

# Lecture 25

## Physical Properties of Molecular Clouds

1. Giant Molecular Clouds
2. Nearby Clouds
3. Empirical Correlations
4. The Astrophysics of the X-Factor

### References

- Blitz & Williams, "Molecular Clouds"  
Myers, "Physical Conditions in Molecular Clouds"  
in *Origins of Stars & Planetary Systems* eds. Lada & Kylafis  
<http://www.cfa.harvard.edu/events/1999/crete>  
Bergin & Tafalla, ARAA 45 339 2007 - observations of cold dark clouds  
McKee & Ostriker, ARAA 45 565 2007 - summary of observations and theoretical interpretations

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## 1. Giant Molecular Clouds

- An important motivation for studying molecular clouds is that's where stars form
- Understanding star formation starts with understanding molecular clouds
- In addition to their molecular character, large and massive molecular clouds are dynamical systems that are

**Self-Gravitating  
Magnetized  
Turbulent**

- The central role of gravity high distinguishes them from other phases of the ISM.

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## What is a Molecular Cloud?

- Molecular clouds have dense regions where the gas is primarily molecular.
- *Giant molecular clouds* (GMCs) are large clouds with  $10^4 M_{\odot} < M < 6 \times 10^6 M_{\odot}$  sizes in the range 10-100 pc.
- The filling factor of GMCs is low; there about 4000 in the Milky Way). They have as much atomic as molecular gas.
- Mean densities are only  $\sim 100 \text{ cm}^{-3}$ , but molecular clouds are inhomogeneous and have much higher-density regions called clumps and cores.

NB There is no accepted explanation for the sharp upper limit to the mass of GMCs; tidal disruption and the action of massive stars have been suggested.

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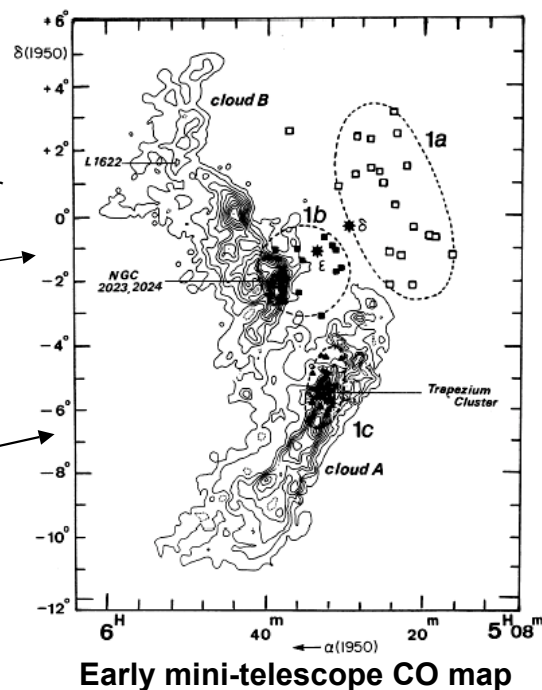
## 2. The Orion Molecular Cloud Complex

These clouds can't be much older than 10-20 Myr, the age of the oldest OB sub-association.

Associations older than 20-30 Myr are not associated with GMCs.

Cloud B

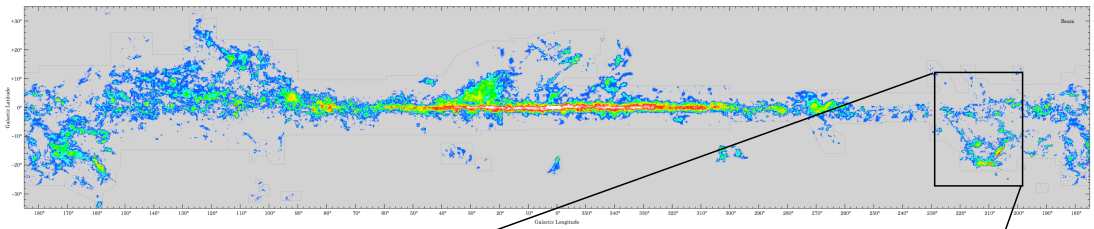
Cloud A



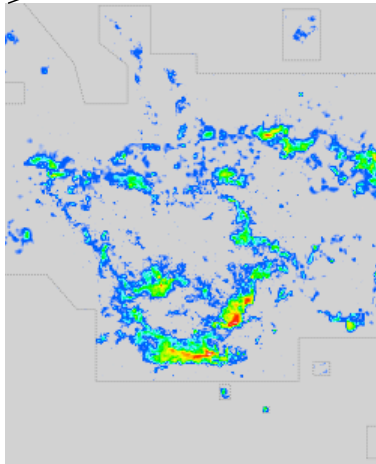
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# Orion: The Very Large Scale Picture

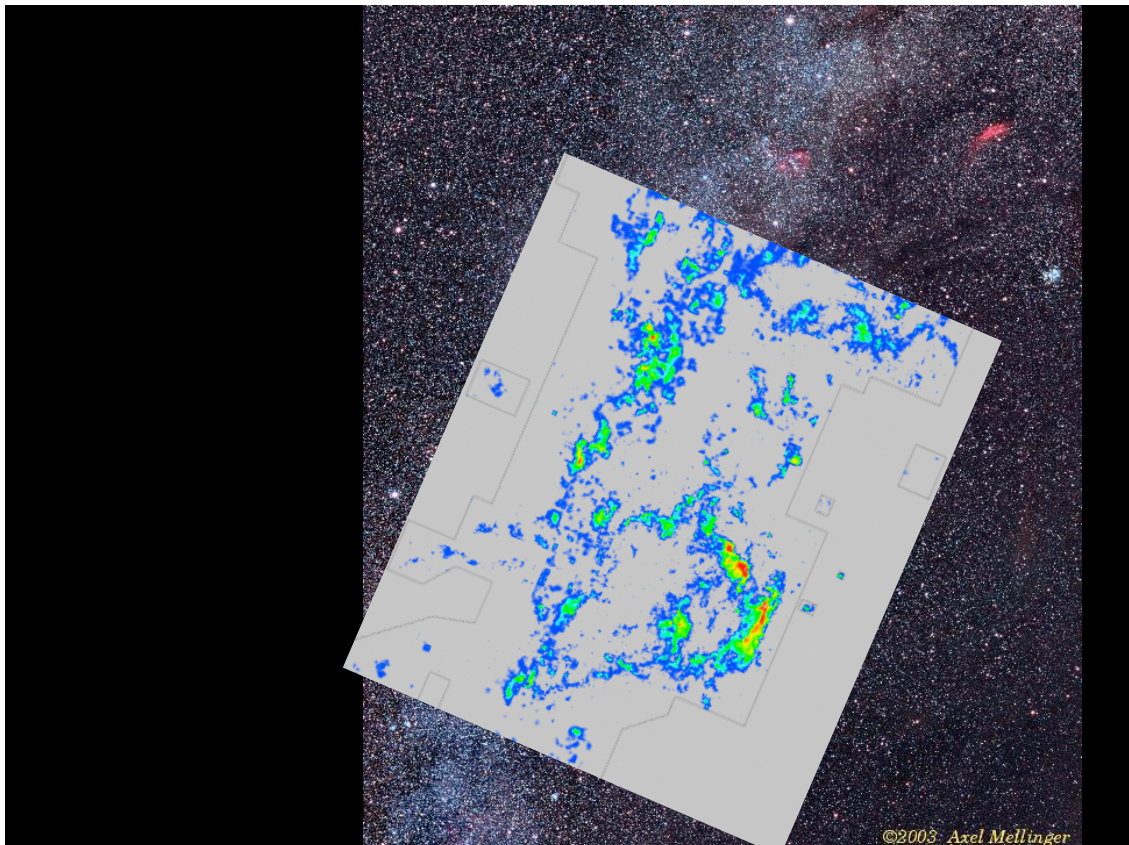


Dame et al. (2001)  
CO survey



See the next slides with stars.

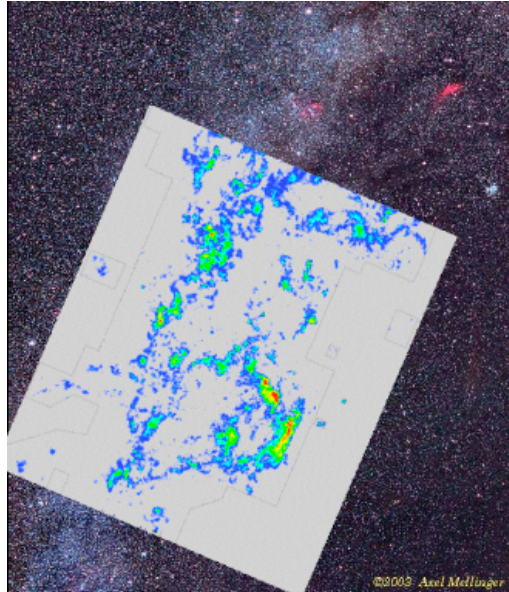
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## Large-scale Optical and CO Images



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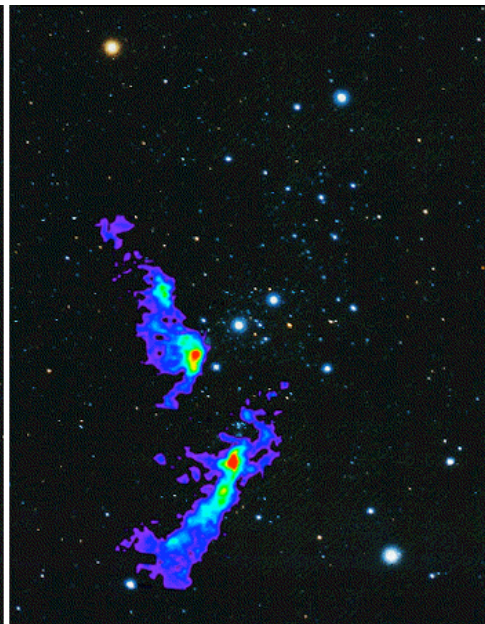


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## Orion Molecular Clouds A and B in CO Constellation Scale Optical and CO Images

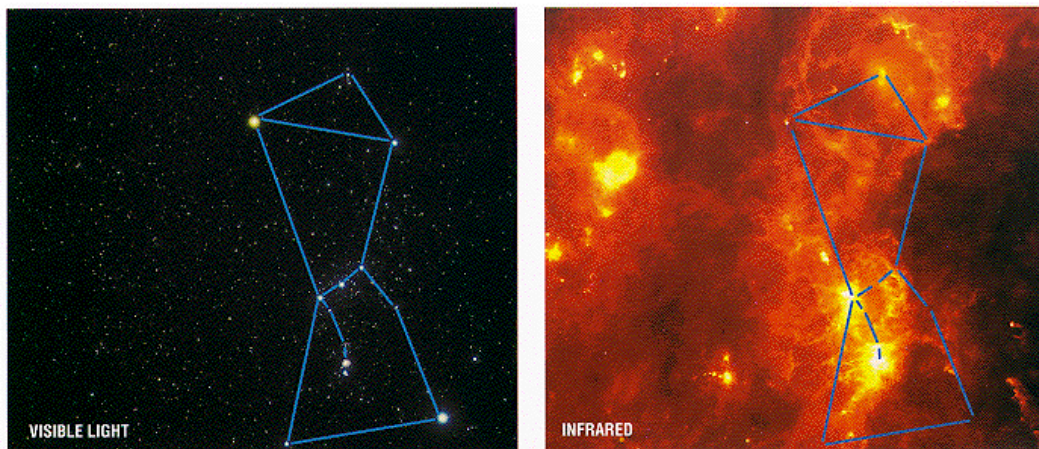


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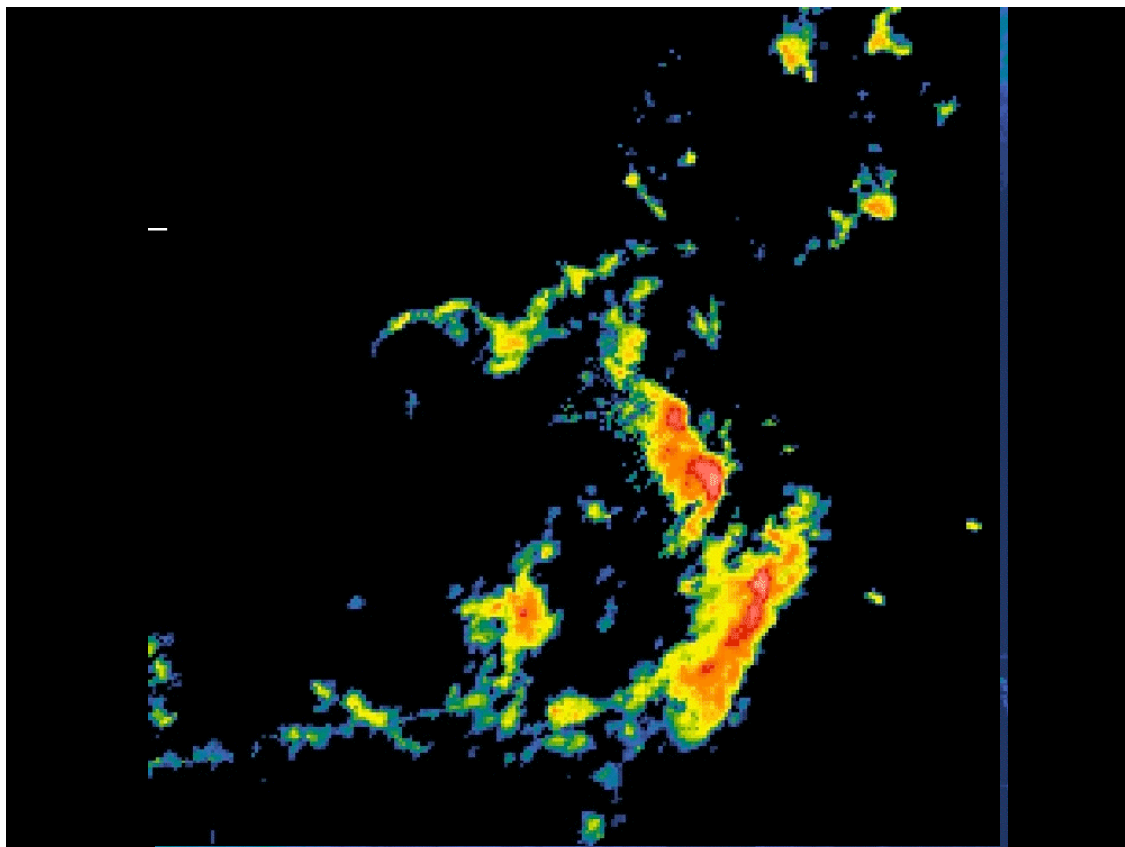
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## Orion Molecular Clouds A and B in IR Constellation Scale Optical and IRAS Images



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## Summary for Orion GMCs

- Cloud A (L1641) exhibits typical features of GMCs:
  - fairly well defined boundaries: GMCs seem to be discrete systems
  - clumpy, but with unit surface filling factor in optically thick  $^{12}\text{CO}$  1-0 in low resolution maps
  - elongated, parallel to the plane of the Galaxy
  - strong velocity gradient (rotation)
- Star clusters form in GMCs
  - no local GMCs ( $d < 1$  kpc) without star formation
  - one nearby GMC ( $d < 3$  kpc) without star formation (Maddalena's cloud  $\sim 10^5 M_{\odot}$ )

***Essentially all star formation occurs in molecular clouds***

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### 3. Basic Properties of Molecular Clouds

- Important deductions can be made from CO studies of molecular clouds by very direct and simple means.
- The relevant data are the line width, the integrated line strength and the linear size of the cloud.

For a Gaussian line, the variance or dispersion  $\sigma$  is related to the Doppler parameter  $b$  and the FWHM as follows:

$$\phi(v) = \frac{1}{\sqrt{2\pi}\sigma} e^{-v^2/2\sigma^2}$$

$$\sigma = b/2^{1/2}, \quad FWHM = 2 \sqrt{2 \ln 2} \sigma \approx 2.355 \sigma$$

For thermal broadening,

$$b_{\text{th}} \approx 0.129 (T/A)^{1/2} \text{ km s}^{-1} \quad (A = \text{atomic mass}).$$

More generally, in the presence of turbulence,

$$\sigma^2 = \frac{kT}{m} + \sigma_{\text{turb}}^2$$

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## Application of the Virial Theorem

A key step in the elementary interpretation of the CO observations by Solomon, Scoville, Sanders et al. uses the virial theorem, which assumes that

**GMCs are gravitationally bound in virial equilibrium,**

The virial theorem with only gravitational forces reads:

$$-\langle V \rangle = 2\langle K \rangle = \langle mv^2 \rangle \quad \text{or} \quad \left\langle \frac{GM}{R} \right\rangle = \langle v^2 \rangle = \sigma^2$$

Measurements of the size  $R$  and the velocity dispersion  $\sigma$  can then be used to estimate the mass of the GMC:

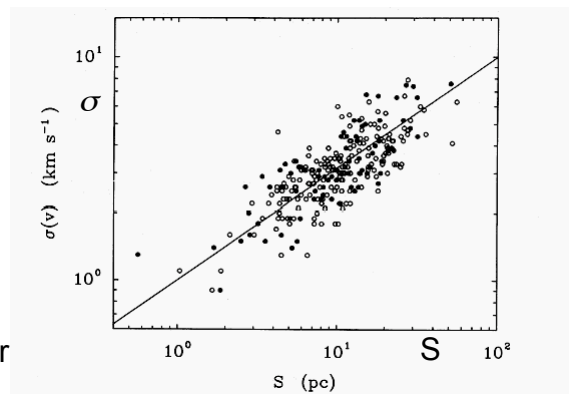
$$M \approx \frac{R\sigma^2}{G}$$

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## The Linewidth-Size Correlation

- $T_{\text{kin}} \sim 20 \text{ K} \Rightarrow \sigma < 0.1 \text{ km/s}$  (from low- $J$  CO lines)
- Linewidths are suprathermal
- Noticed by Larson (MNRAS 194 809 1981), who fitted  $\sigma \sim S^{0.38}$  close to Kolmogorov 1/3.
- Others found  $\sigma \sim S^{0.5}$  ( $\sigma$  in  $\text{km s}^{-1}$  and  $S$  in pc).
- The correlation extends to smaller clouds and smaller length scales within GMCs (Heyer & Brunt, ApJ 615 L15 2004), but not to cores
- If the linewidth is a signature for turbulence\*, this correlation is an **empirical statement about turbulence in molecular clouds.**



**Linewidth-size correlation**  
for 273 molecular clouds  
Solomon et al. ApJ 319 730 1987

$$\sigma = (0.72 \pm 0.03) \left( \frac{R}{\text{pc}} \right)^{0.5 \pm 0.05} \text{ km s}^{-1}$$

\* For an introduction to interstellar turbulence, see Sec 2. McKee & Ostriker (2007)

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## The Luminosity-Mass Correlation

$$I_{CO} = \int_{line} T_A(v) dv$$

is the line integrated intensity for optically thick  $^{12}\text{CO}$ .

The CO luminosity of a cloud at distance  $d$  is

$$L_{CO} = d^2 \int_{cloud} I_{CO} d\Omega ; \quad \text{hence} \quad \boxed{L_{CO} \approx T_{CO} \Delta v \pi R^2}$$

where  $T_{CO}$  is the peak brightness temperature,  $\Delta v$  is the velocity line width and  $R$  is the cloud radius.

Substituting  $\Delta v^2 \approx \frac{GM}{R}$  (virial equilibrium) and  $M = \frac{4\pi}{3} \rho R^3$

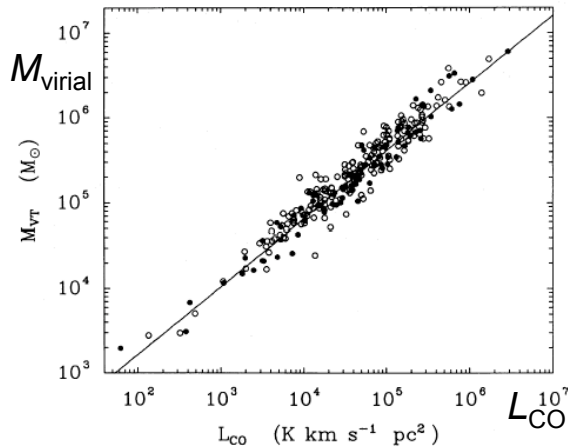
yields

$$\boxed{L_{CO} \approx \sqrt{3\pi G/4\rho} T_{CO} M}$$

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## The Mass-CO Luminosity Correlation



Solomon, Rivolo, Barrett & Yahil  
ApJ 319 730 1987

The good correlation over 4 dex supported the assumption that GMCs are in virial equilibrium.

$M_{\text{virial}}$  may be an underestimate because it is based on optically thick CO. GMCs have diffuse regions that are not optically thick. And there are observational problems as well.

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## c. Correlations

The two observationally based correlations for GMCs are:

$$\sigma \approx R^{1/2} \quad (\text{line width size relation})$$

$$\frac{M}{R} \approx \sigma^2 \quad (\text{virial equilibrium})$$

Substitution leads to another

$$N \approx \frac{M}{R^2} \approx \frac{\sigma^2}{R} \quad (\text{constant surface density})$$

and as well two more

$$\rho \approx \frac{M}{R^3} \approx \frac{\sigma^2}{R^2} \approx \frac{1}{R} \quad \text{and} \quad M \approx \sigma^2 R \approx R^2 \approx \sigma^4$$

The first three are often referred to as Larson's Laws

**Is it really true that the surface densities of GMCs are all about the same?** Many have so assumed following Solomon et al.:

$$N_H \sim 1.5 \times 10^{22} \text{ cm}^{-2}$$

$$A_V \sim 10$$

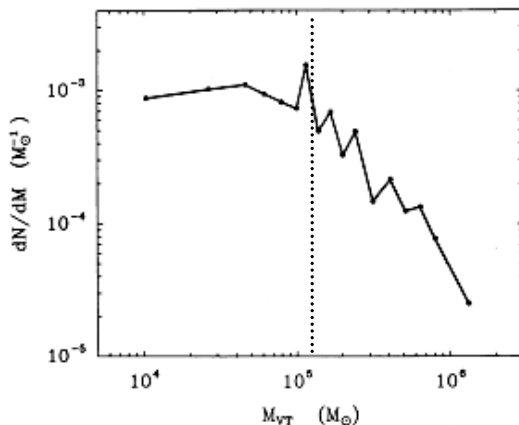
$$\Sigma \sim 150 M_\odot \text{ pc}^{-2}$$

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## GMC Mass Spectrum

Solomon et al. ApJ 319 730 1987



- The spectrum is incomplete for  $M < 10^5 M_{\text{sun}}$  (dotted line).
- $dN/dM \propto M^{-3/2}$  for large  $M$

To be addressed later:

1. What is the mass spectrum for clumps and cores?
2. How are cloud mass functions related to the stellar initial mass function (IMF)?

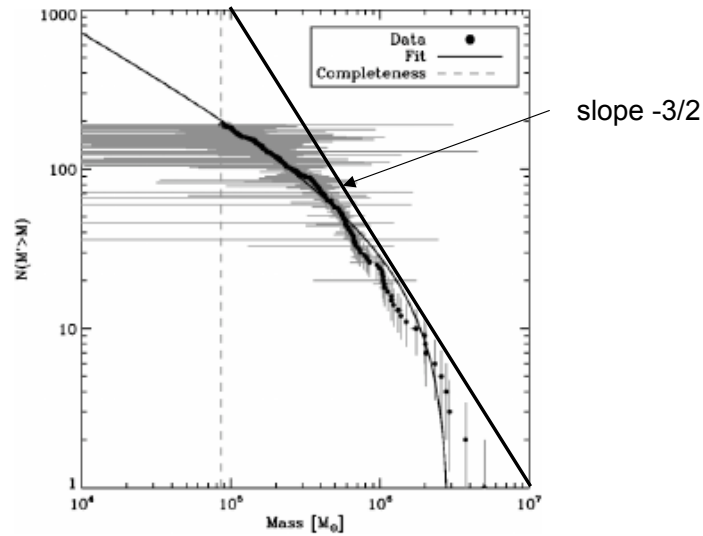
FIG. 3.—The molecular cloud mass spectrum  $dN/dM$ . A fit to the data above  $M = 7 \times 10^4 M_\odot$  gives  $dN/dM \propto M^{-3/2}$ . There are 15 clouds in each bin and the standard deviation is  $\pm 24\%$ . The turnover at low mass is due to undercounting of smaller clouds in the more distant parts of the galactic disk.

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# Reanalysis of Solomon et al. (1987)

Rosolowsky PASP 117 1403 2005



There is a sharp cutoff at  $M = 3 \times 10^6 M_{\text{sun}}$

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## Typical Properties of Local GMCs

Based on Solomon et al. (1987)

Number	4000
Mass	$2 \times 10^5 M_{\odot}$
Mean diameter	45 pc
Projected surface area	$2000 \text{ pc}^2$
Volume	$10^5 \text{ pc}^3$
Volume density ( $\text{H}_2$ )	$300 \text{ cm}^{-3}$
Mean mass surface density	$1.5 \times 10^{22} \text{ cm}^{-2}$
Number surface density	$4 \text{ kpc}^{-2}$
Mean separation	500 pc

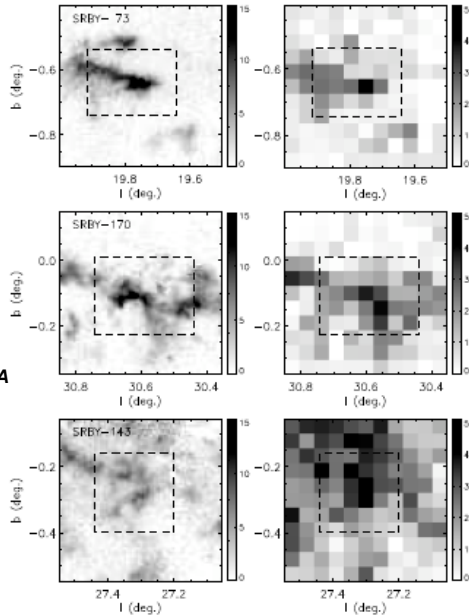
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# Re-examining Larson's Laws

Based on the UMASS-BU Galactic Ring Survey of  $^{13}\text{CO}$  (1-0)  
 Heyer et al. arXiv:0809:1397v1

Left side: UMass-BU  $^{13}\text{CO}$  (1-0) maps  
 Heyer et al 2009



Right side: UMass-SUNY  $^{12}\text{CO}$  (1-0) defined GMCs (dashed lines)  
 Solomon et al 1987

New data:

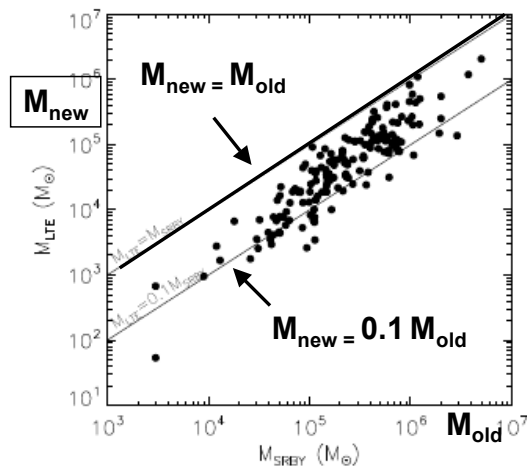
- better resolution
- densely sampled
- more optically thin
- sensitive to lower  $T_A$

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## NEW AND OLD GMC MASSES

Using Solomon definition of GMCs



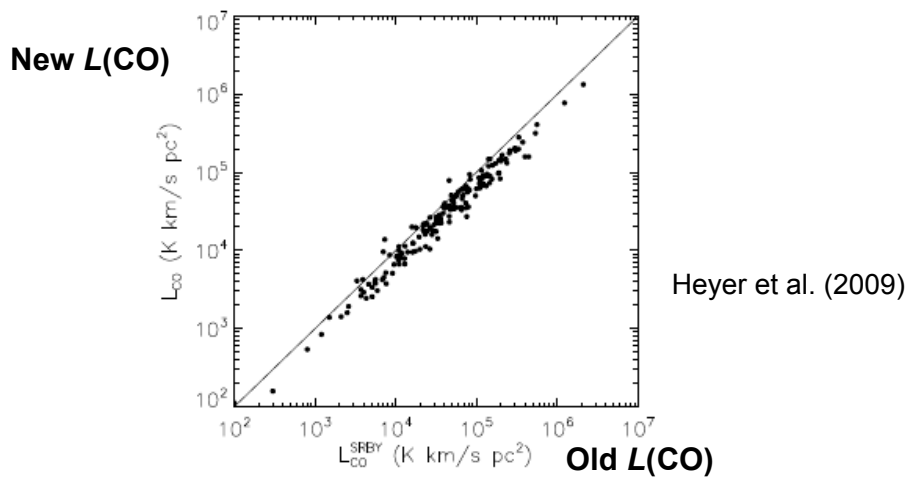
Old = Solomon et al. (1987)      New = Heyer et al. (2009)

**New GMC masses are ~ factor of 5 smaller than the old virial theorem masses**

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# NEW AND OLD CO LUMINOSITIES



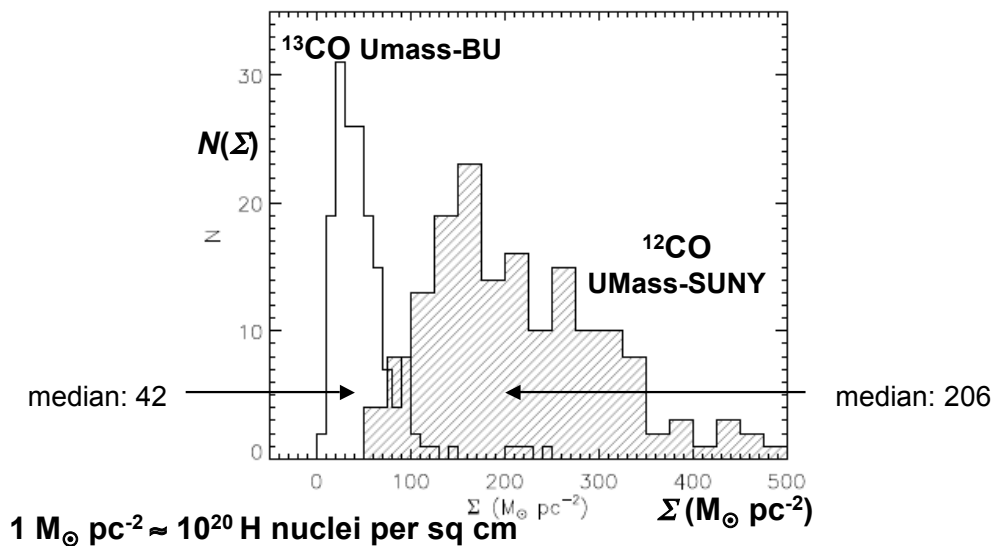
~ 50% overestimate by Solomon et al. indicates that their extrapolation below the detectable brightness temperature over-estimates the luminosity

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# Surface Density Distribution of GMCs

Heyer et al. (2009)



Not only is there is a factor of 5 difference in the medians, but *GMCs do not all have the same surface densities.*

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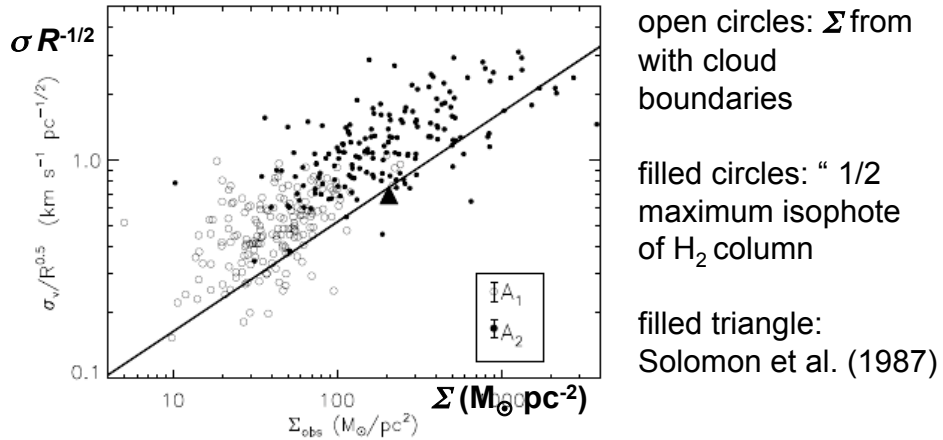
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## Alternate Approach to Correlations

Start with the virial mass relation and the definition of surface density (with  $N = \Sigma$ ), rather than with Larson's law, and then:

substitute  $N \approx \frac{M}{R^2}$  into  $\frac{M}{R} \approx \sigma^2$  to get :  $\sigma \approx N^{1/2} R^{1/2}$

This indeed is what's observed:



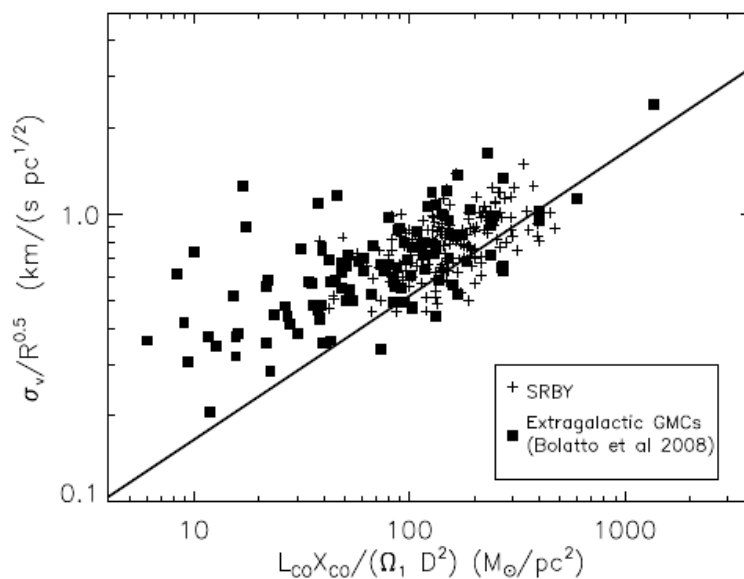
This is not the universal scaling law indicative of turbulence.

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## Similarity of the Extragalactic Correlation

Bolatto et al. IAU Symposium 255 274 2008



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## Understanding GMC Masses and Linewidths

### 1. Observe with better resolution, sampling, and sensitivity.

See Goldsmith et al. ApJ 680 428 2008 for a 20", 32 pixel focal plane array study of the TMC, analyzed with a variable CO abundance model for diffuse regions. They obtain twice the mass compared to the fixed abundance model, with half the mass in diffuse regions.

### 2. Observe the HI with comparable resolution.

### 3. Observe and include magnetic fields and other measures of the velocity field in the analysis

### 4. The origin of the supersonic linewidths seen in GMCs

If it is not hydrodynamic turbulence, is it magnetic?

- We show in Lecture 27 that the magnetic virial theorem gives  $M \sim BR^2$  or  $\Sigma \sim B$ .
- If the linewidths come from Alfvén waves,  $\sigma^2 \sim B^2 / \rho$ .
- Replace  $\rho$  by  $M / R^3$  and use  $M \sim BR^2$  to get  $\sigma^2 \sim RB$ , or  
$$\sigma \sim \Sigma^{1/2} R^{1/2}.$$

This is Heyer's result which he ascribes to Mouschovias (1987).

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## 4. The CO / H<sub>2</sub> Conversion Factor

- Measuring the CO mass or column density is *not* the same as measuring the total gas, which is dominated by H<sub>2</sub> and He and are effectively invisible in cool clouds.
- The integrated CO intensity  $I_{\text{CO}} = \int T_A(\nu) d\nu$  can be calibrated to yield the average H<sub>2</sub> column density. This is surprising because <sup>12</sup>CO is optically thick and because the CO / H<sub>2</sub> ratio might be expected to vary within a cloud and from cloud to cloud.
- It is surprising that a single conversion factor between H<sub>2</sub> column density and  $I_{\text{CO}}$  (the **X-factor**) applies on average to all molecular clouds in the Galaxy.
- That several calibration methods agree to within factors of a few should provide insights into the properties of the clouds.

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## X-factor Method 1: $I_{CO}$ and Virial Theorem

- Measured line intensity:  $I_{CO} \equiv I(^{12}CO) \approx \langle T_A \rangle \Delta v_{FWHM}$

- Virial theorem:  $\frac{GM}{R} \approx \sigma^2 = \left(\frac{\Delta v}{2.35}\right)^2$

- Mass estimate:

$$M = \frac{4\pi}{3} R^3 n(H_2) m \quad \text{and} \quad N(H_2) = \left(\frac{4\pi}{3}\right)^{-1} \frac{M/m}{R^2}$$

- $\Delta v_{FWHM} = 2.35 \sigma \sim (GM/R)^{1/2}$

$$\frac{N(H_2)}{I_{CO}} \approx 3 \times 10^{20} \text{ cm}^{-2} \text{ K}^{-1} \text{ km s}^{-1} \frac{10\text{K}}{T} \left(\frac{n(H_2)}{1000 \text{ cm}^{-3}}\right)^{1/2}$$

### Problems:

Assumes virial equilibrium

Depends on  $n(H_2)$  and  $T$

Measures only mass within  $\tau = 1$  surface

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## X-factor Method 2: $I_{CO}$ and NIR Extinction

- Measure  $I_{CO}$  for regions with high  $A_V$
- Determine  $A_V$  from IR star counts
- Extrapolate  $N_H/A_V$  from diffuse clouds
- Assume all hydrogen is molecular

Result:

$$N(H_2) / I_{CO} \approx 4 \times 10^{20} \text{ cm}^{-2} / (\text{K km s}^{-1})$$

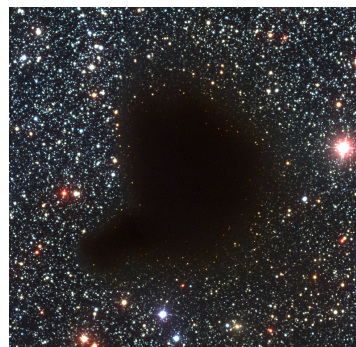
Problems:

Inaccuracies in star-count  $A_V$

Variable dust properties

Variable  $N_H / A_V$

Best for dark clouds



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Lada et al.  
ApJ 586 286 2003

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### X-factor Method 3: $I(^{13}\text{CO})$ vs. $A_V$

- Determine  $A_V$  as in method 2
- Measure  $^{13}\text{CO}$  line intensity
- Assume  $^{13}\text{CO}$  optically thin,  $^{12}\text{CO}$  optically thick
- Assume  $T_{\text{ex}}(^{13}\text{CO}) = T_{\text{ex}}(^{12}\text{CO})$
- Assume  $^{12}\text{CO}/^{13}\text{CO} \approx 40 \dots 60 \Rightarrow \tau(^{13}\text{CO}) \Rightarrow N(^{13}\text{CO})$

Problems:

Accuracy of  $A_V$  determination

Often  $T_{\text{ex}}(^{13}\text{CO}) < T_{\text{ex}}(^{12}\text{CO})$

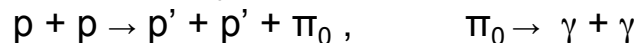
$^{13}\text{CO}$  may not be optically thin

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### X-factor Method 4: $I_{\text{CO}}$ and $\gamma$ -Rays

- High energy cosmic rays ( $> 1$  GeV) produce neutral pions in collisions with protons in H and  $\text{H}_2$ , which then decay into two  $\gamma$ -rays



- The  $\gamma$ -ray emission depends on the product of the cosmic ray density and the density of all protons ( $n_{\text{H}}$ ).
- Hunter et al. ApJ 481 205 1997 combine  $\gamma$ -ray measurements from COMPTON/EGRET with the Columbia-CfA CO survey and obtain,

$$N(\text{H}_2)/I_{\text{CO}} = (1.56 \pm 0.05) \times 10^{20} \text{ cm}^{-2} \text{ K}^{-1} \text{ kms}^{-1},$$

presumably assuming all hydrogen is molecular.

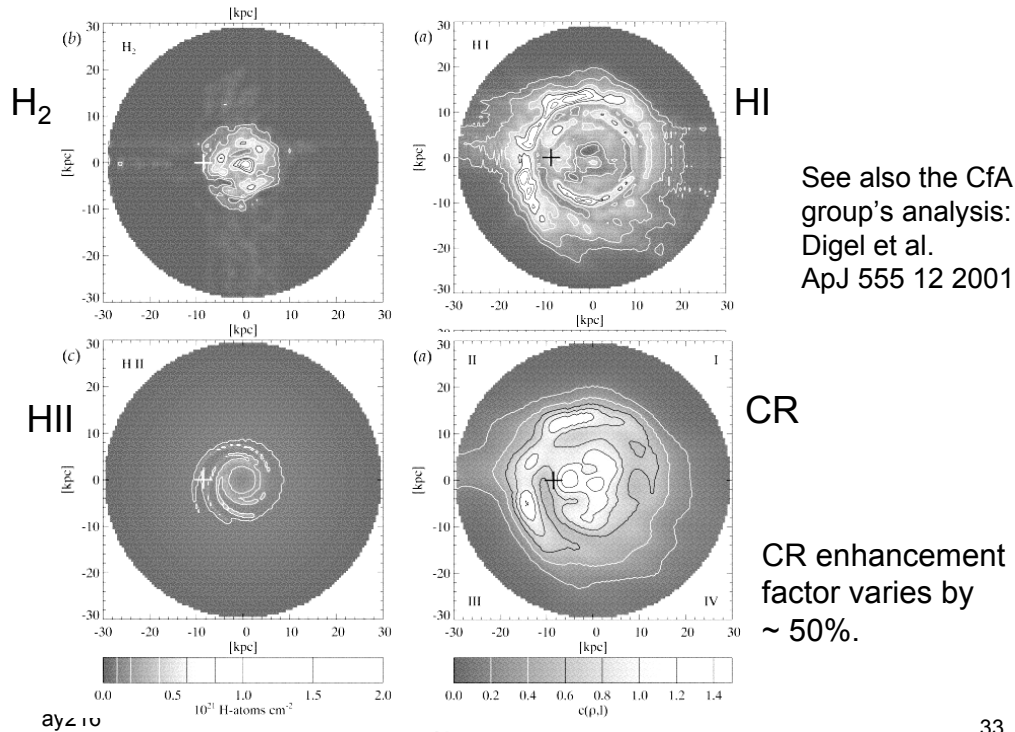
**NB** The modulation correction for high energy CRs is small.  
Hunter et al. assume that the CR density is proportional to  $n_{\text{H}}$ .

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## Hunter et al. ApJ 481 205 1997



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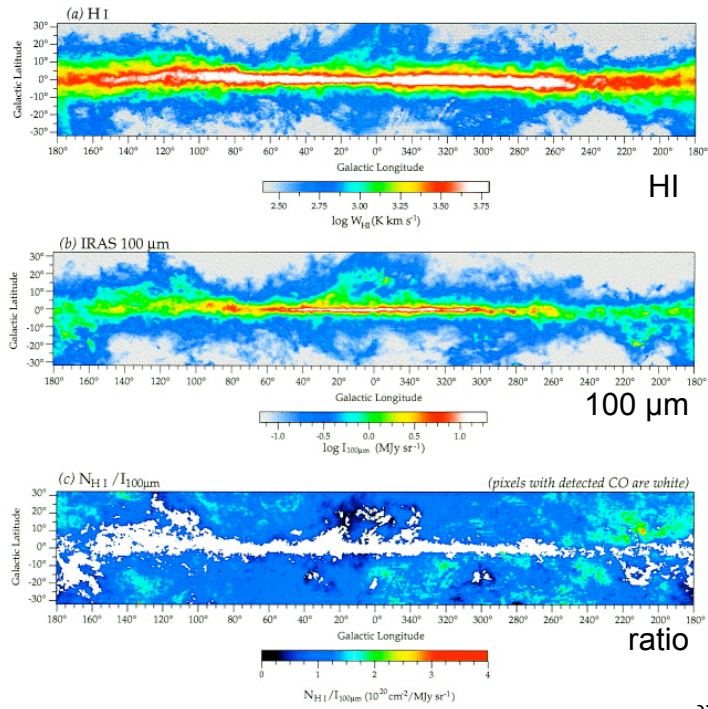
## X-factor Method 5: HI/IRAS/CO

- Dame et al. (ApJ 547 792 2001) used IRAS far-IR emission as a tracer of total gas column density
- Calibrated with the Leiden-Dwingeloo 21-cm HI survey in regions free of CO emission
- Total gas map differenced with the HI map to obtain a complete and unbiased predicted map of H<sub>2</sub>
  - Close agreement between this map and observed CO implies that few molecular clouds at  $|b| < 30^\circ$  have been missed by CO surveys
- The ratio of the observed CO map to the predicted molecular map provides a measure of the local average X-factor for  $|b| > 5^\circ$ :

$$N(\text{H}_2)/I_{\text{CO}} = 1.8 \pm 0.3 \times 10^{20} \text{ cm}^{-2} \text{ K}^{-1} / \text{km s}^{-1}$$

# Method 5: HI/IRAS/CO

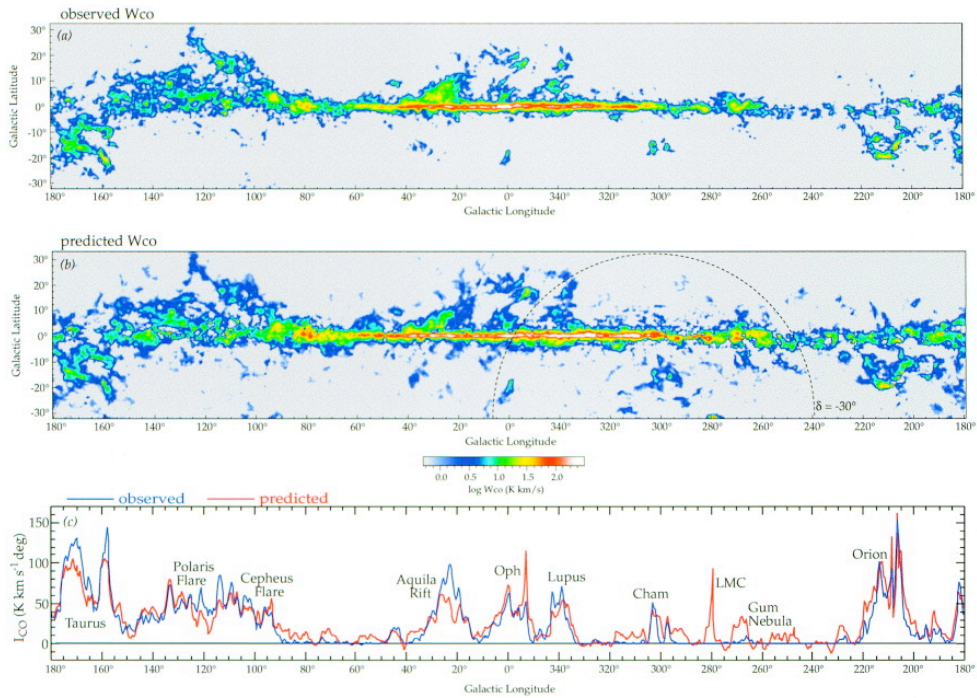
Dame et al. compared IRAS far-IR (dust). 21 cm (HI) and 2.6 mm (CO).



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# Verification of Method 5

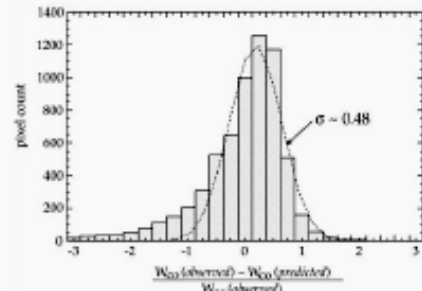
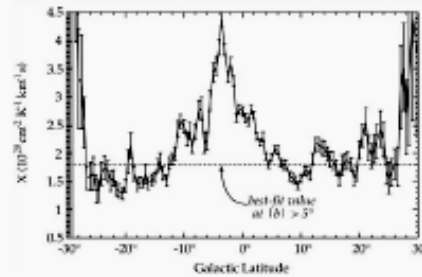


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## James Graham's Critique of Method 5

- Average  $X$  varies with latitude
  - High  $X$  at  $l \sim 0^\circ$  may be spurious since, the lack of CO-free regions toward the inner plane mean  $I_{100}/N_{\text{tot}}$  cannot be properly determined
- Point-to-point dispersion is significantly larger than can be accounted for by instrumental noise
  - Excluding the plane ( $|b| < 5^\circ$ ), where the prediction is expected to break down owing to dust temperature variations along the line of sight, the dispersion is  $\approx 50\%$
  - The high dispersion may be due to variations in the gas-to-dust ratio, and by dust temperature variations not accounted for by the simple IRAS color correction



c.f. JRG Lecture 18 (2006)

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## CO/H<sub>2</sub> Conversion Factors: Summary

- Various methods agree remarkably well
- Relevant on global scales, not locally
- Limits on applicability are unclear
- No information on  $N(\text{H}_2) / N(\text{CO})$  is obtained
- Conversion factors should depend on  $T$ ,  $n$  and metallicity
- Conversion factor derived for Milky Way disk is not valid for galactic nuclei (including our own Galactic Center) or for metal-poor systems
- Blitz et al. (PPV) find that  $X_{\text{CO}} \approx 4 \times 10^{20} \text{ cm}^{-2} (\text{K km s}^{-1})^{-1}$  holds approximately for members of the local group, but not the SMC, where  $X_{\text{CO}} \approx 13.5 \times 10^{20} \text{ cm}^{-2} (\text{K km s}^{-1})^{-1}$ . The conversion for the LMC,  $X_{\text{CO}} \approx 9.0 \times 10^{20} \text{ cm}^{-2} (\text{K km s}^{-1})^{-1}$ , also reflects the reduced abundances of the clouds .

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## CO/H<sub>2</sub> Conversion Factor: Summary

Source	X
Early work	2-5
$\gamma$ -rays (Hunter et al. 1997)	$1.56 \pm 0.05$
HI/IRAS/CO (Dame et al. 2001)	$1.8 \pm 0.3$
IR extinction (Lada et al. 2003)	$\sim 4$

Units for X:  $10^{20} \text{ cm}^{-2} / \text{K km s}^{-1}$