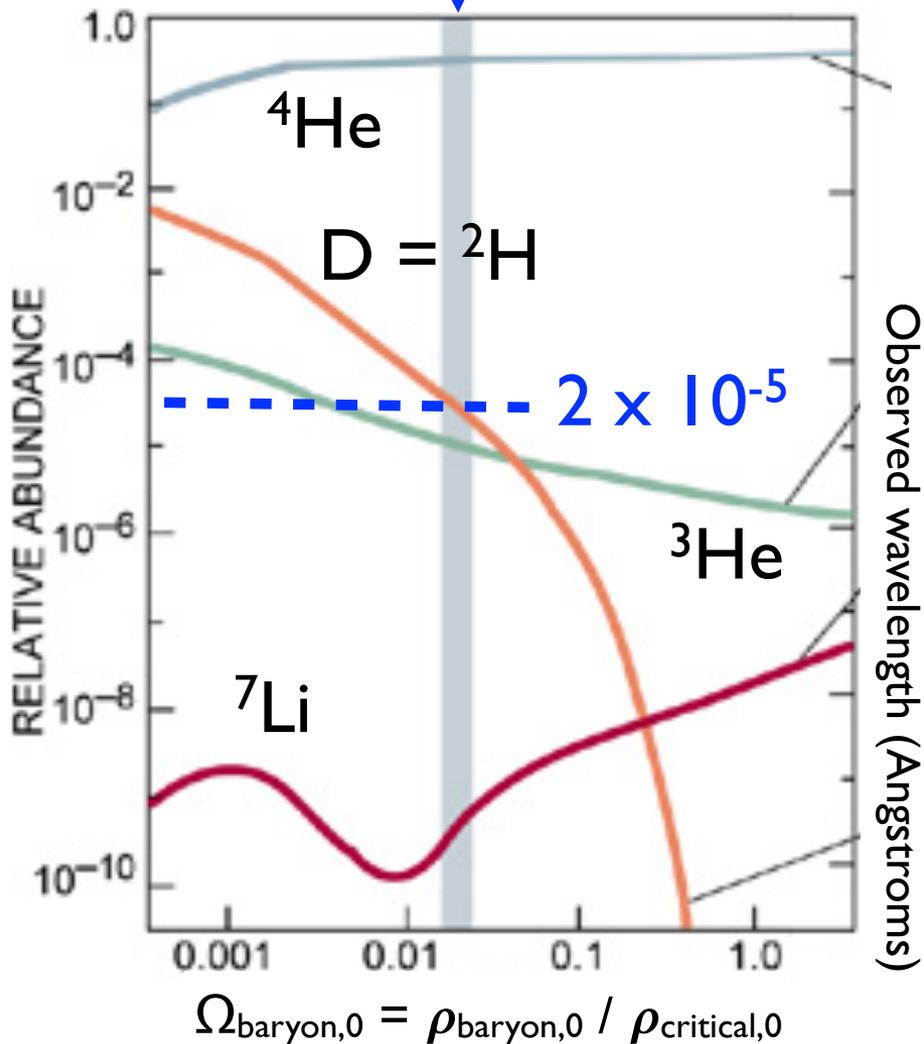
$$\tau_{\text{inflation}} \sim 10^{-36} \text{ s}$$

# Big Bang Nucleosynthesis



More baryons  
 $\Rightarrow$  More fusion reactions  
 $\Rightarrow$  Less deuterium

$\Omega_{\text{baryon},0} = 0.04$



Primordial deuterium abundance measured from Ly- $\alpha$  forest

