Lecture 20. Three Special Molecules: OH, H₂O and NH₃

- 1. Introduction
- 2. OH
- 3. H₂O
- 4. NH₃
- 5. Summary

References

Stahler & Palla, "The Formation of Stars" (Wiley 2004):

Ch. 5 & 6 - Molecular Transitions

Ch. 14 - Masers

Ho & Townes, ARAA, 21, 239, 1983 (NH₃)

Lo, ARAA, 43, 625, 2005 (megamasers)

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1. Introduction

These molecules are of high and varied astrophysical interest. They were discovered in space by Townes and his collaborators at MIT and UC Berkeley.

- OH paramagnetic radical that permits magnetic field measurements chemical precursor to more stable molecules like CO and H₂O first cosmic maser
- H₂O the most interesting molecule in the universe immensely complex with a trillion lines maser
- NH₃ tracer of high density gas and thermometer mysterious chemistry

Basic Properties

Quantity	OH	H ₂ O	NH ₃
<i>IP</i> (eV)	13.02	12.62	10.07
D(eV)	4.411	5.101	4.392
ΔH^{a}	9.32	-57.8	-11.0
pa ^a	141.8	165	204.0
gr. state	$X^{2}\Pi_{3/2,1/2}$	¹ A ₁	X ² A ₂
μ^{b}	1.66	1.85	1.47

- a. Units for chemical energy, kcal/mol: 1 eV = 23.06 kcal/mol
- b. Units for Dipole moment: 1 Debye = 10^{-18} cm

The enthalpy ΔH and proton affinity pa are the energy and the proton binding, as discussed in the next lecture.

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2. The OH Radical

OH has 9 electrons: $(1s\sigma)^2(2s\sigma)^2(2p\sigma)^2(2p\pi)^3$ The unpaired $2p\pi$ electron determines the electronic state:

$$\lambda$$
 = 1 and σ = 1/2 \Rightarrow $^{2}\Pi_{1/2/3/2}$

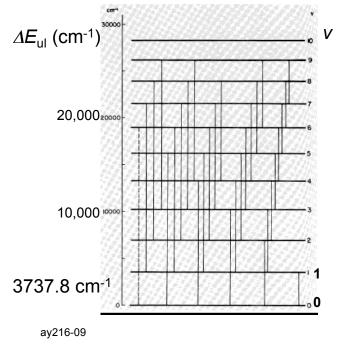
The zero-order vibrational and rotational constants in the ground state are:

 $v = 3738 \text{ cm}^{-1}$ or 2.68 μ m (NIR) $B = 18.91 \text{ cm}^{-1}$ or 119 μ m (FIR)

OH is a diatomic molecule with finite electron angular momentum, i.e., $\lambda = 1$.

Here we sketch the vibrational and rotational levels and mention briefly some astrophysical applications.

OH Ground State Vibrational Levels



The first excited electronic level (not shown) is A $^2\Sigma^+$.

The transition A ${}^2\Sigma^+ \rightarrow X {}^2\Pi_{1/2,3/2}$ occurs in the UV near 3060 Å.

1 cm⁻¹ = 1.4883 K. 1 eV = 11604 K

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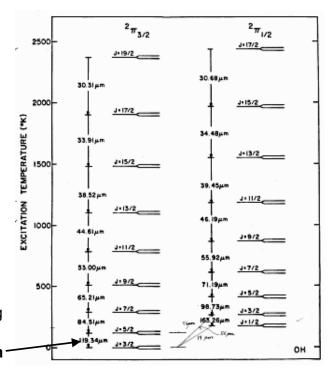
OH Ground State Rotational Levels

The rotational ladders are split by electron spin-orbit coupling.

 λ -doubling arises from the the magnetic interaction of the orbital and rotational motions (not to scale in the figure).

The hfs of the lowest rotational level is shown in the next slide; hfs plus λ -doubling splits the ground state into four.

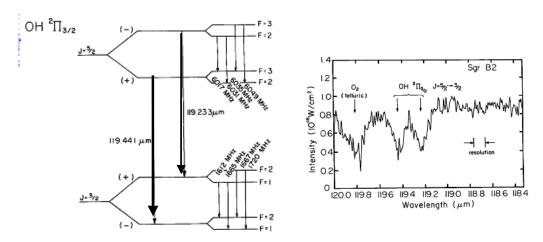
The large rotational spacing Is due to the small mass of the H atom. 119.34 µm



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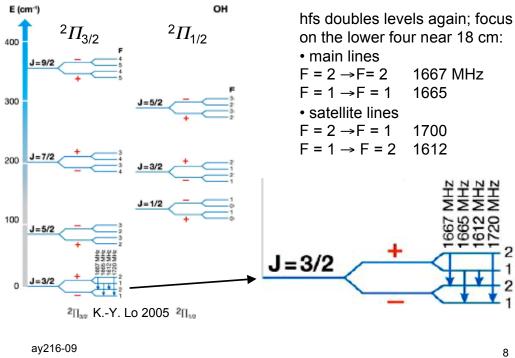
First Detection of OH Rotational Lines



Storey, Watson & Townes (ApJ, 244 L47, 1981) detected the split 119µm line in using the KAO (Kuiper Airborne Observatory)

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OH Hyperfine Masing Levels



Interstellar and Circumstellar OH

Discovered by Weinreb et al. (1963) in absorption - first cosmic radio molecule.

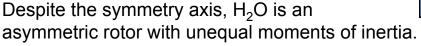
Anomalous emission detected by Weaver et al. (1965) - identified as the first cosmic maser.

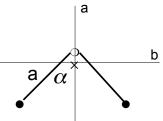
OH is pervasive: It is observed in comets, in the ISM, and as masers around young and old stars.

- OH is strongly paramagnetic. It is used to measure magnetic fields in both atomic and molecular gas (first by Verschuur 1969, and recently in cloud cores by Troland & Crutcher, ApJ, 680, 547, 2008)
- OH masers are often seen at very large distances, ("megamasers"), including near AGN (reviewed by Lo 2005)
- Robishaw et al. (ApJ 680, 981, 2008) measured the Zeeman effect in extragalactic megamasers.
- Robishaw, Heiles, and Crutcher (2009) detected a megamaser in the Cas A SNR.
- The relative abundance of OH to H₂O provides a key test of astrochemistry (it is often much larger than expected).

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3. H₂O





Bond length a = 0.958 ÅBond half-angle $\alpha = 52.25^{\circ}$

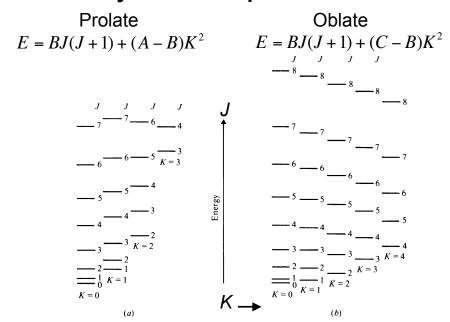
Axis a : symmetry axis Axis b : parallel to the 2H Axis c : perpendicular to plane

Crude estimate:
$$I_a = \cos^2 \alpha \ 2m_H a^2$$

 $I_b = \sin^2 \alpha \ 2m_H a^2$
 $I_c = 2m_H a^2$
 $\sin \alpha = .793$
 $\cos \alpha = .609$

Measured rotational constants:
$$A = 835.783 \text{ GHz}$$
 (27.878 cm⁻¹)
 $B = 435.044 \text{ GHz}$ (14.512 cm⁻¹)
 $C = 278.447 \text{ GHz}$ (9.288 cm⁻¹)

Symmetric Top Recall



Allowed transitions: $\Delta K = 0$, $\Delta J = \pm 1$ (on K ladders)

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Quantum Numbers for Water

- Conserved quantities: total angular momentum and one projection on a space-fixed axis (J,M) plus the energy E.
- Angular momentum components along the moving principal axes are not conserved.
- Water is closer to prolate than oblate.
- An asymmetry parameter is defined as

$$\kappa = \frac{2B - A - C}{A - C}$$

with limiting values of -1 (prolate) and +1 (oblate).

- $\kappa(H_2O) = -0.44$
- The energy varies continuously with κ from -1 to +1

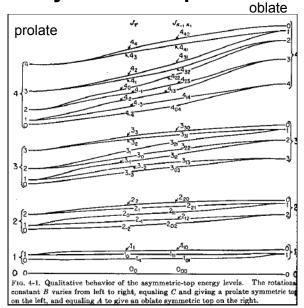
Example: J = 1 allows three states, K = 0, 1 (doubly degenerate). For $\kappa = -1$, the lowest level is (1,0), whereas for $\kappa = 1$, it is (1,1). For intermediate κ , the degeneracy is lifted, and one of the prolate state K = 1 states eventually becomes the oblate K = 0 state.

Levels for the Asymmetric Top

The levels are specified by the limiting K- values, K_p and K_o , or more usually, K_{-1} and K_{+1} .

Molecular spectroscopists use $J_{K-1,K+1}$ to label the rotational states of asymmetric molecules like water, even though K_{-1} and K_{+1} are pseudo quantum numbers.

There is no general closedformed formula for the energy, hence the need for huge line lists.



Levels vs. asymmetry parameter Townes & Schawlow, Fig 4-1

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Role of Nuclear Exchange

Like H_2 , the identity of the protons leads to two distinct families, ortho (spins aligned, I = 0) and para (spins anti-parallel, I = 0),

The exchange properties of the wave function function are similar to those of a symmetric top: $(-1)^{K-1-K+1}$.

Nuclear para (I = 0) have either K_{-1} or K_{+1} odd, not both. Nuclear ortho (I = 1) have K_{-1} and K_{+} both odd or both even.

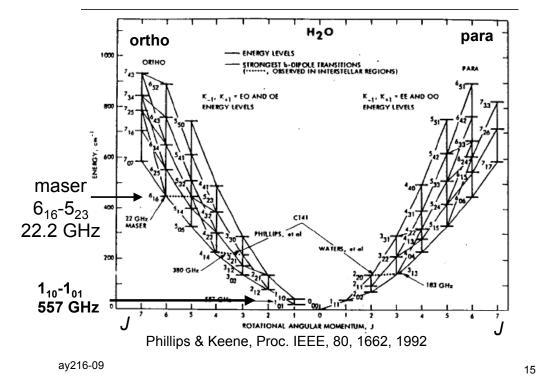
The result is two branches when plotted in an *E-J* diagram (next slide). In this diagram, transitions with both $\Delta K \neq 0$ and $\Delta K = 0$ occur, unlike the case of the symmetric top

The energy diagram is very complicated and becomes even more so when ro-vibrational transitions are included. The most complete line list (BT2) has 0.5 trillion lines:

Barber and Tennyson, MNRAS, 368, 1007, 2000

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Famous Old H₂O Level Diagram



The 22 GHz (1.35 cm) H₂O Megamaser

First observations of galactic H₂O masers Cheung et al. (1969): Sgr B2, Orion and W 49

First *megamaser* detected in M 33 (Churchwell 1977)

Megamasers have luminosities up to ~ 10^6 greater than galactic masers, i.e., ~ 10^2 - 10^4 L_{\odot} .

Thousands of galaxies have been searched; the megamaser detection rate is ~ 5%, with a preference for Seyfert 2 AGN.

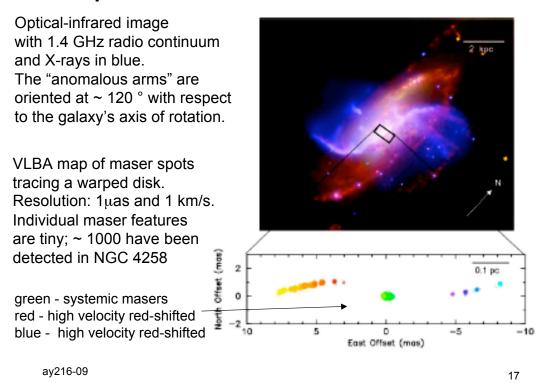
The maser transition is ortho $\mathbf{6_{16}}$ ----- $\mathbf{447.252}$ cm⁻¹ $\mathbf{5_{23}}$ ----- $\mathbf{446.511}$ cm⁻¹

These levels lie 650 K above ground. The maser is believed to be collisionally excited in warm, dense gas

The maser system in NGC 4258 (M 106) is the poster child for megamasers. The masers are in a near edge-on disk that permits clear deduction of many important results. There are only a few others with comparable promise.

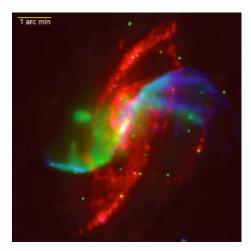
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Composite Photo of NGC 4258 With Masers

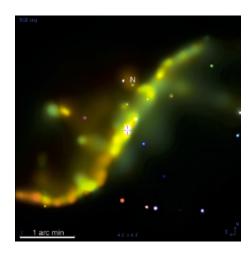


X-ray Observations of NGC 4258 (M 106)

Yang et al. ApJ 660, 1106, 2007



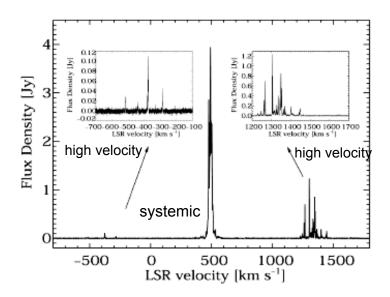
Blue - 1.46 GHz VLA Green - Chandra X-rays Red - Spitzer 8μm



Red - 0.4-0.7 keV Green - 0.7-1.4 keV Blue - 1.4-2.0 keV

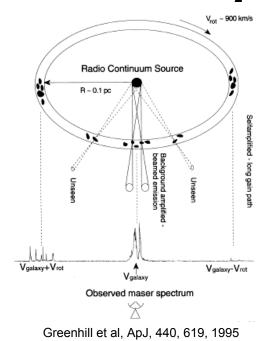
Green Bank Maser Spectrum of NGC 4258

Moran ASPC 395, 87, 2008



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Cartoon of the H₂O Masers in NGC 4258



 $v_K = \sqrt{\frac{GM_c}{R}}$

Keplerian behavior of the high-velocity masers. Moran et al. PNAS 92, 11427, 1995

- Masers detected where velocity gradient is smallest.
- They pass out of sight in ~ 12 yrs; beaming is suggested.
- Amplification is supposed to occur along the line of sight

Dynamical Conclusions on NGC 4258

Thin warped disk, radius \sim 0.1 pc and thickness \sim 40 AU, corresponding to an isothermal atmosphere of \sim 600 K

Central mass based on radius and rotational speed (~ 900 km s⁻¹): $M \sim 2\text{--}3 \times 10^7 M_{\odot}$

Distance of NGC 4258, based on disk velocity model and measured proper motion (displacement angle and time): $d = 7.2 \pm 0.4\,$ Mpc., which is close to latest Cepheid distance of $d = 7.5 \pm 0.2\,$ Mpc, which relies on $d(LMC) = 50\,$ kpc.

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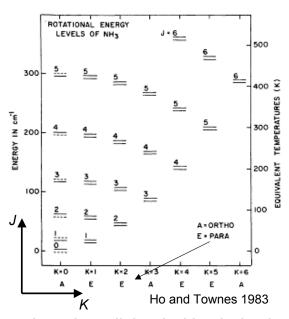
4. NH₃

Pyramid with three H at base and N on top.

Exclusion Principle requires: **ortho** states - all 3 spins aligned **para** states -1 misaligned spin *K*=3n for ortho; otherwise para

Dipole moment aligned with symmetry axis; allowed transitions satisfy

$$\Delta K = 0$$
, $\Delta J = 0, \pm 1$
These rotational transitions are in the FIR near 200 μ m



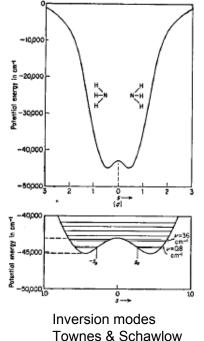
inversion splitting doubles the levels

Inversion Splitting of NH₃



In the ground state, the N atom Is located on either side of the 3 H atoms in the plane. To reach the other side, it has to tunnel through the potential barrier, whose height is ~ 2,000 cm⁻¹.

The tunneling frequency is small (cm band), whereas the allowed rotational transitions are in the farinfrared and require observations from space.



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NH₃ Inversion Thermometer

The inversion splitting of the rotational levels is ~ 25 GHz (~ 1 cm). They are usually observed at the bottom of a K-ladder. The splittings of the (K,K) levels are:

- (1,1)23.694 GHz
- (2,2)23.723
- (3,3)23.870
- (4,4)24.139
- 24.533 (5,5)
- 25.056 (6,6)

The big advantage of the NH₃ inversion transitions is they can be measured with a single telescope, even simultaneously

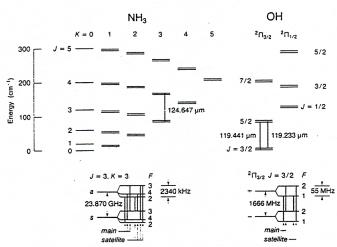
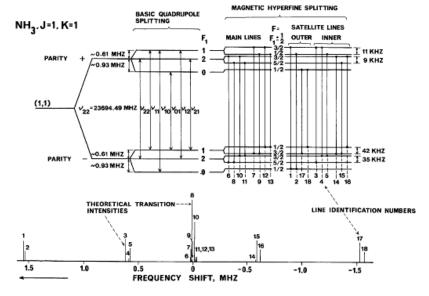


Fig. 5. Rotational levels of NH₃ (left) and OH (right) (adapted from Watson 1982) Inversion splitting for NH3 λ -doubling for OH both with hfs

Hyperfine Structure of the NH₃ (1,1) Level

Complication: the inversion transition is split into 18 hyperfine components

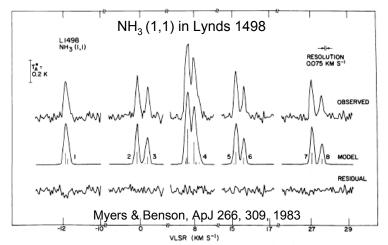


Rydbeck et al. ApJ, 215, L35, 1977

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Measured Hfs of the Inversion Transition



The rotational temperature is obtained by modeling the optical depth of the hfs transitions assuming they are thermalized, Following Barrett et al. (ApJ, 211, L239, 1977) and described in Ho and Townes (1983) and Sec. 6.2 of Stahler and Palla. See also the appendix to Ungerechts et al. A&A, 157, 207, 1986.

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